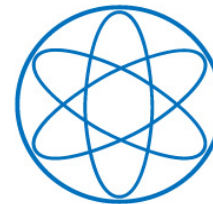


Current Status of Dark Matter

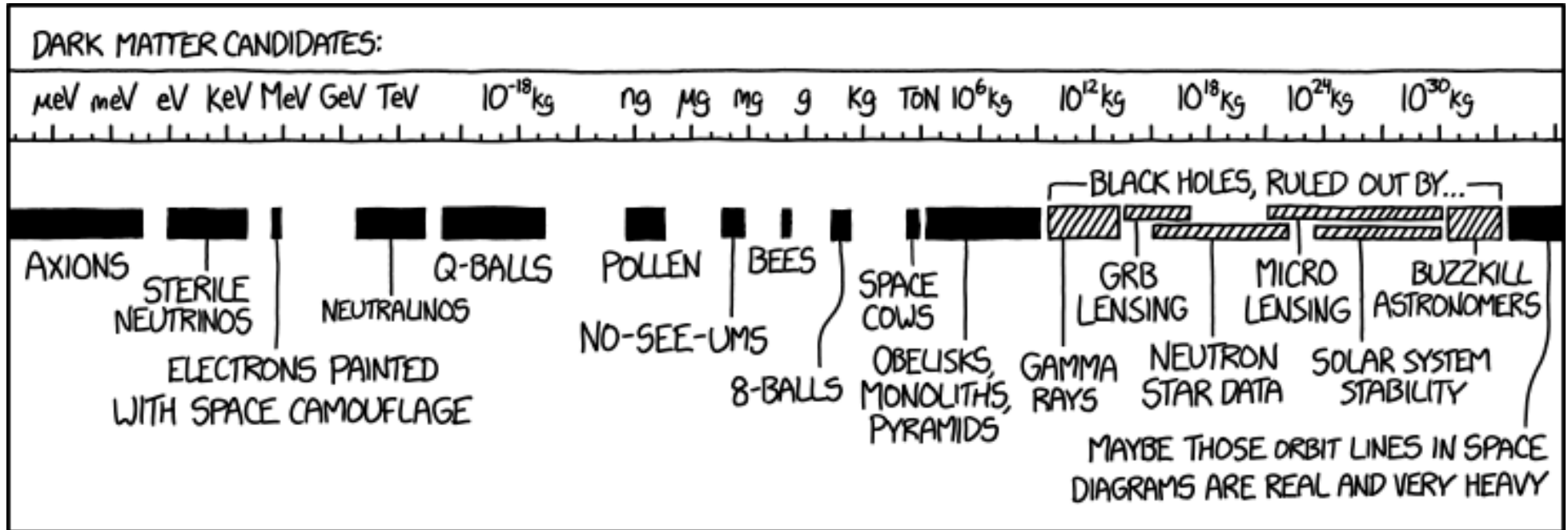
Alejandro Ibarra



ICHEPAP 2023,
Kolkata
December 2023

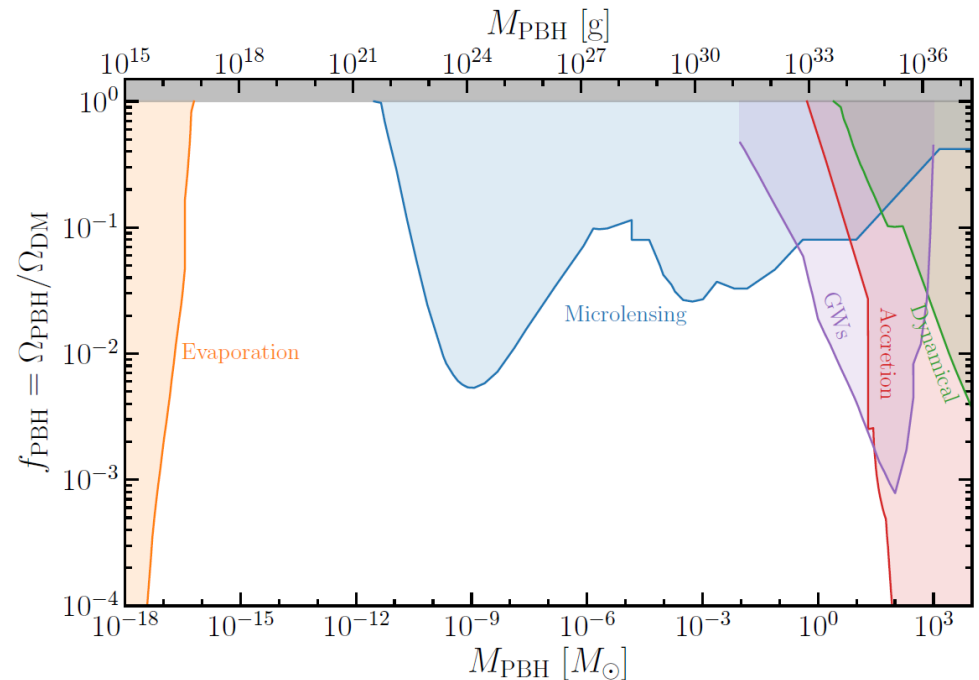
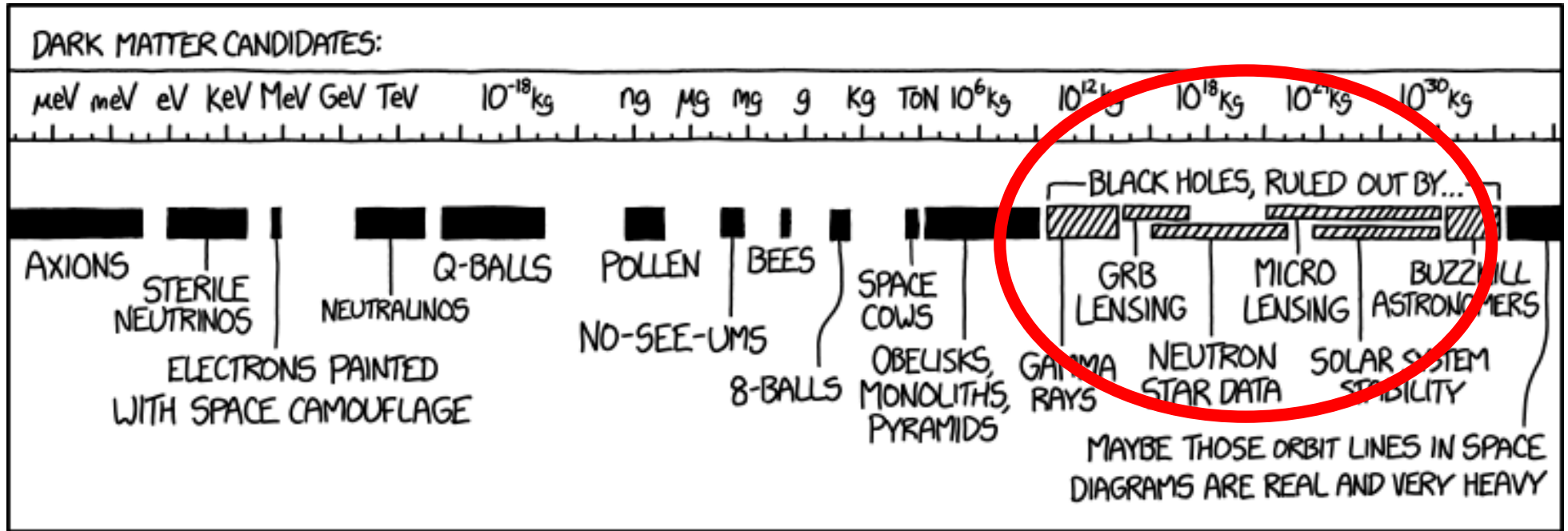
The dark matter zoo

Explain xkcd



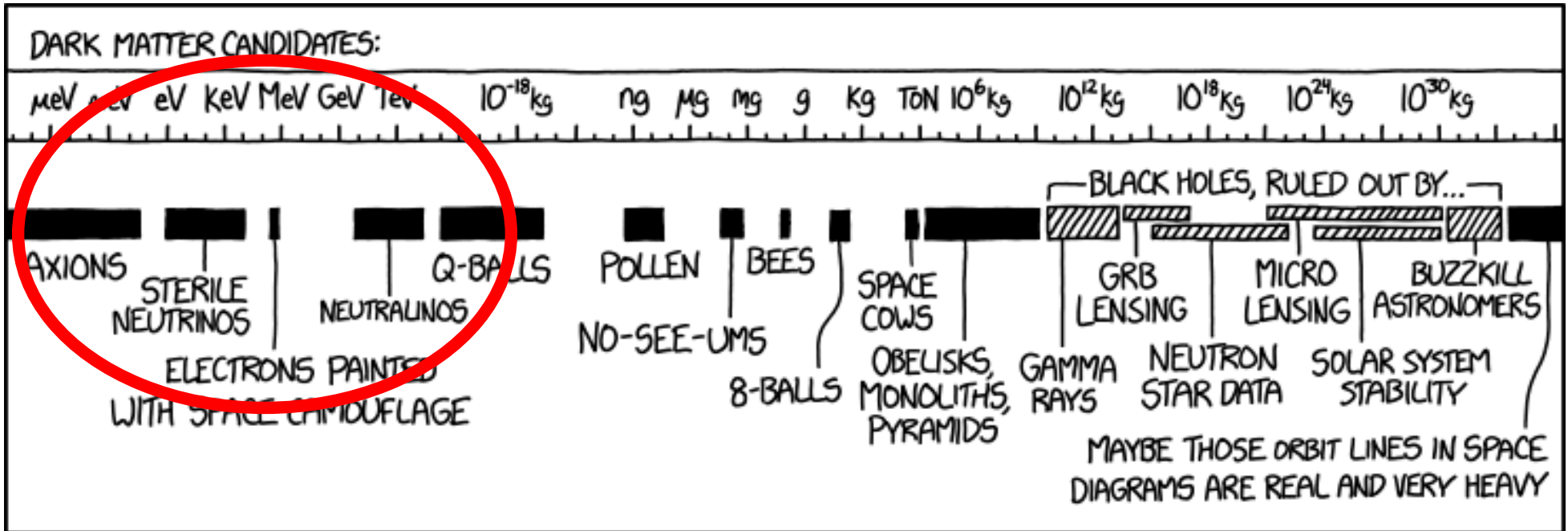
The dark matter zoo

Explain xkcd



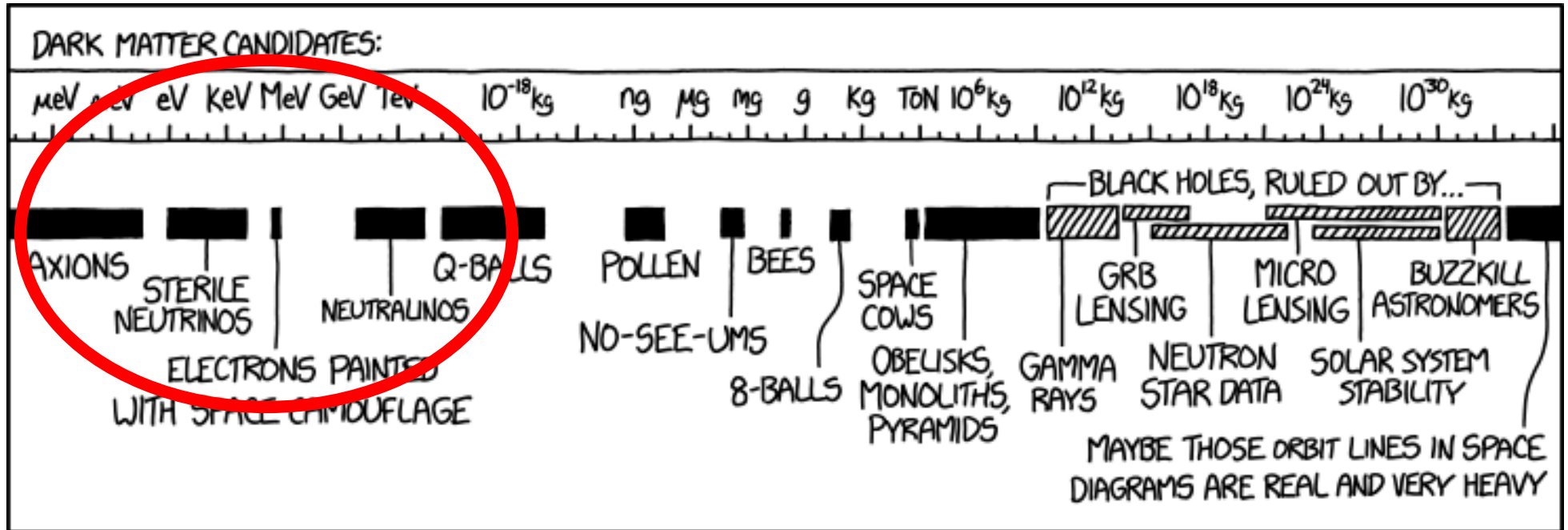
The dark matter zoo

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The dark matter zoo

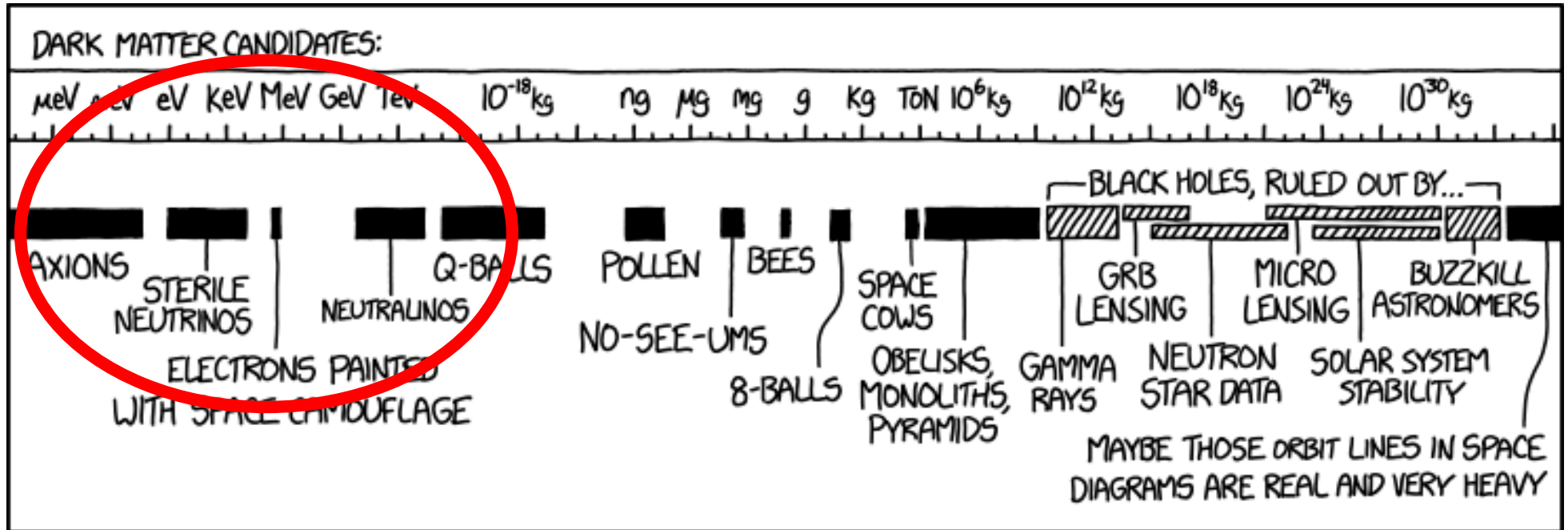
Explain xkcd



→ spin

The dark matter zoo

Explain xkcd

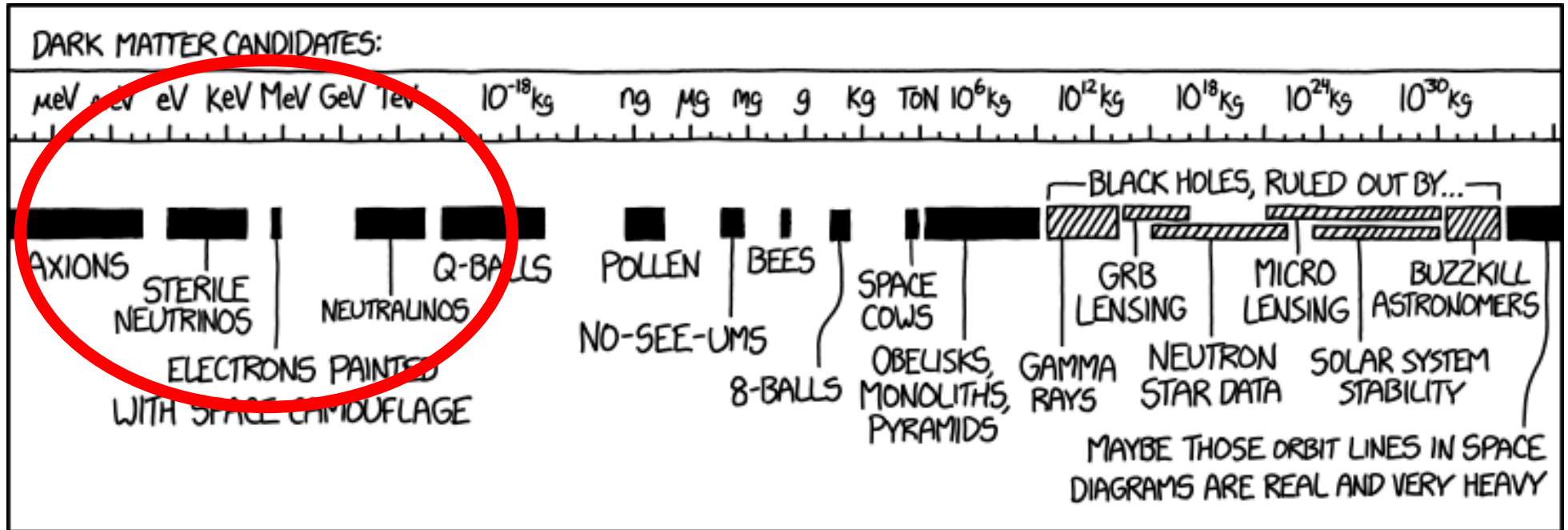


spin

Scattering cross-section off SM particles

The dark matter zoo

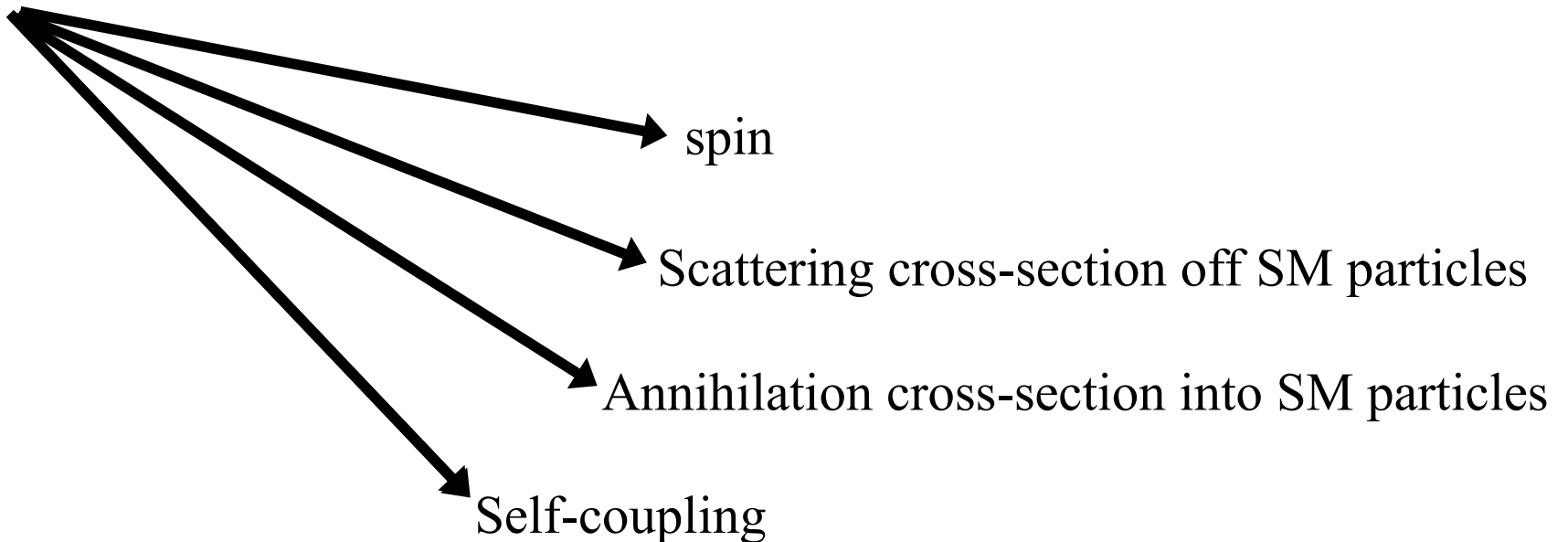
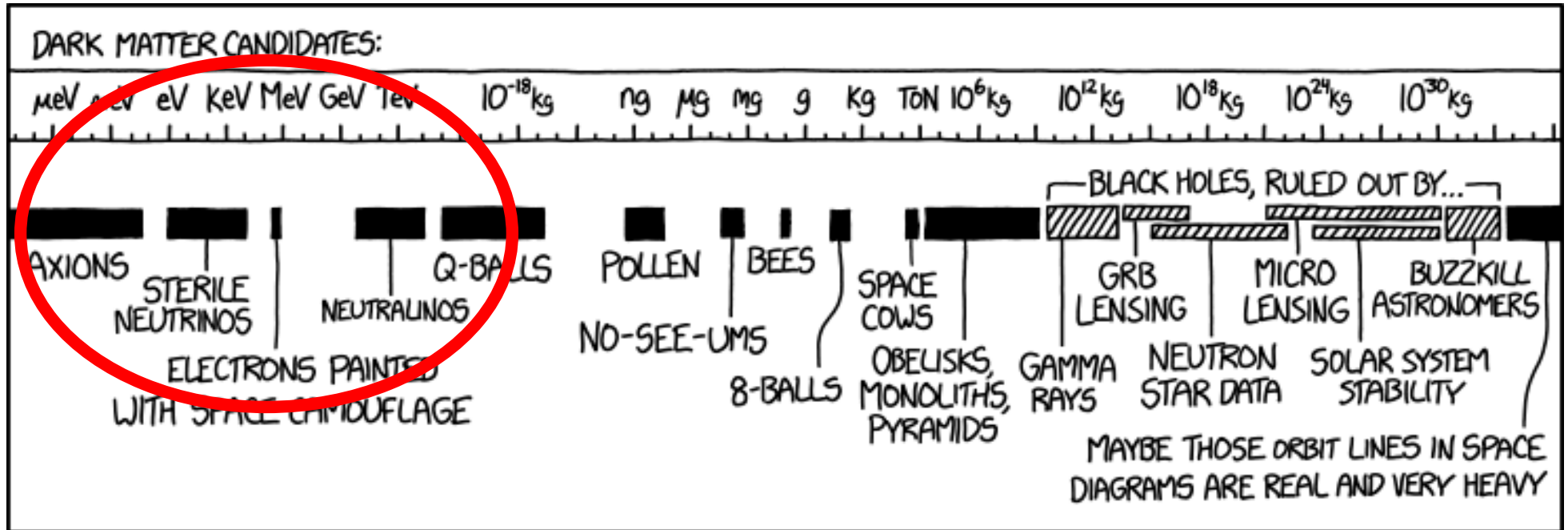
Explain xkcd



- spin
- Scattering cross-section off SM particles
- Annihilation cross-section into SM particles

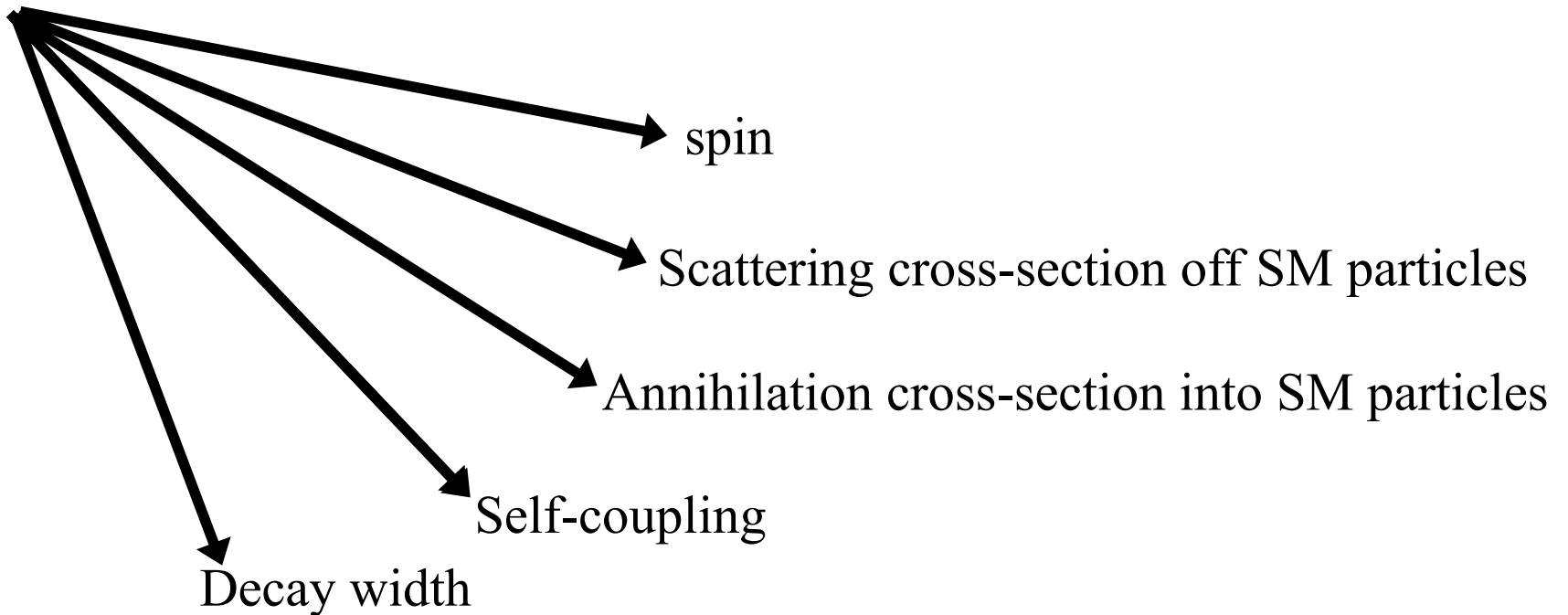
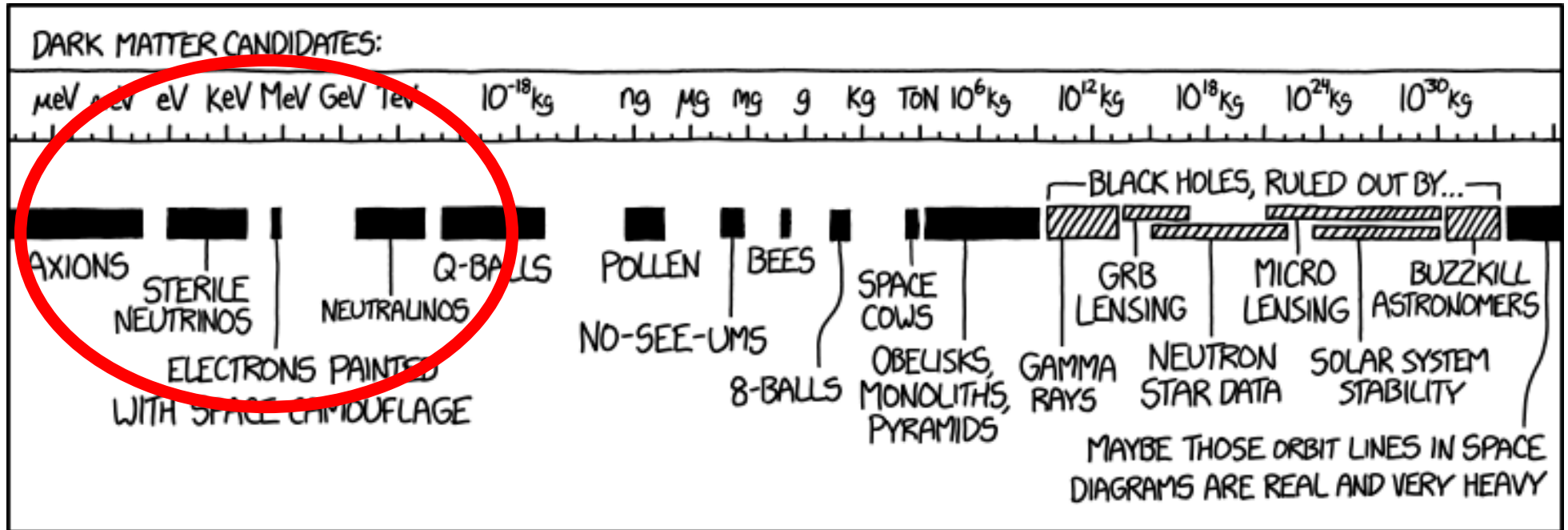
The dark matter zoo

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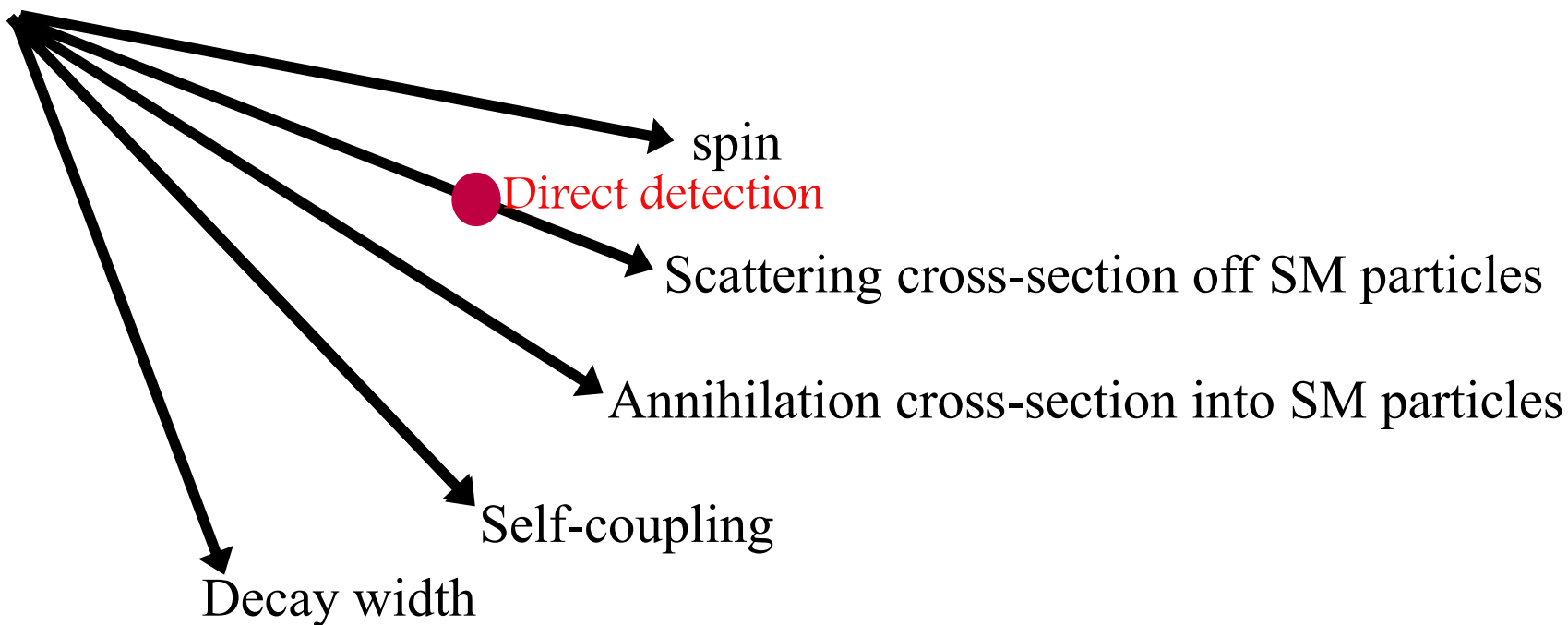
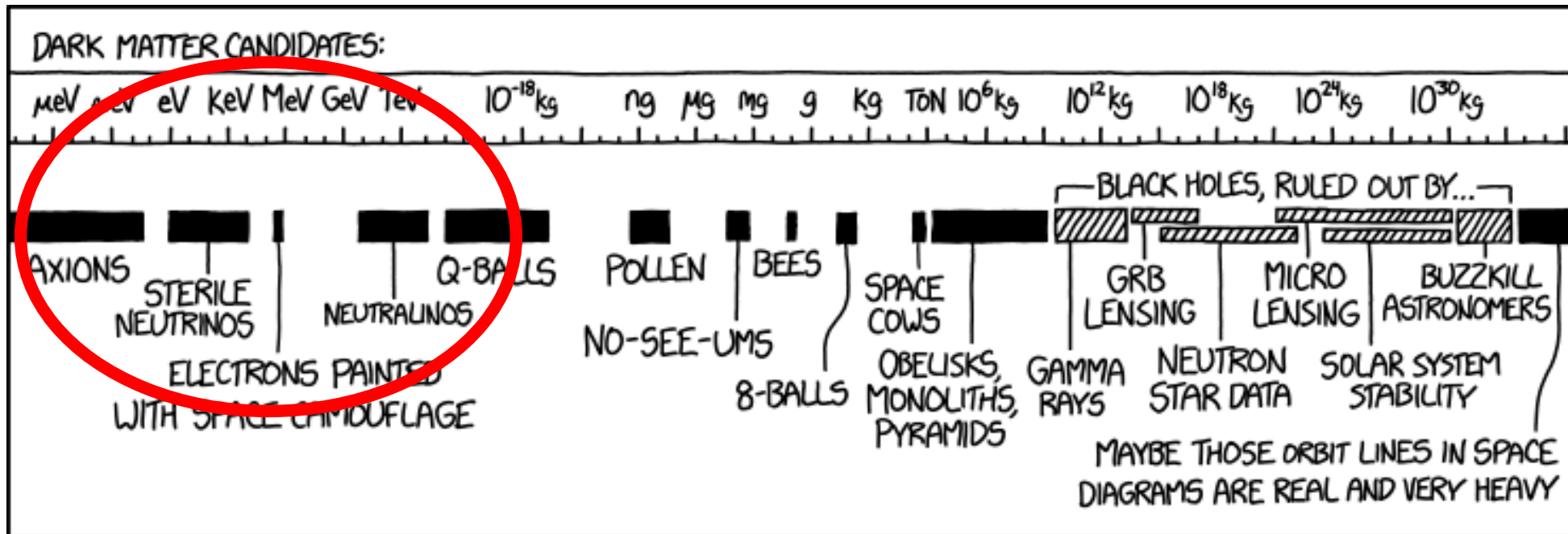
The dark matter zoo

Explain xkcd



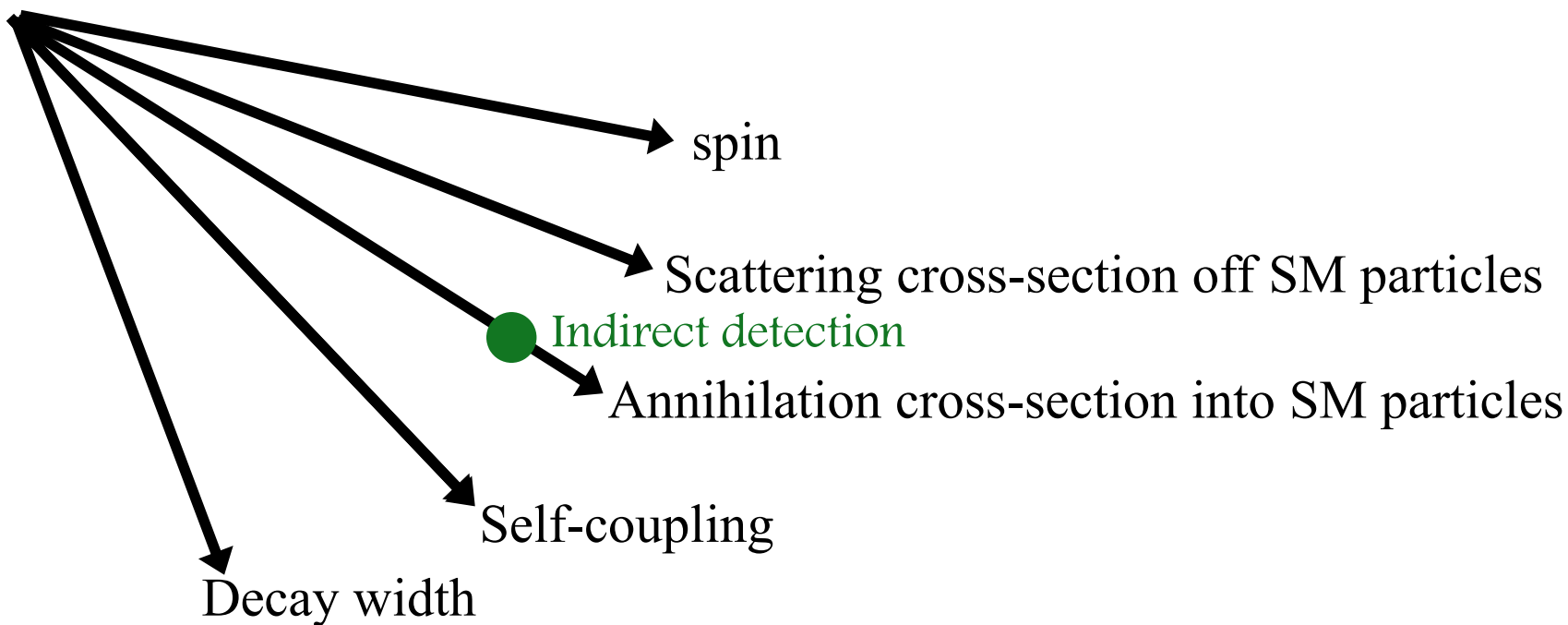
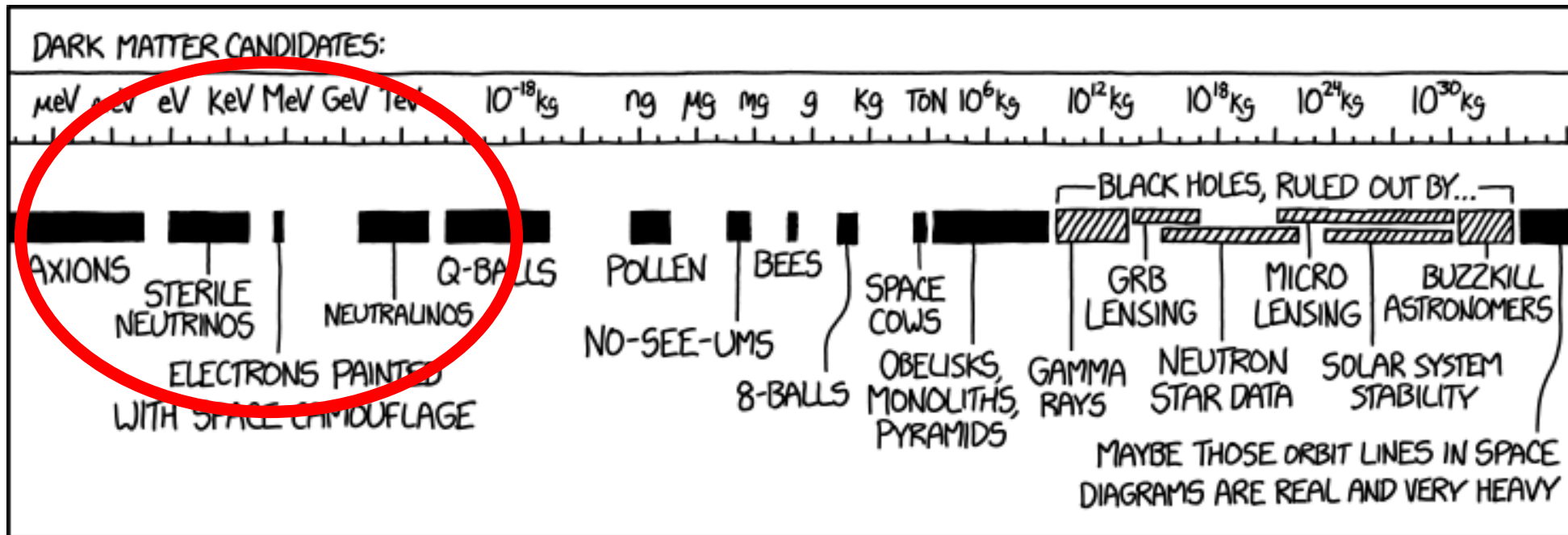
The dark matter zoo

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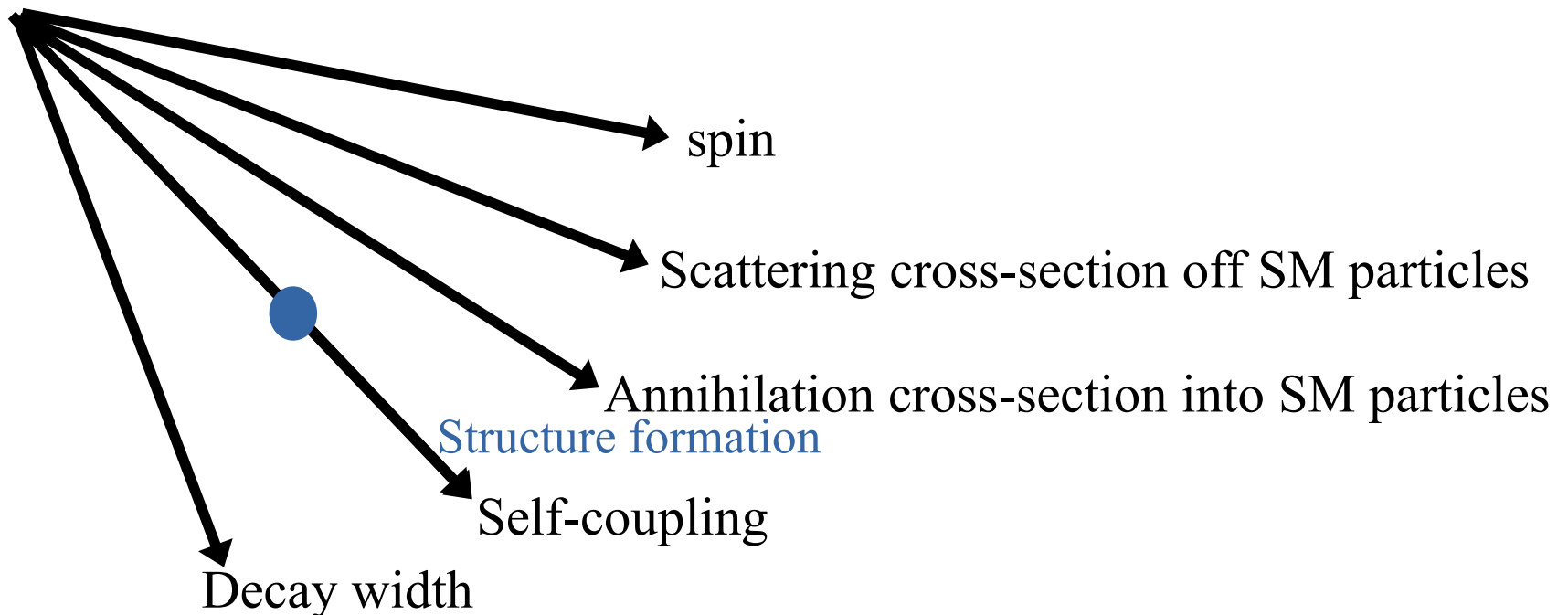
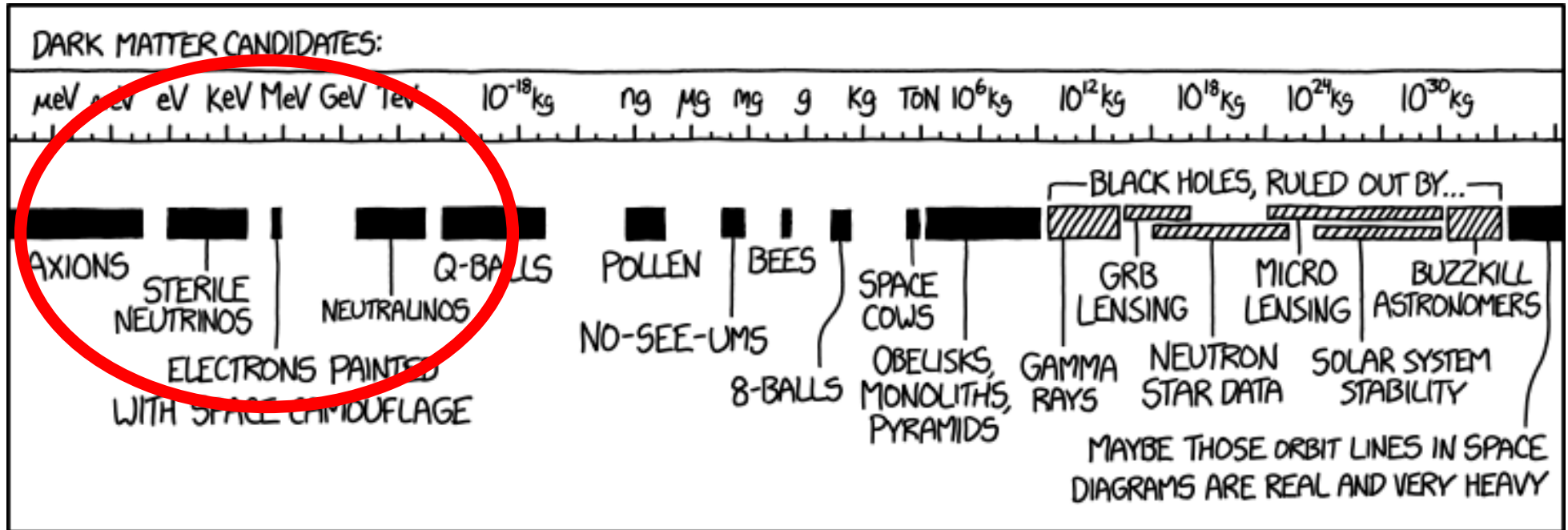
The dark matter zoo

Explain xkcd



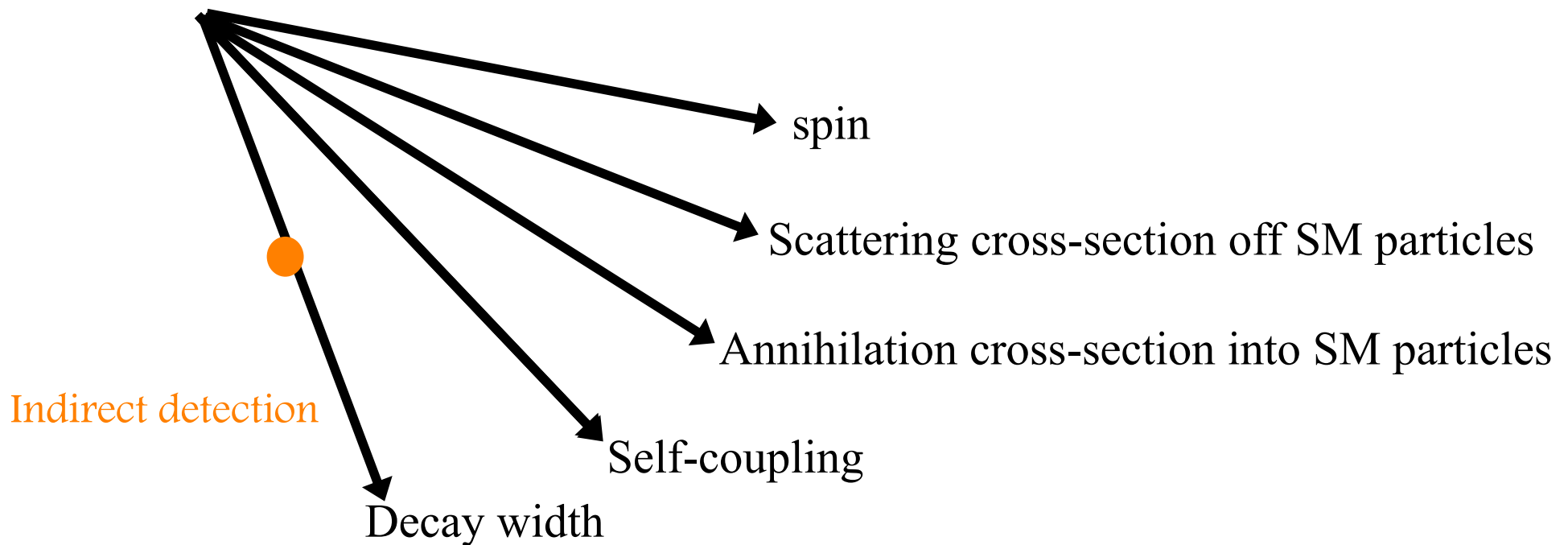
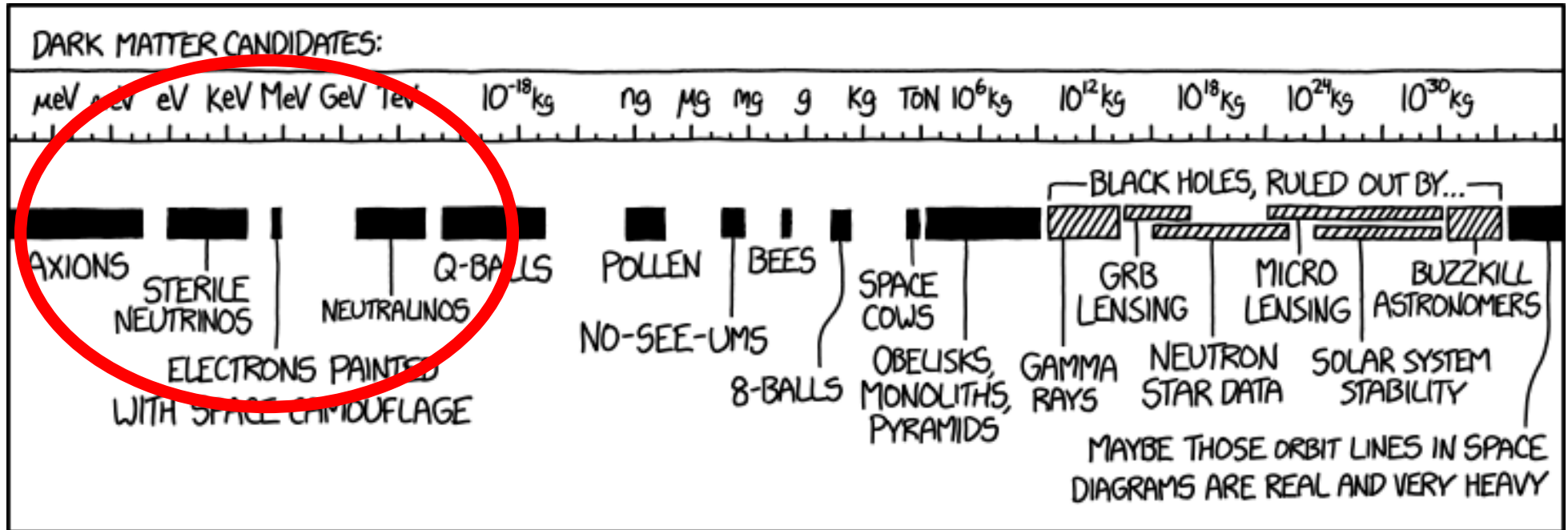
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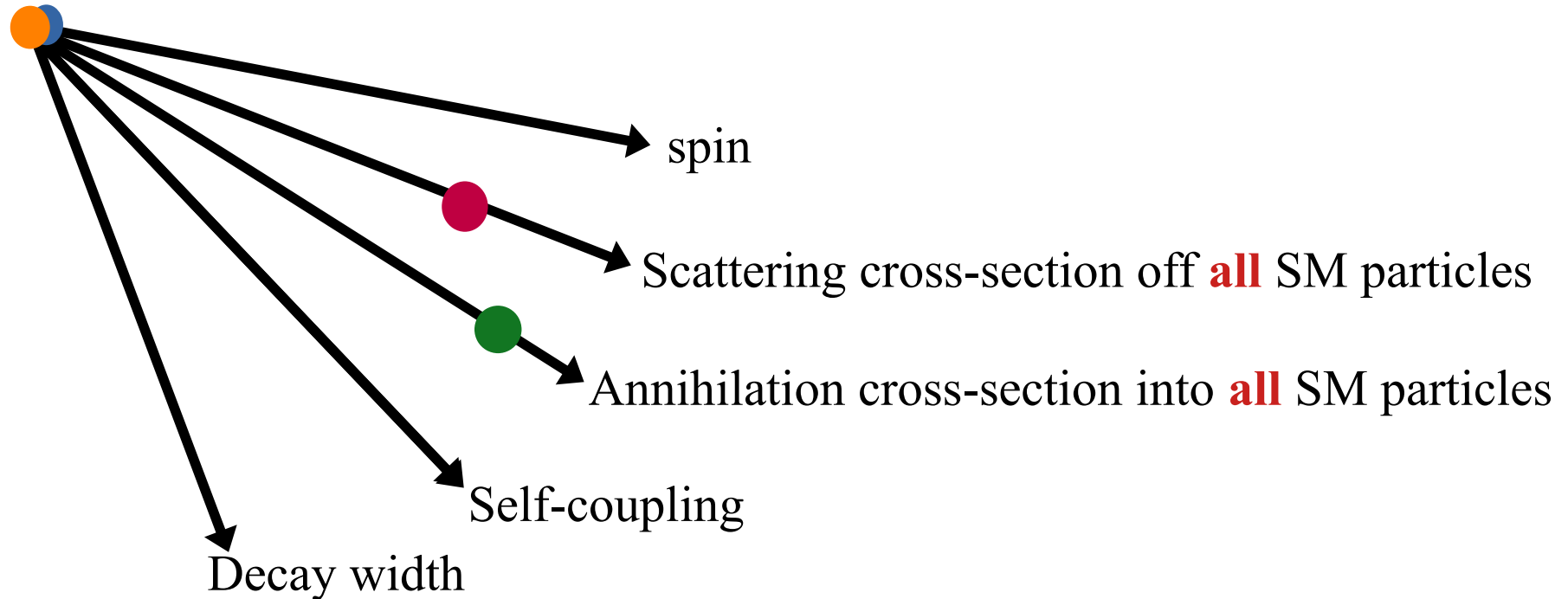
The dark matter zoo

Explain xkcd



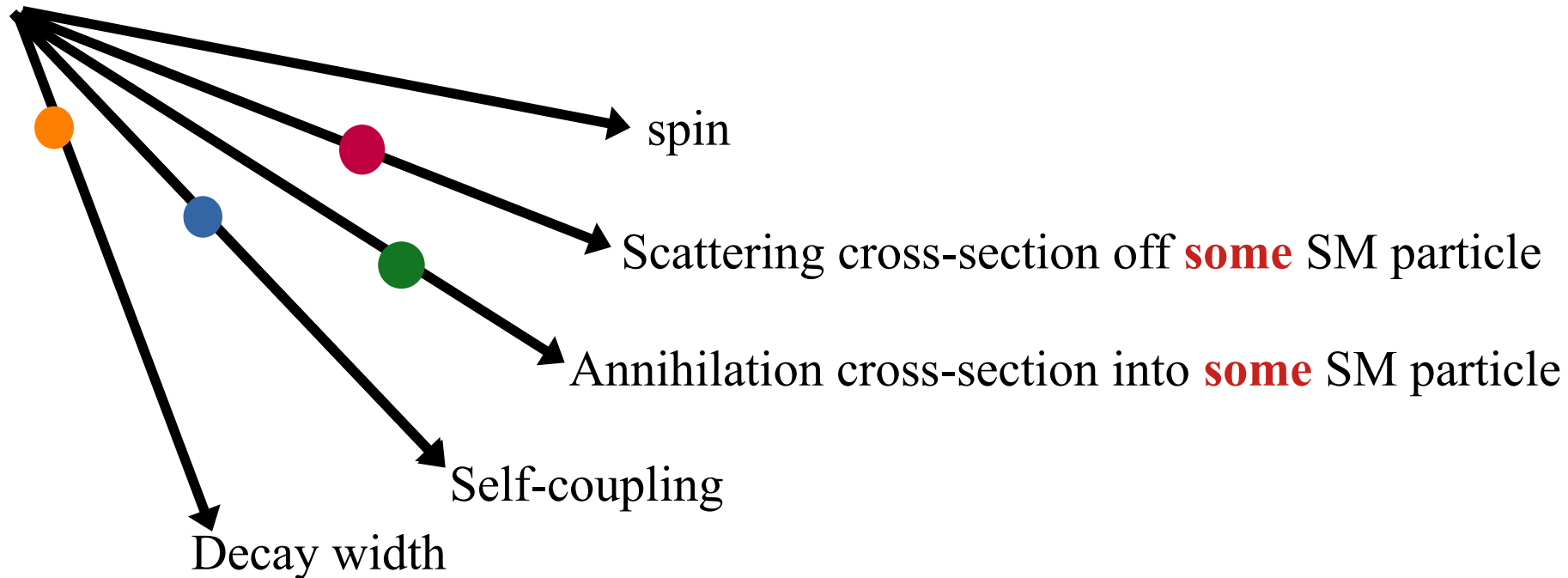
The dark matter zoo

Traditional dark matter searches optimized to detect the lightest neutralino of the Minimal Supersymmetric Standard Model.



The dark matter zoo

Traditional dark matter searches optimized to detect the lightest neutralino of the Minimal Supersymmetric Standard Model.

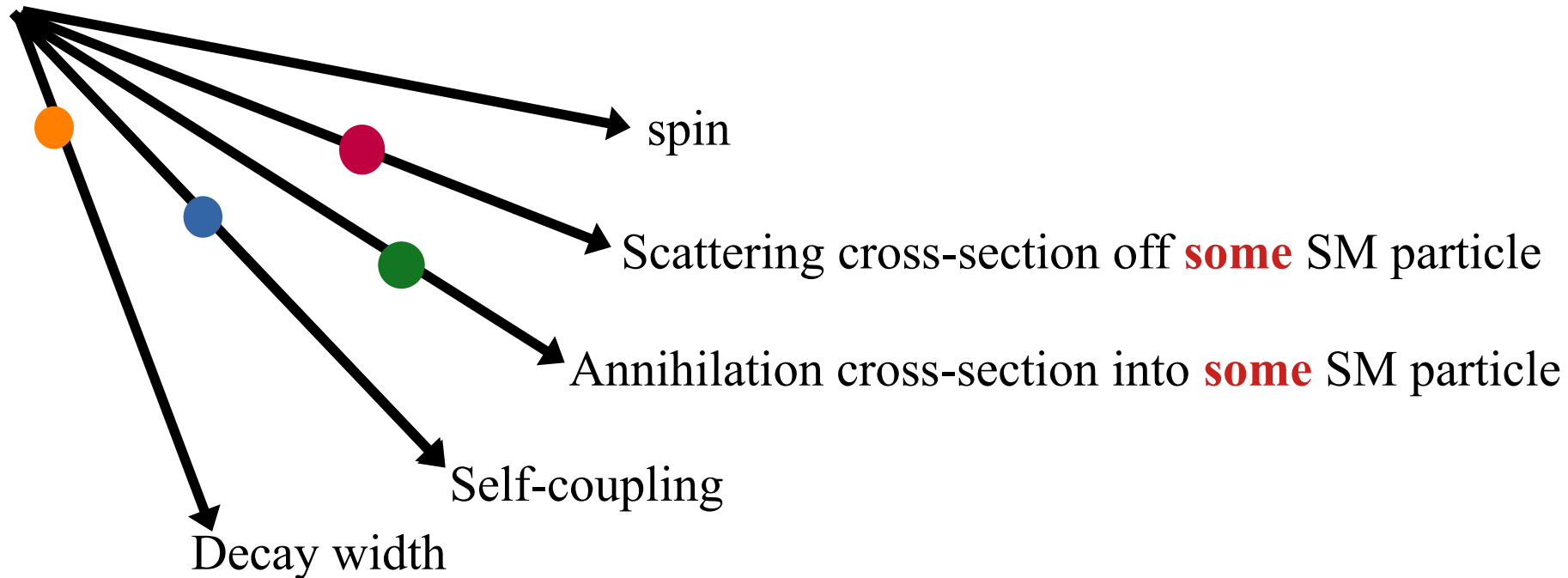


Modern approach:

- Be agnostic about the model.
- Identify distinct DM signals that allow to explore as much parameter space as possible.

The dark matter zoo

Traditional dark matter searches optimized to detect the lightest neutralino of the Minimal Supersymmetric Standard Model.

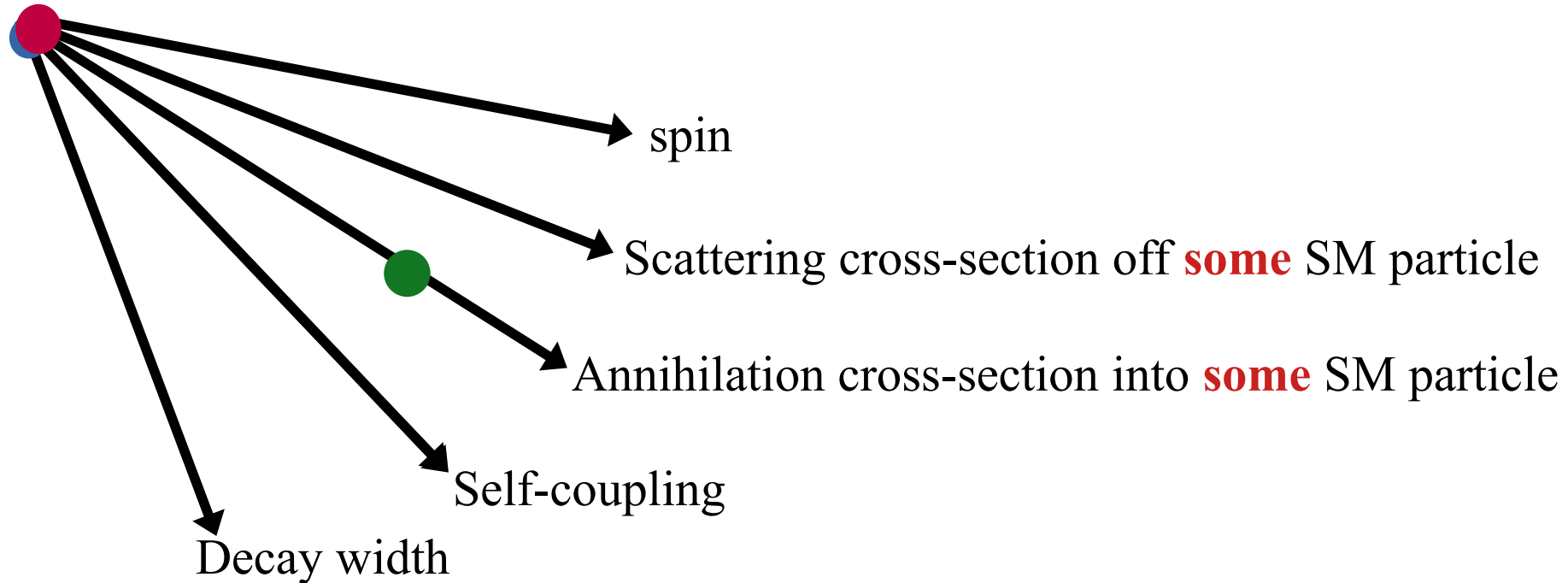


Modern approach:

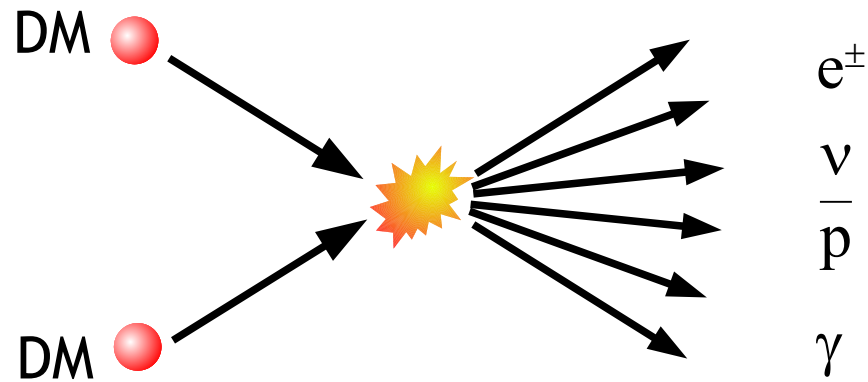
- Be agnostic about the model.
- Identify distinct DM signals that allow to explore as much parameter space as possible.

No stone must be left unturned!

Probing the annihilation cross-section



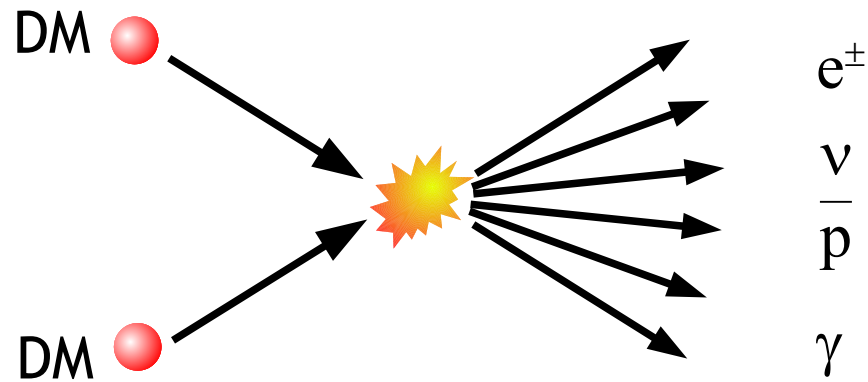
Probing the annihilation cross-section



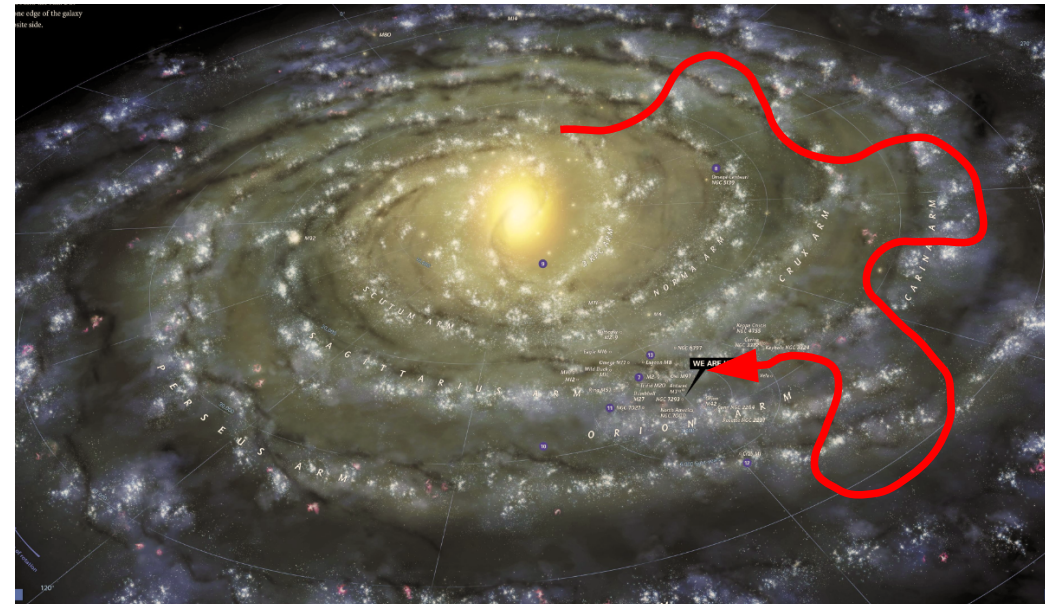
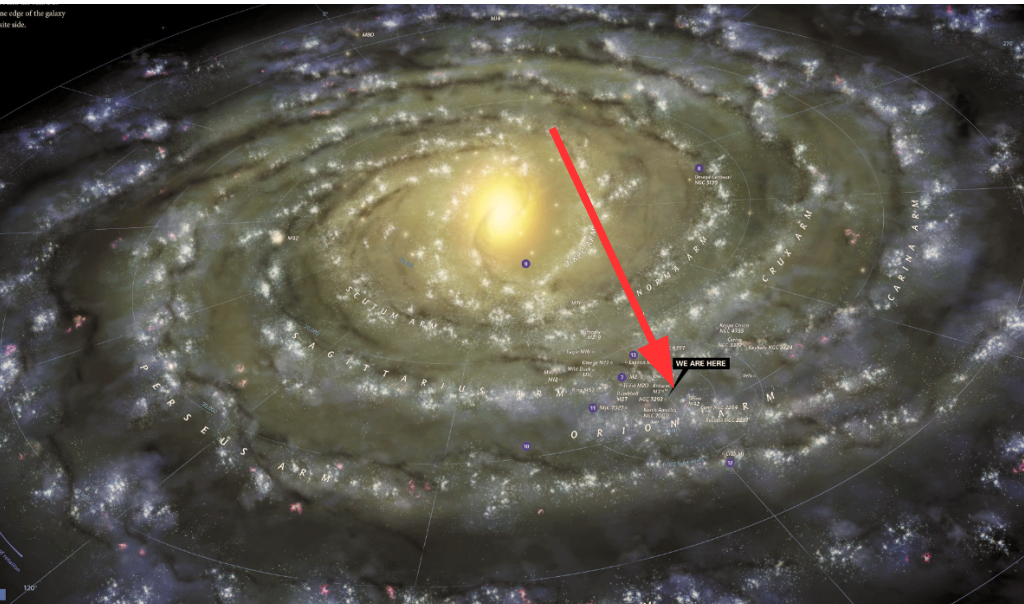
Number of particles of the type “i” produced at the position r per unit time and unit volume:

$$Q(T, \vec{r}) = \frac{1}{2} \frac{\rho_\chi^2(\vec{r})}{m_\chi^2} \sum_i (\sigma v)_i \frac{dN^i}{dT}$$

Probing the annihilation cross-section

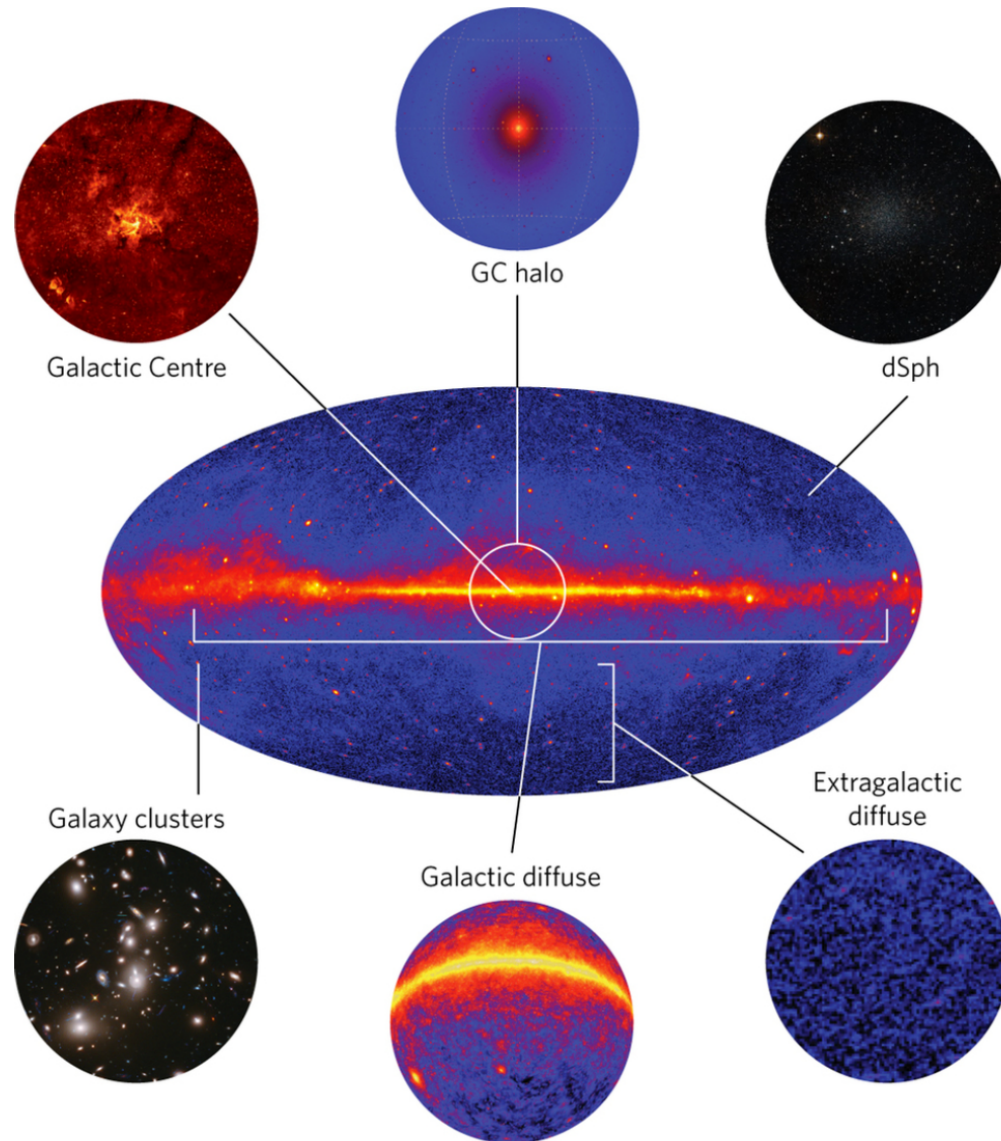


Neutral particles propagate in straight lines practically without losing energy. Charged particles, on the other hand, propagate in a complicated way through the tangled magnetic field of our Galaxy.



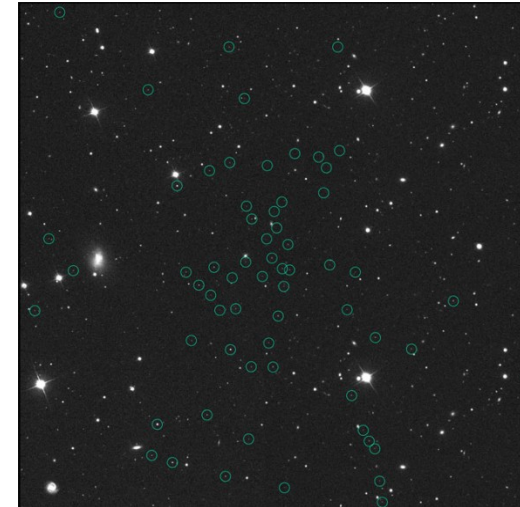
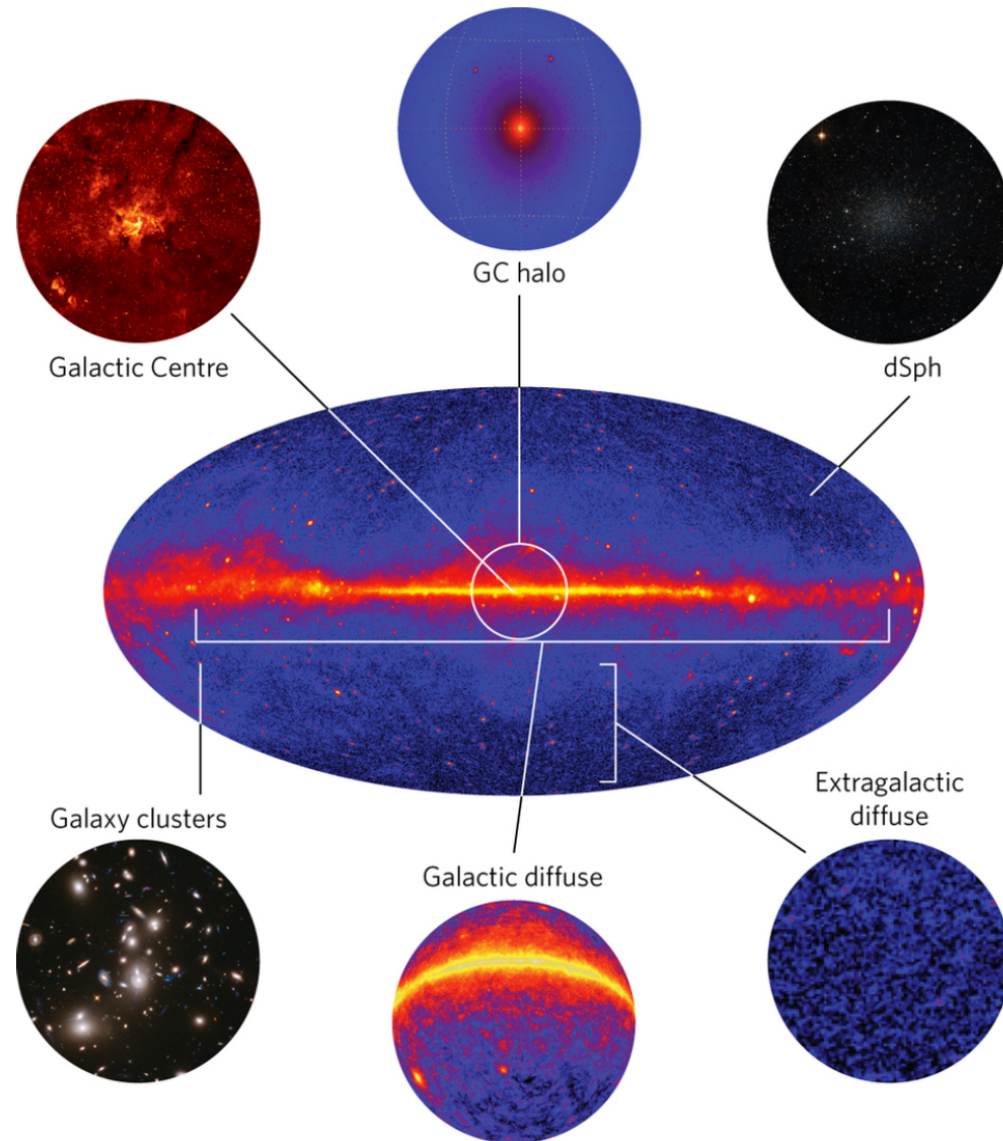
Gamma-rays from dark matter annihilation

Possible targets for detection of gamma-rays from annihilation



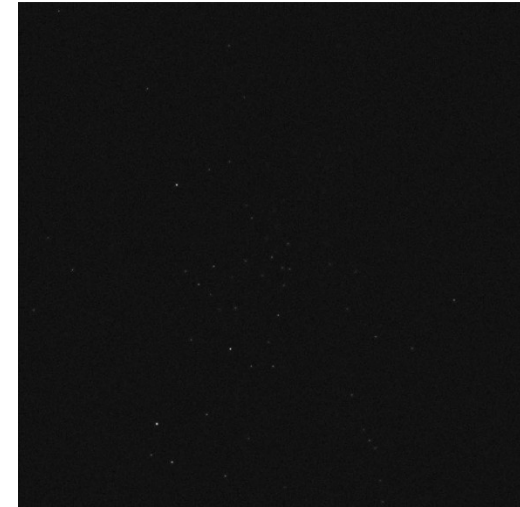
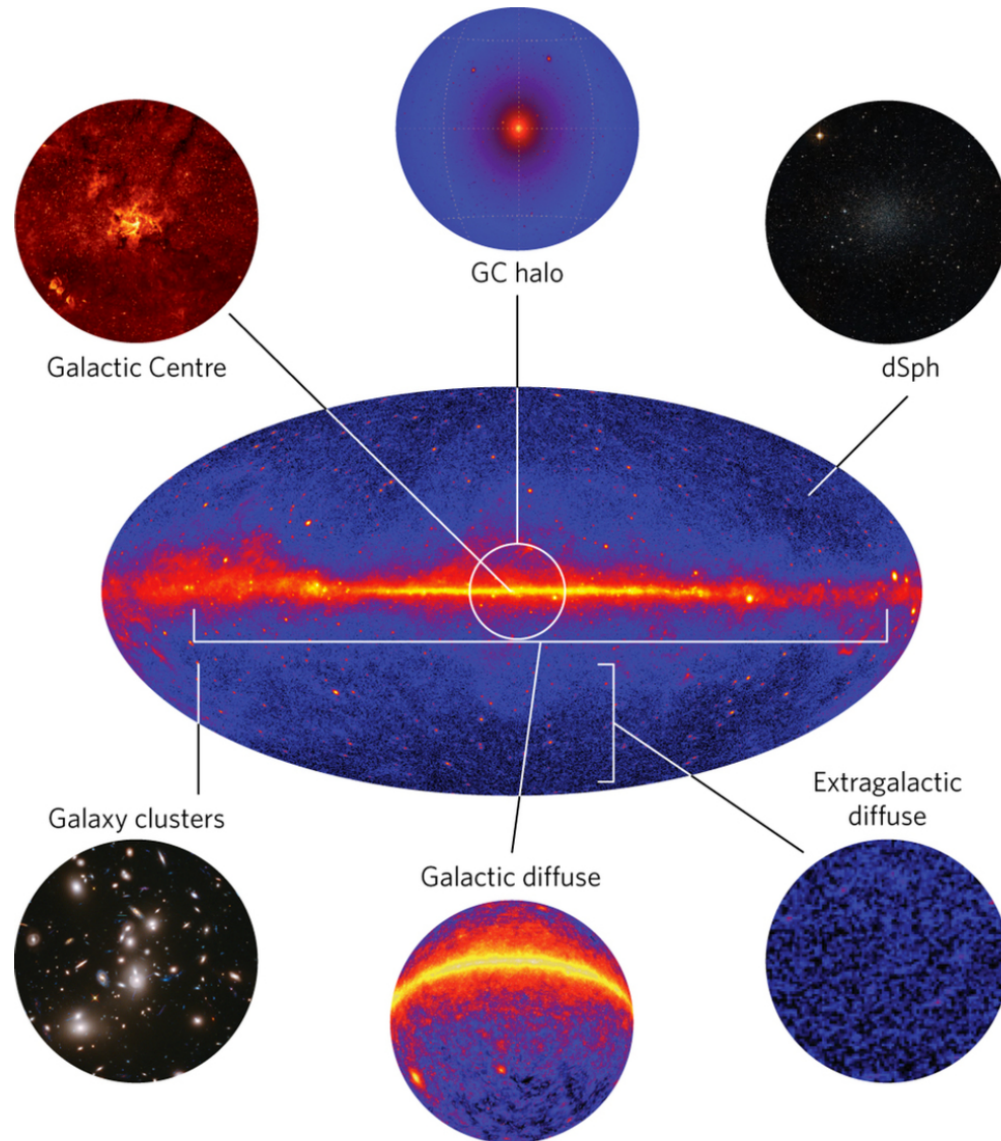
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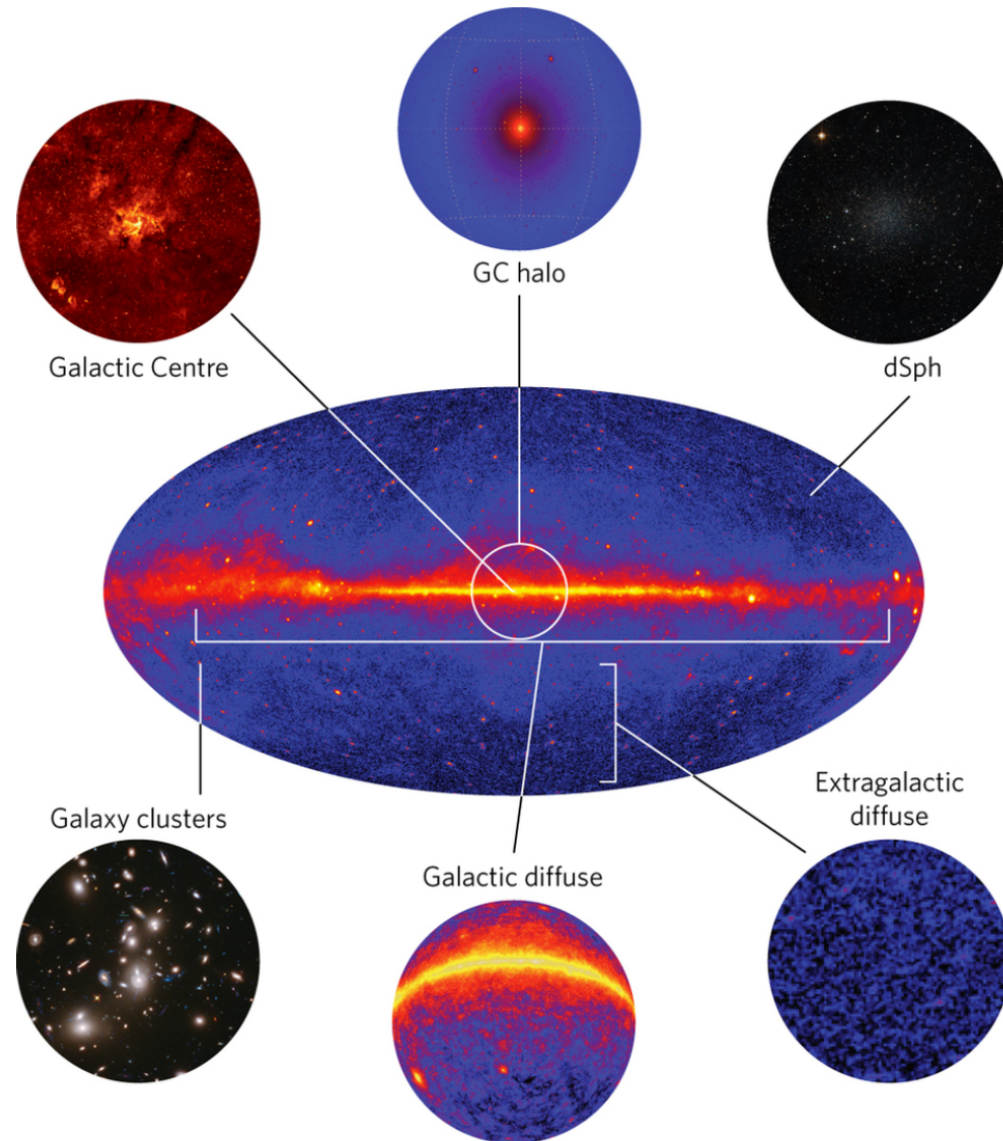
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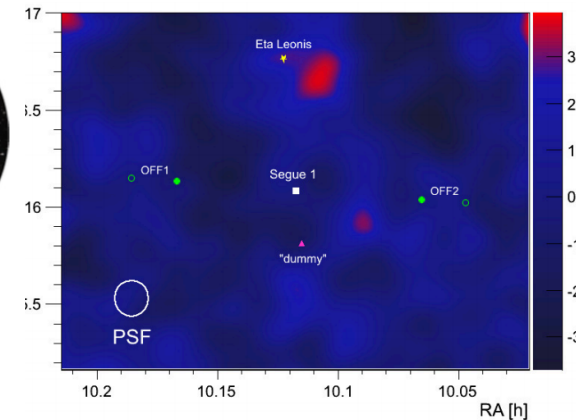
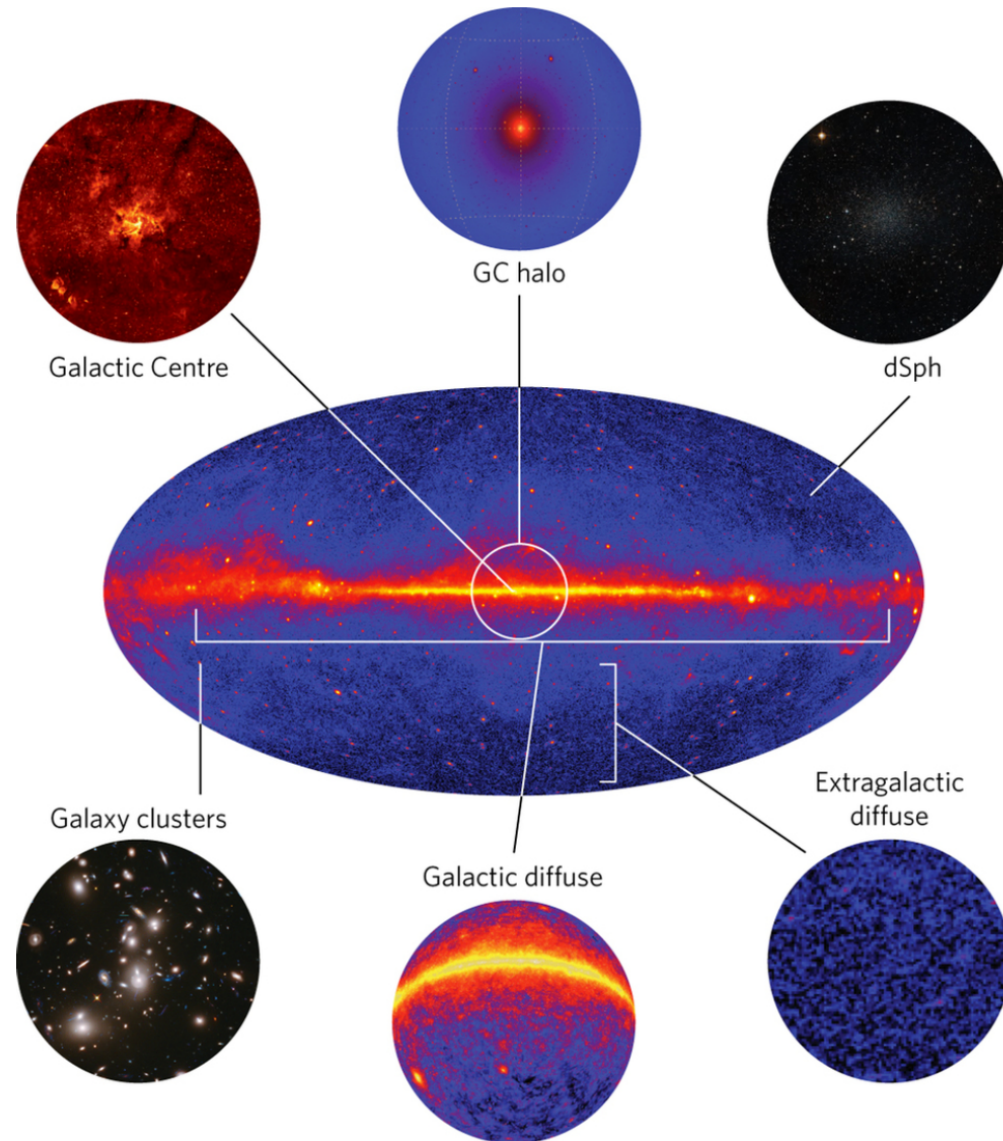
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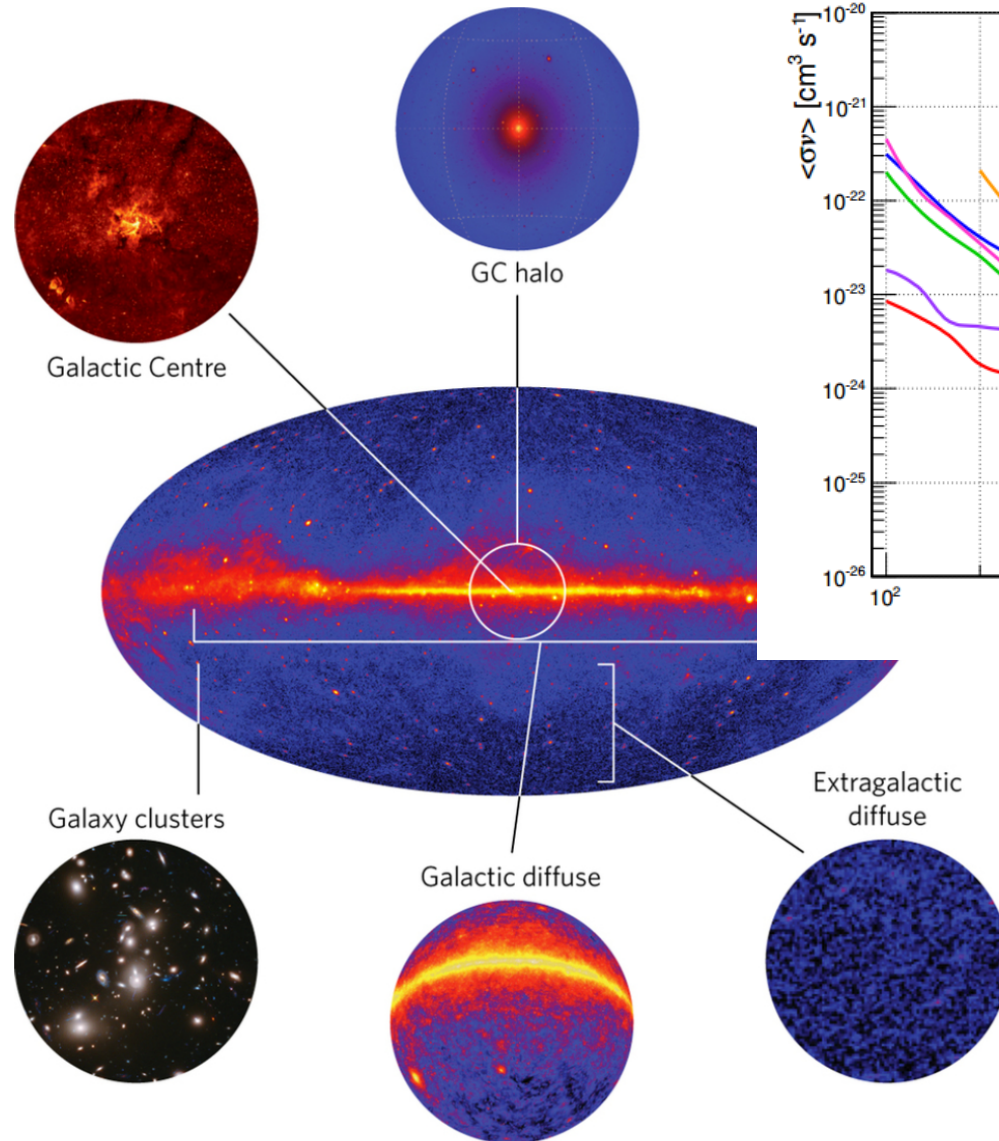
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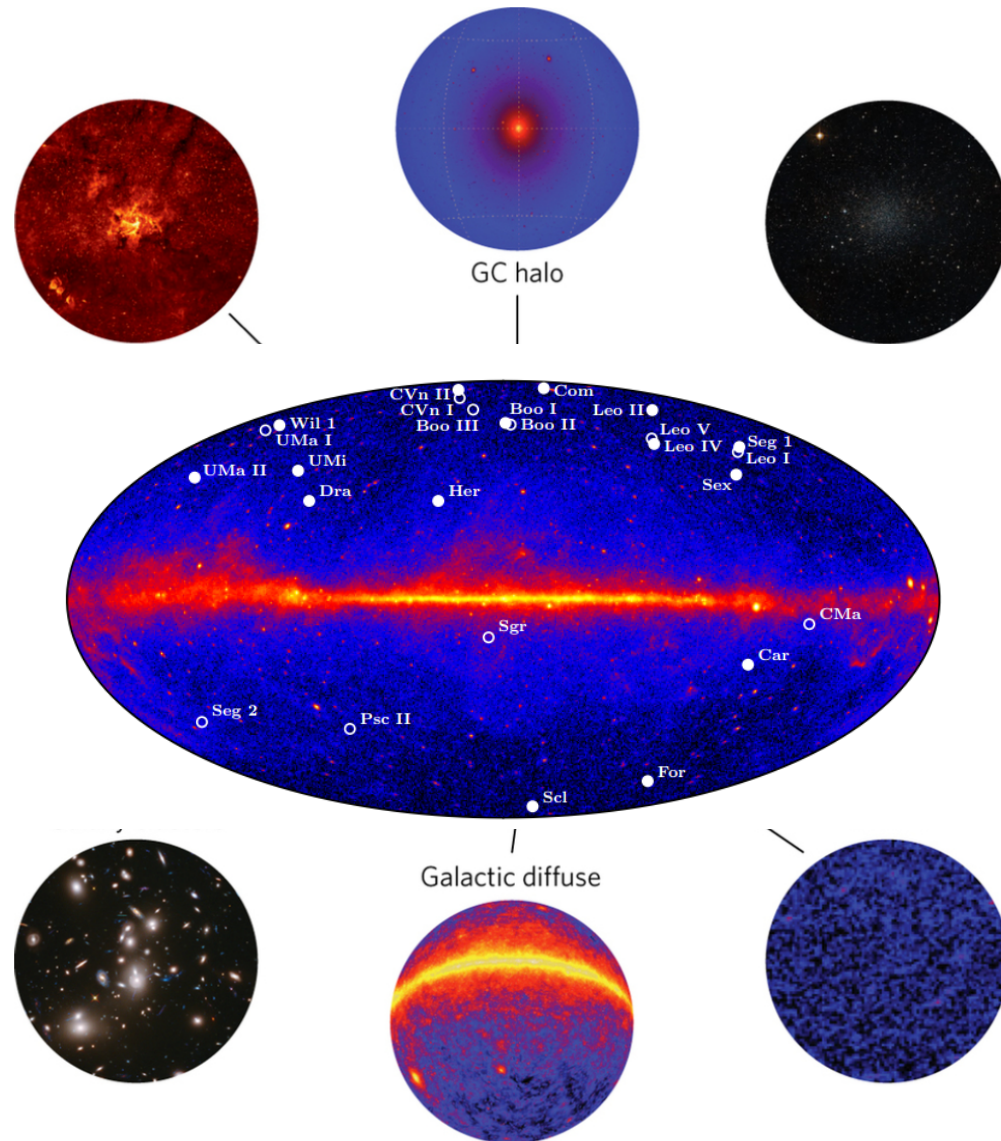
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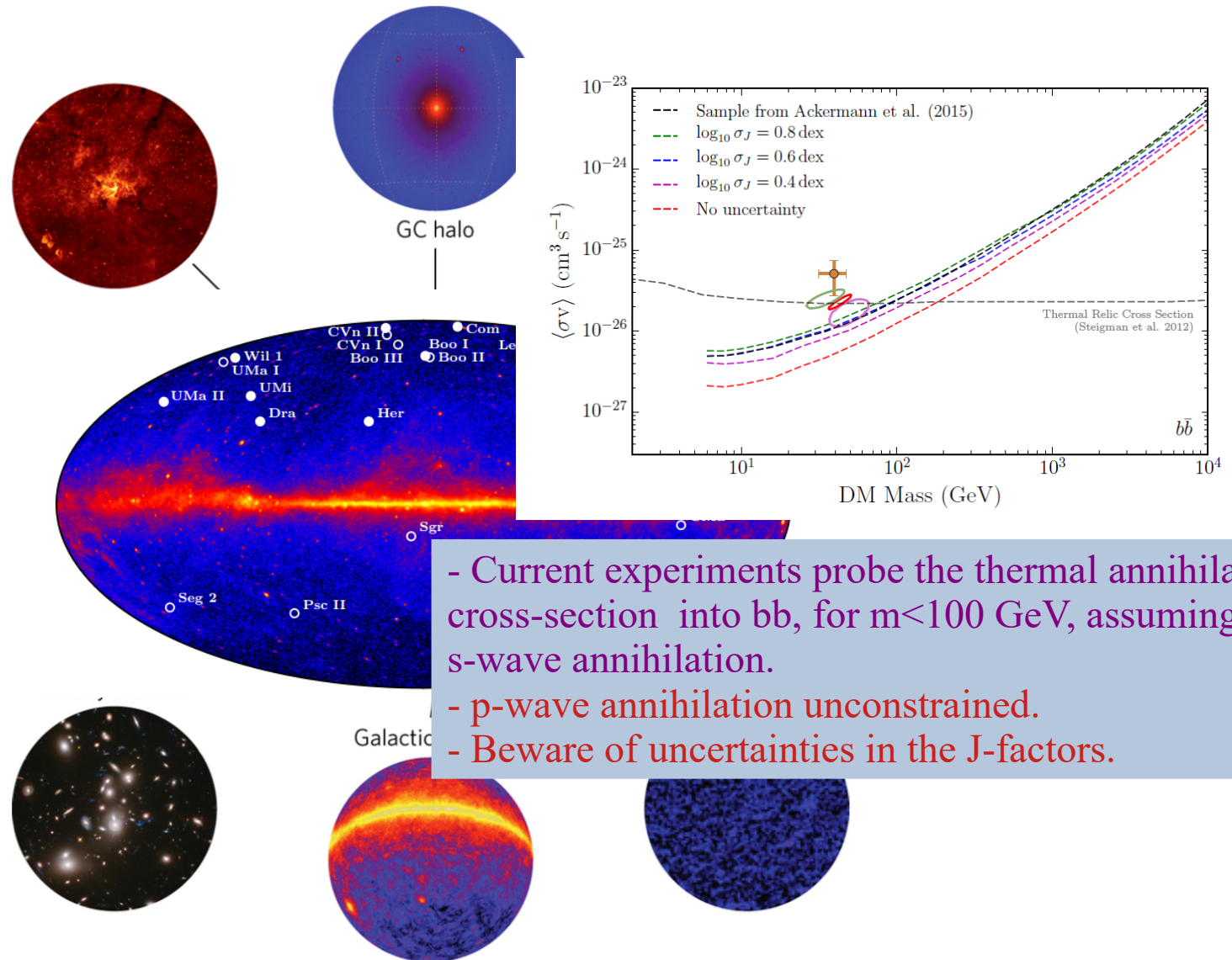
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Gamma-rays from dark matter annihilation

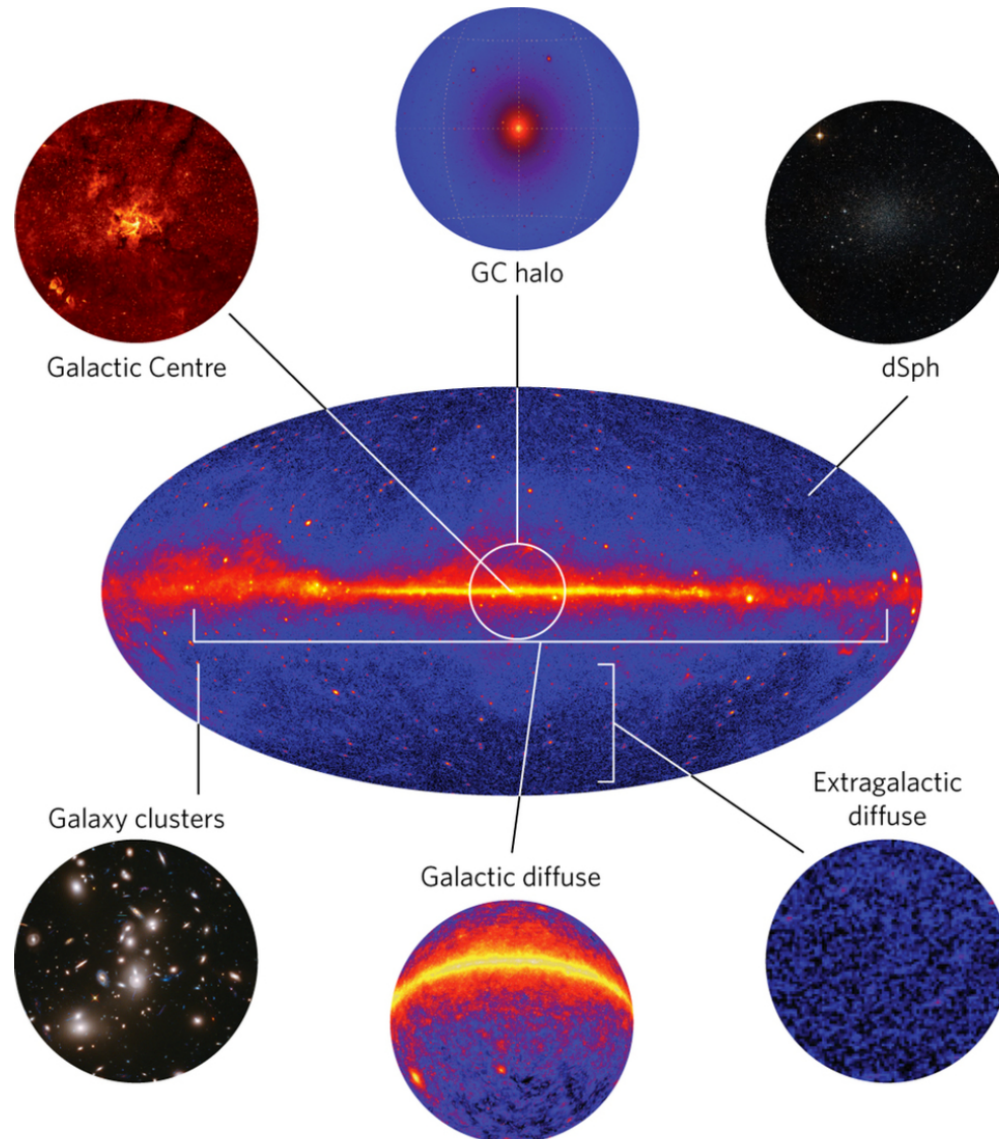
Possible targets for detection of gamma-rays from annihilation



- Current experiments probe the thermal annihilation cross-section into bb , for $m < 100$ GeV, assuming s-wave annihilation.
- p-wave annihilation unconstrained.
- Beware of uncertainties in the J-factors.

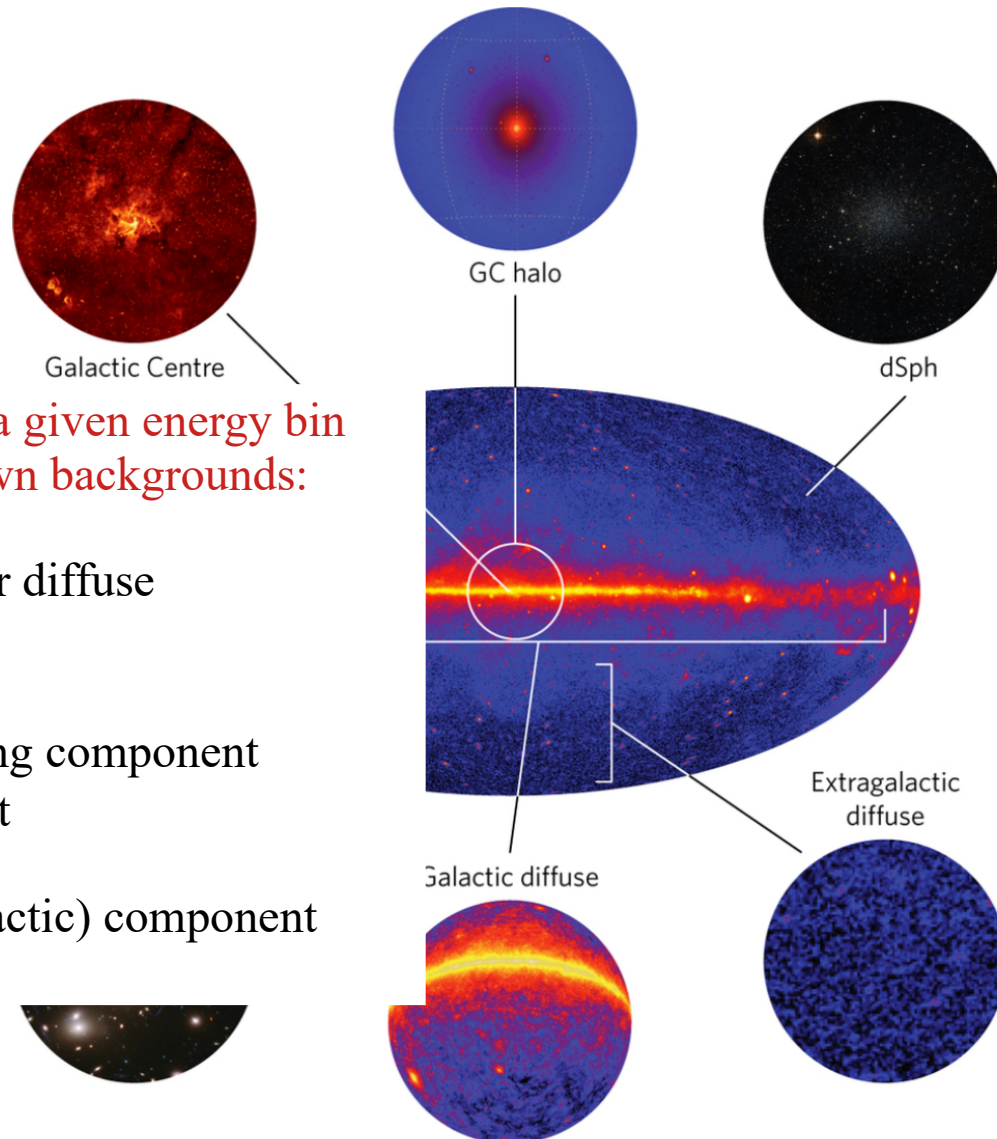
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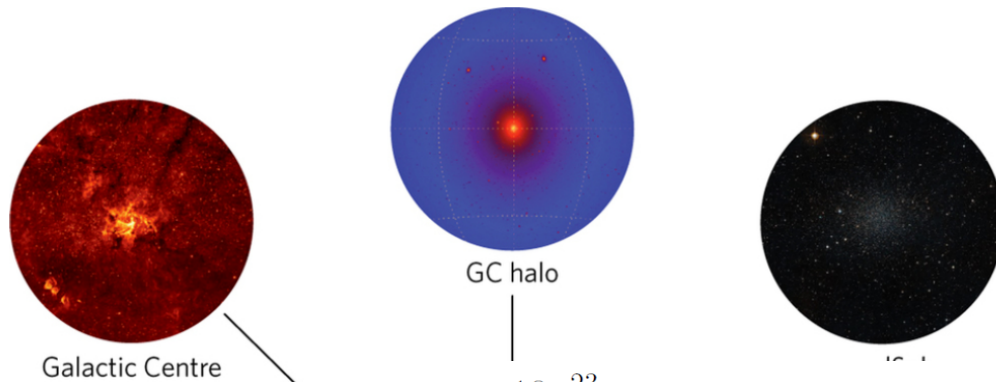


Take the sky map in a given energy bin and subtract the known backgrounds:

- Sources
- Spatial template for diffuse galactic emission
 - π^0 component
 - Bremsstrahlung component
 - ICS component
- Fermi bubbles
- Isotropic (extragalactic) component

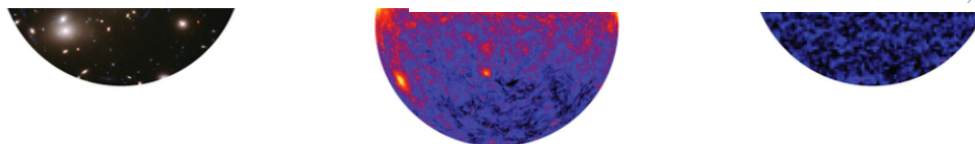
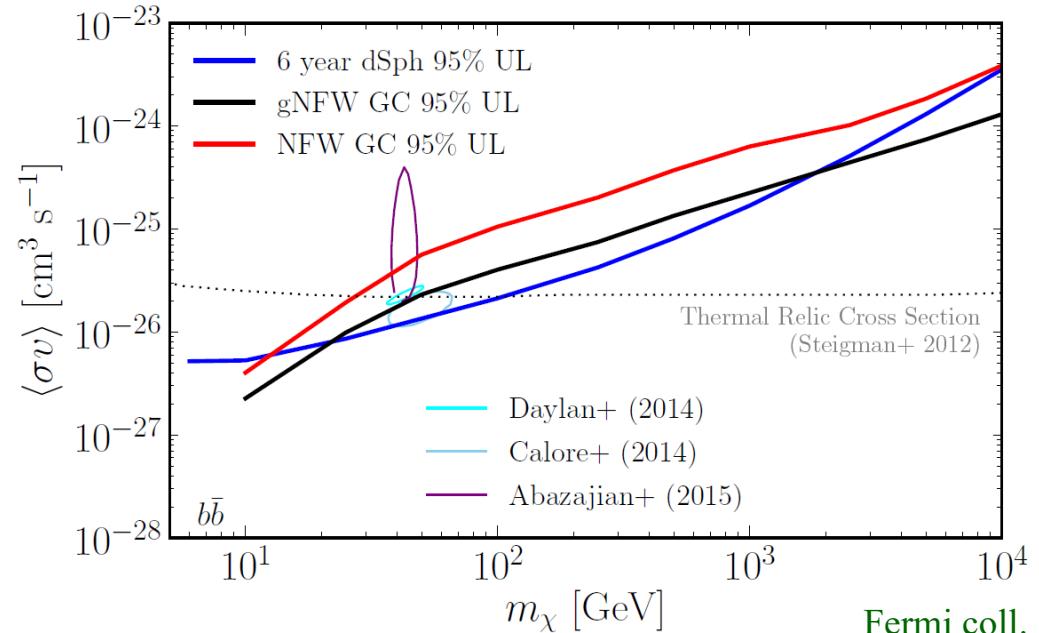
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Fermi coll. '17

Antimatter from dark matter annihilation

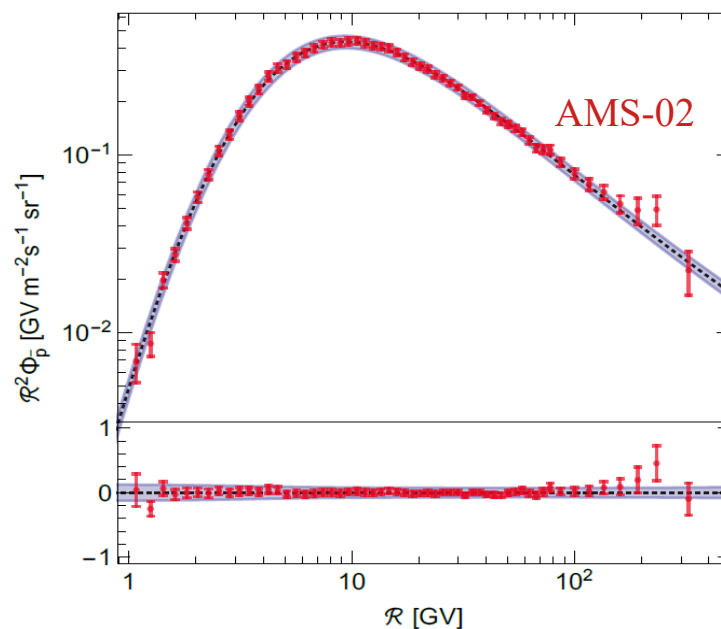
Antimatter particles propagate through the tangled magnetic field of the galaxy in a complicated way, losing energy on their way.

Model the propagation with a diffusion equation:

$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

Assumptions on the quantities entering are necessary.

Still, very good agreement between the expected antiproton flux from collisions of cosmic rays on the nuclei of the interstellar medium, and the antiproton data.



Reinert, Winkler.
1712.00002

Antimatter from dark matter annihilation

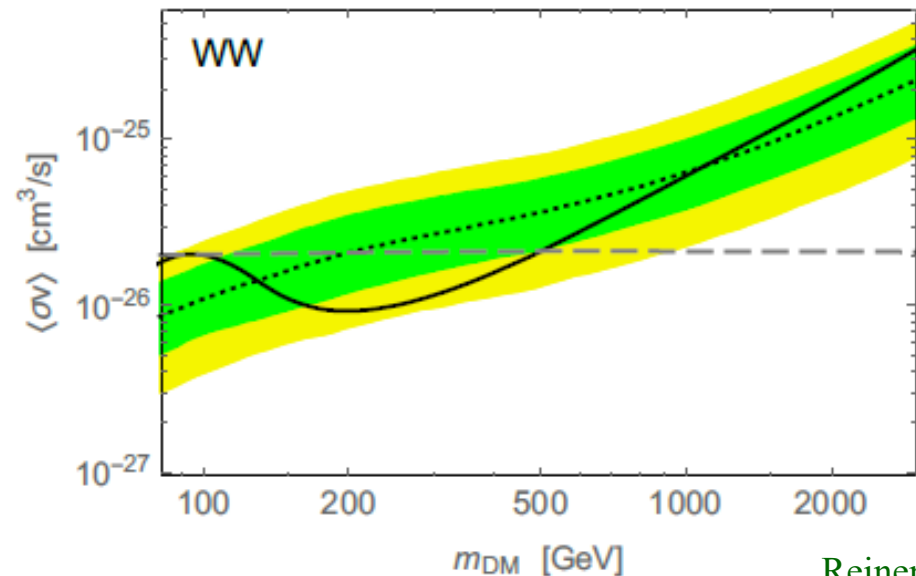
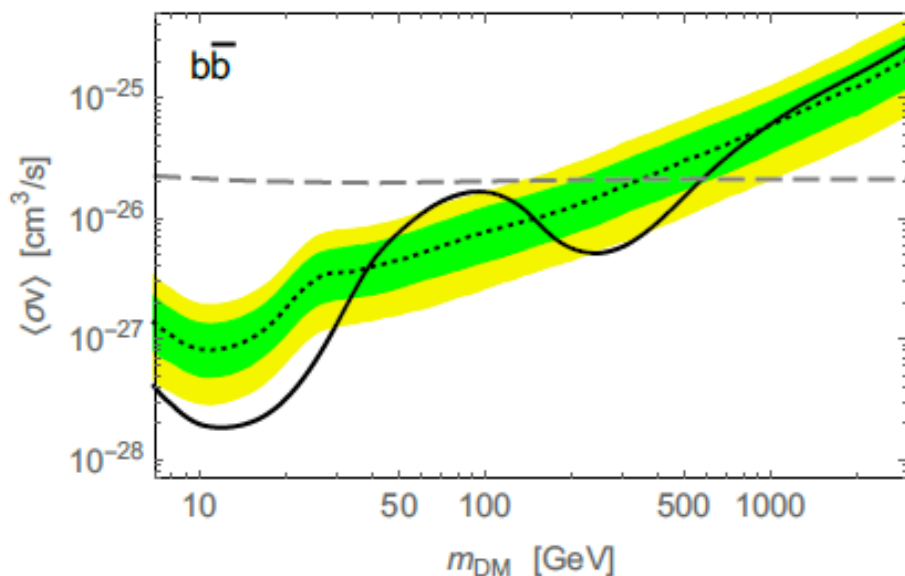
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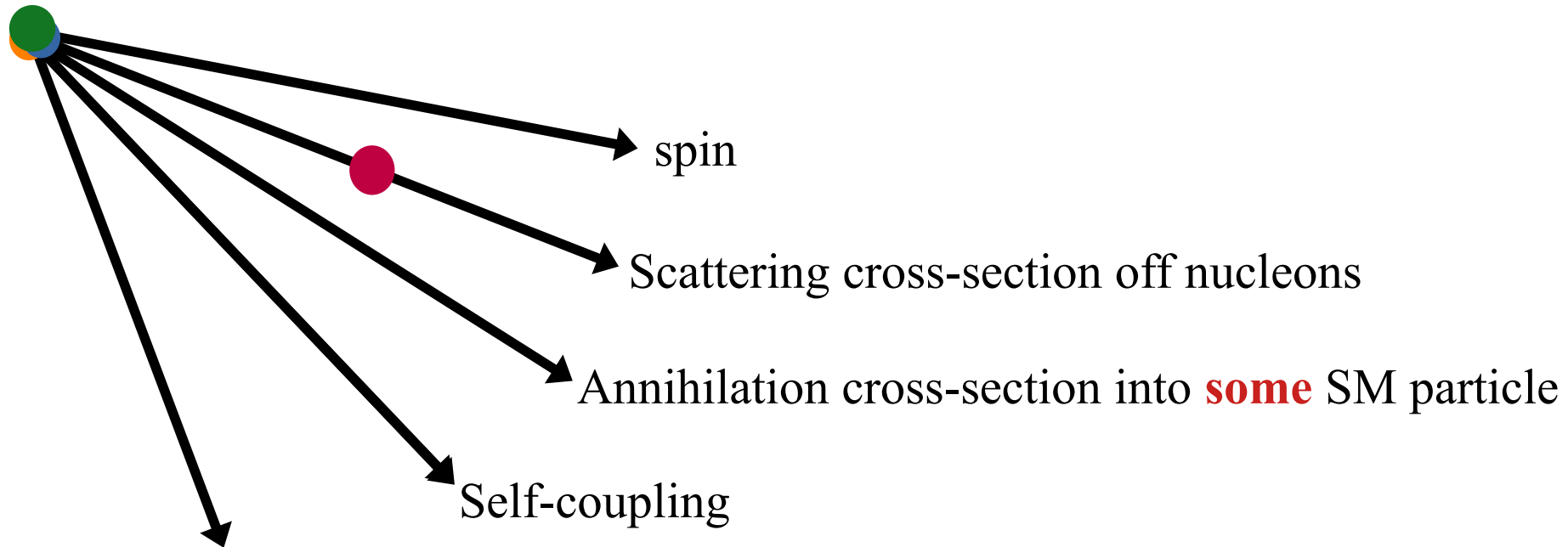
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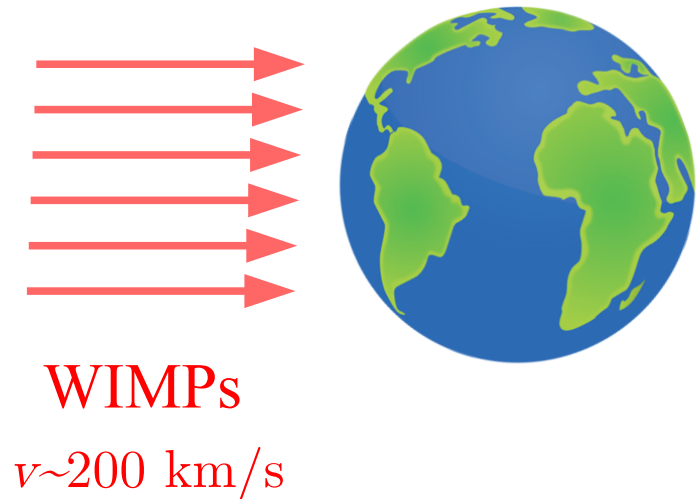
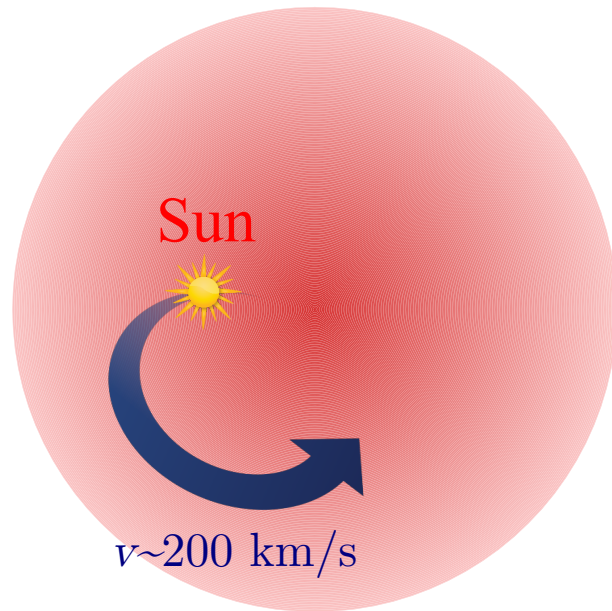


Probing the scattering with nucleons

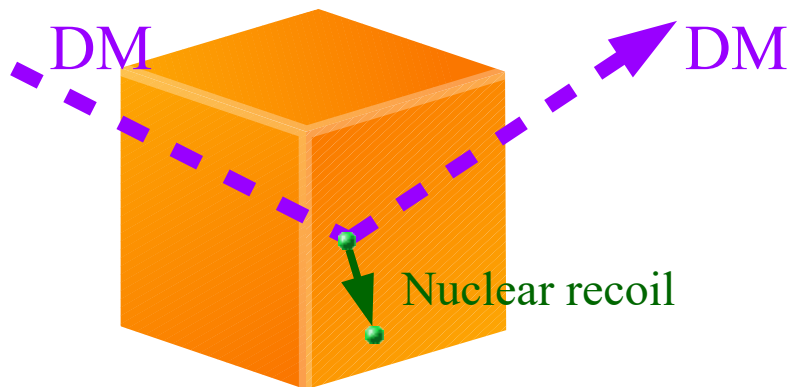


Probing the scattering with nucleons

The Sun (and the Earth) might be moving through a “gas” of dark matter particles.

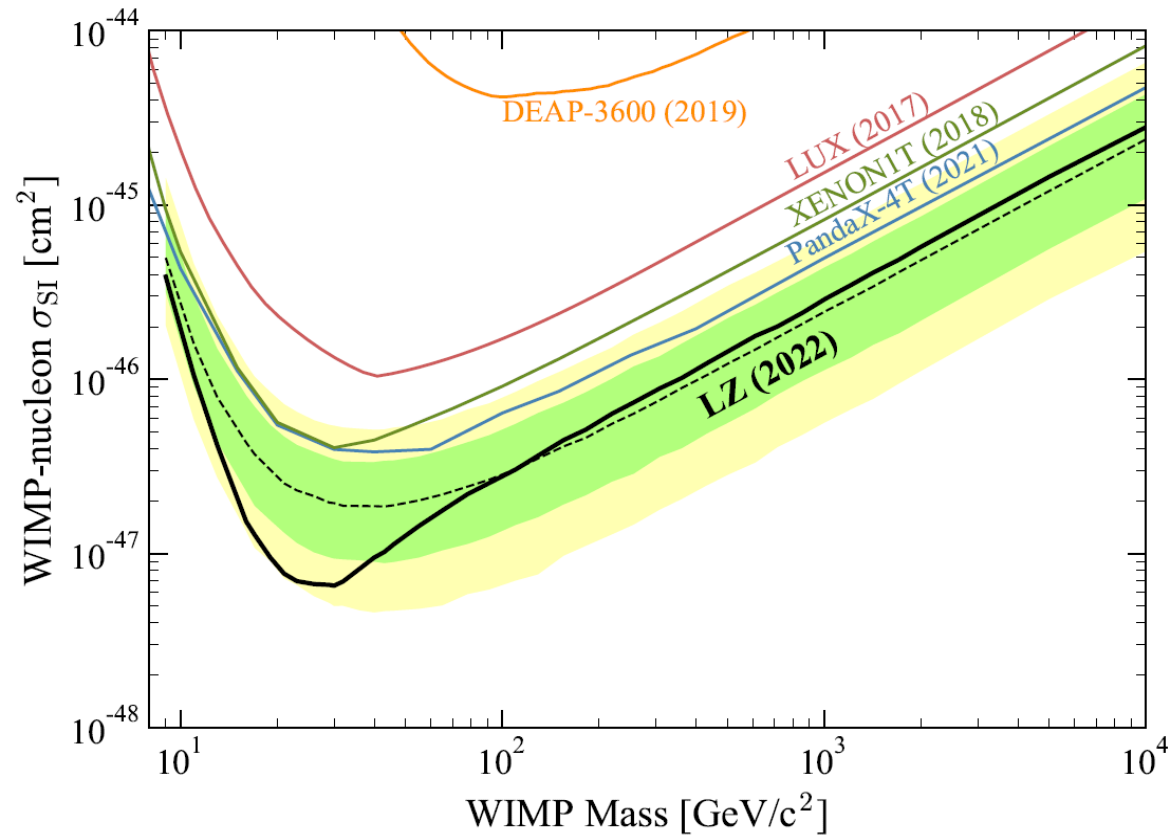


Once in a while a dark matter particle will interact with a nucleus. The nucleus then recoils, producing vibrations, ionizations or scintillation light in the detector.

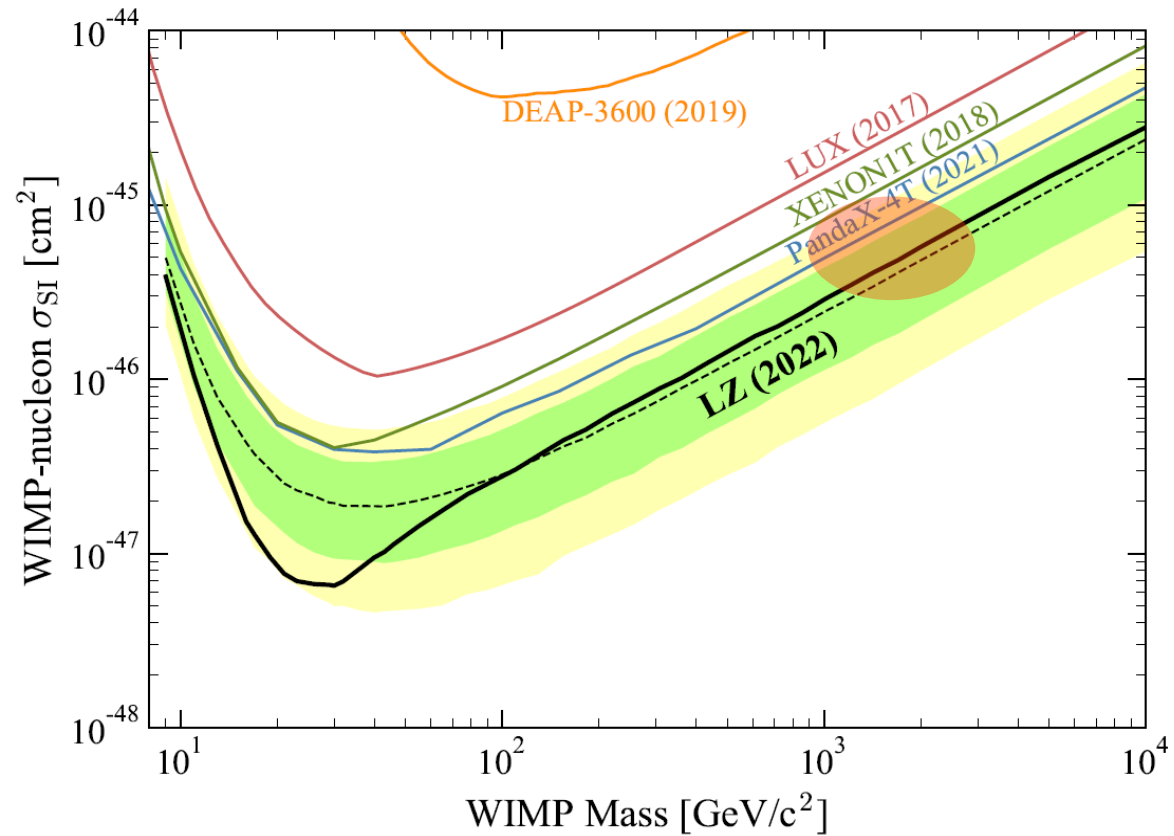


No significant excess detected so far

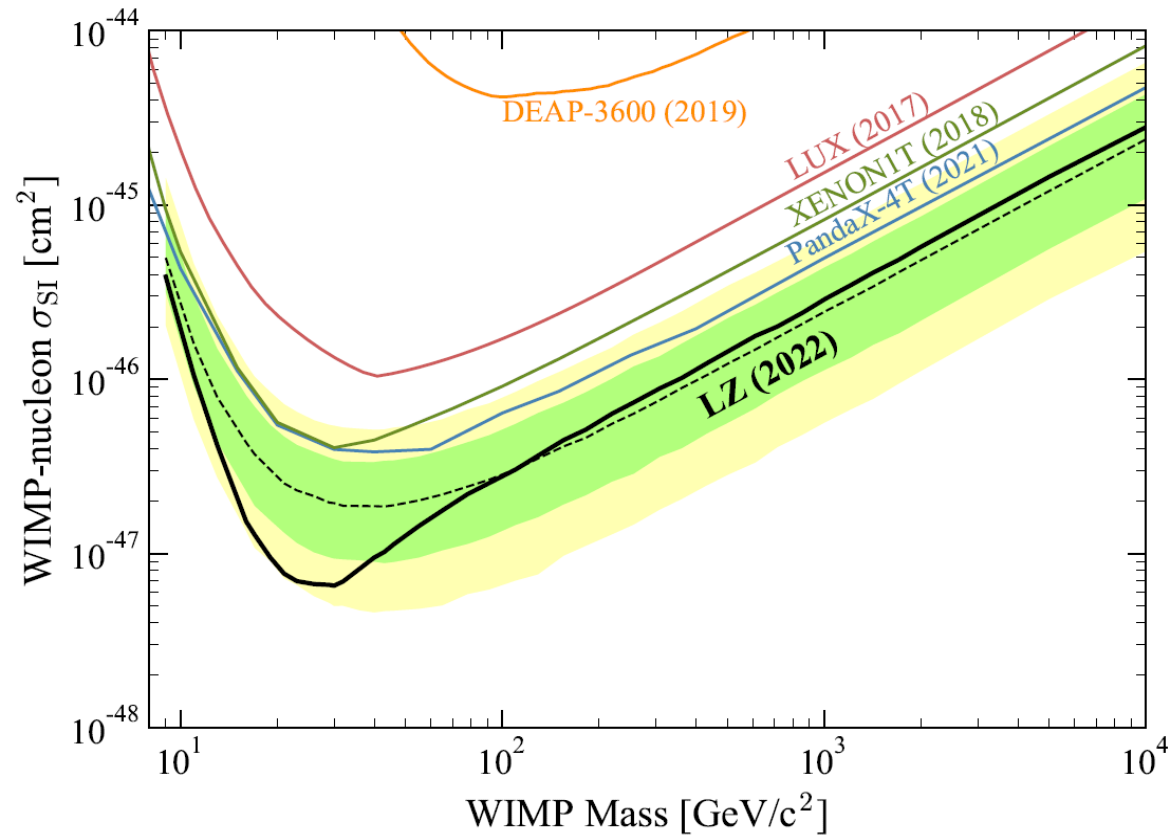
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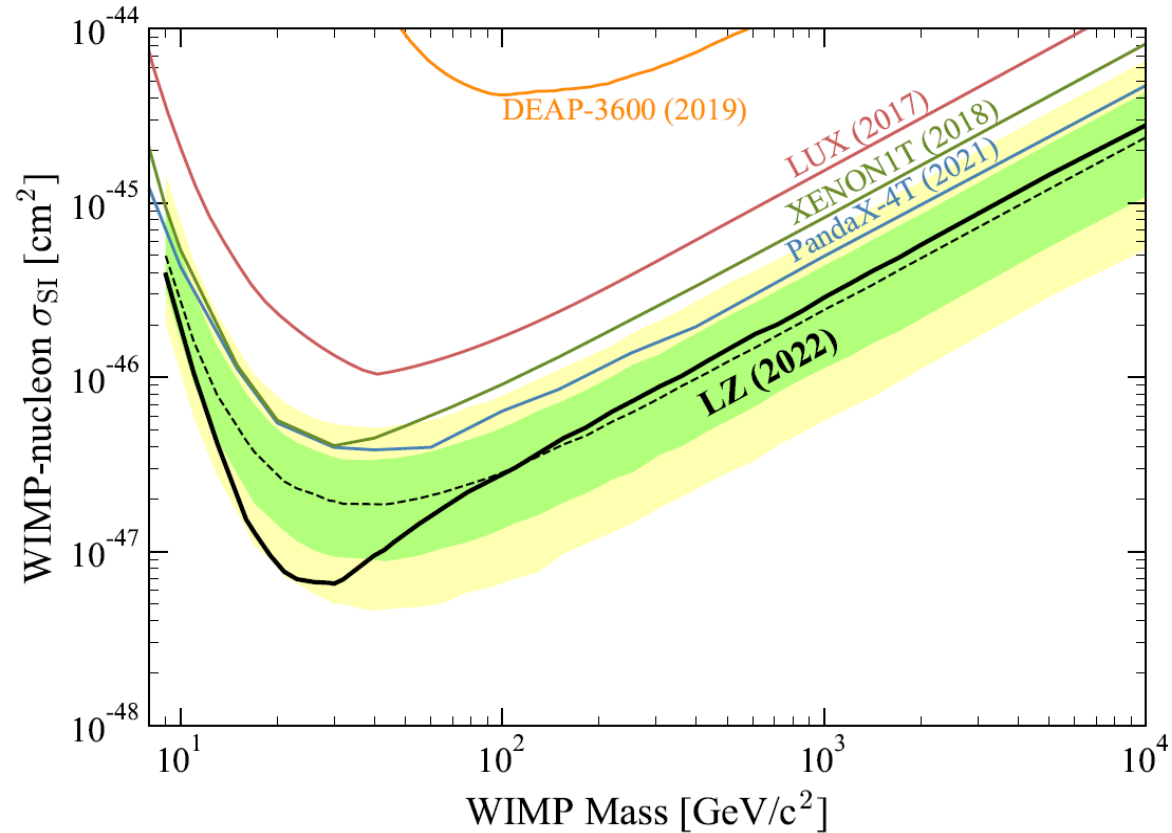
Probing the scattering with nucleons



Smallprint:

- DM interacts only through the spin-independent interaction
- DM couples with equal strength to protons and neutrons (isoscalar)
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- DM velocity distribution given by a Maxwell-Boltzmann, truncated at the escape velocity.

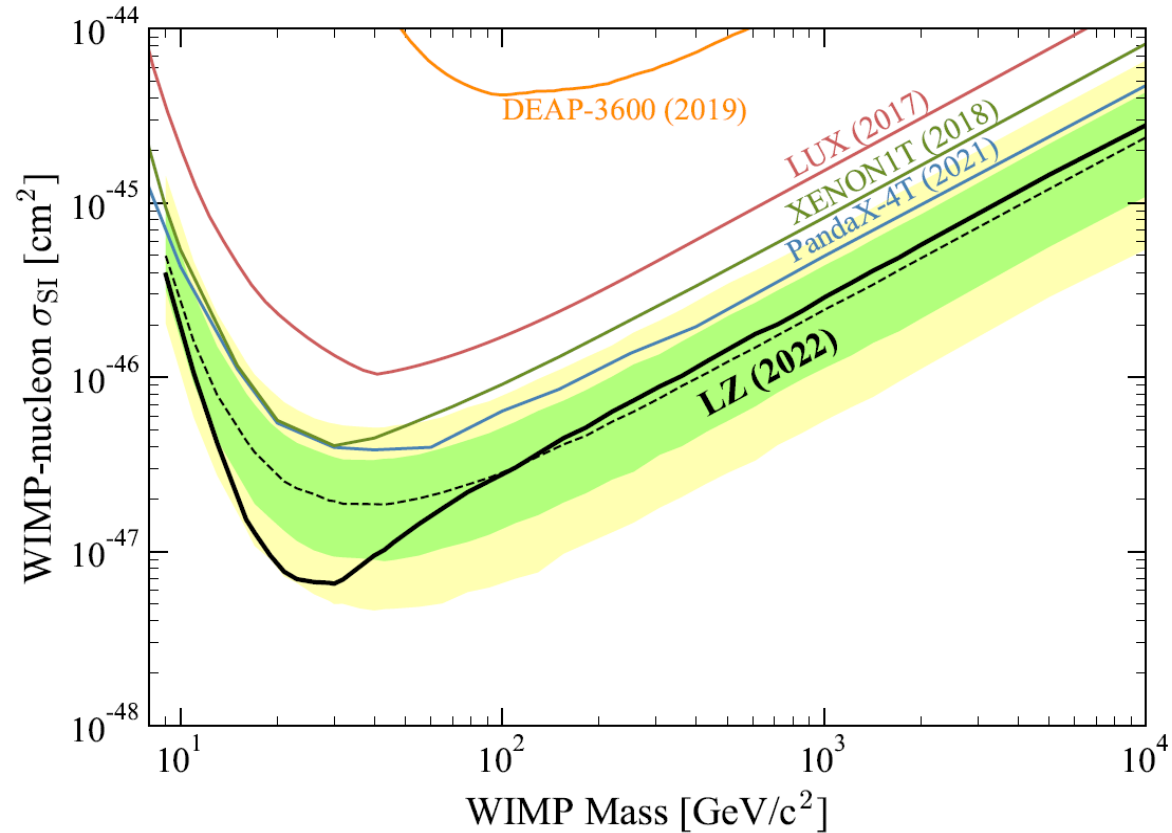
Probing the scattering with nucleons



Differential recoil rate

$$\frac{dR}{dE_R} = \frac{\rho_{\text{loc}}}{m_A m_{\text{DM}}} \int_{v \geq v_{\text{min}}(E_R)} d^3v v f(\vec{v} + \vec{v}_{\text{obs}}(t)) \frac{d\sigma}{dE_R}$$

Probing the scattering with nucleons

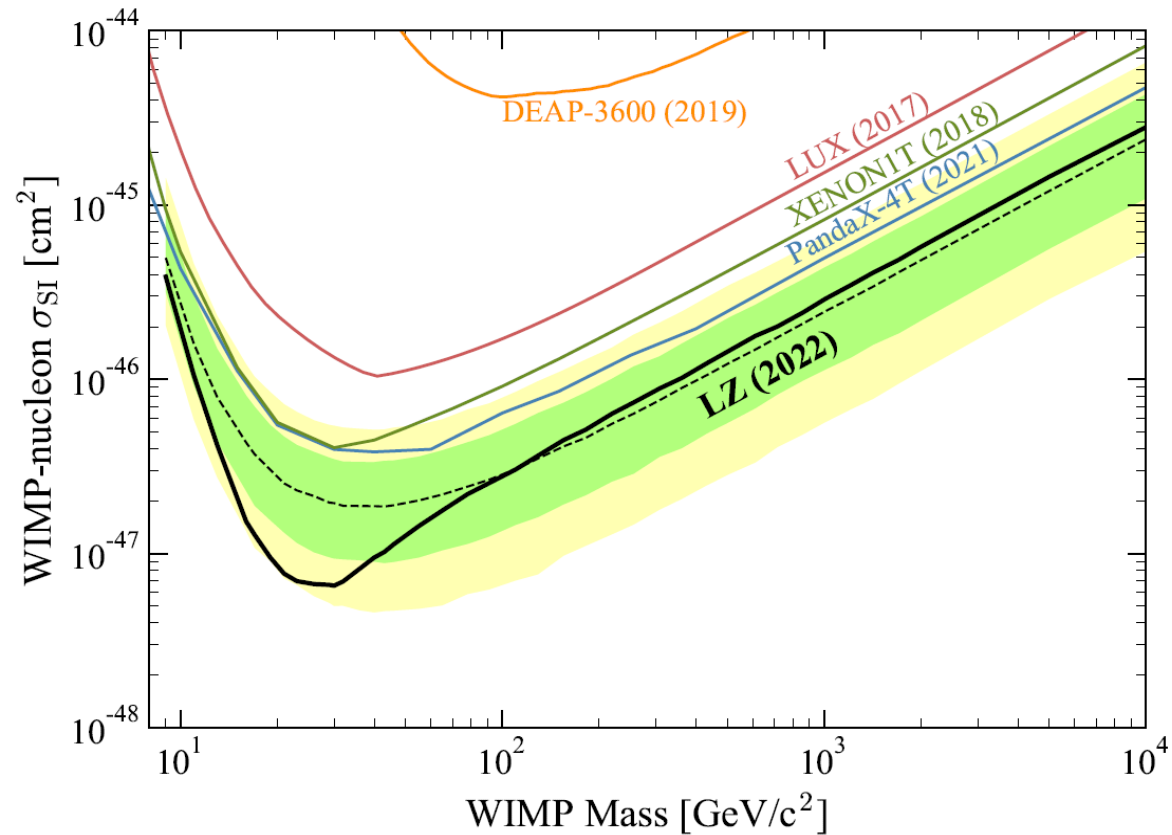


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?

Probing the scattering with nucleons

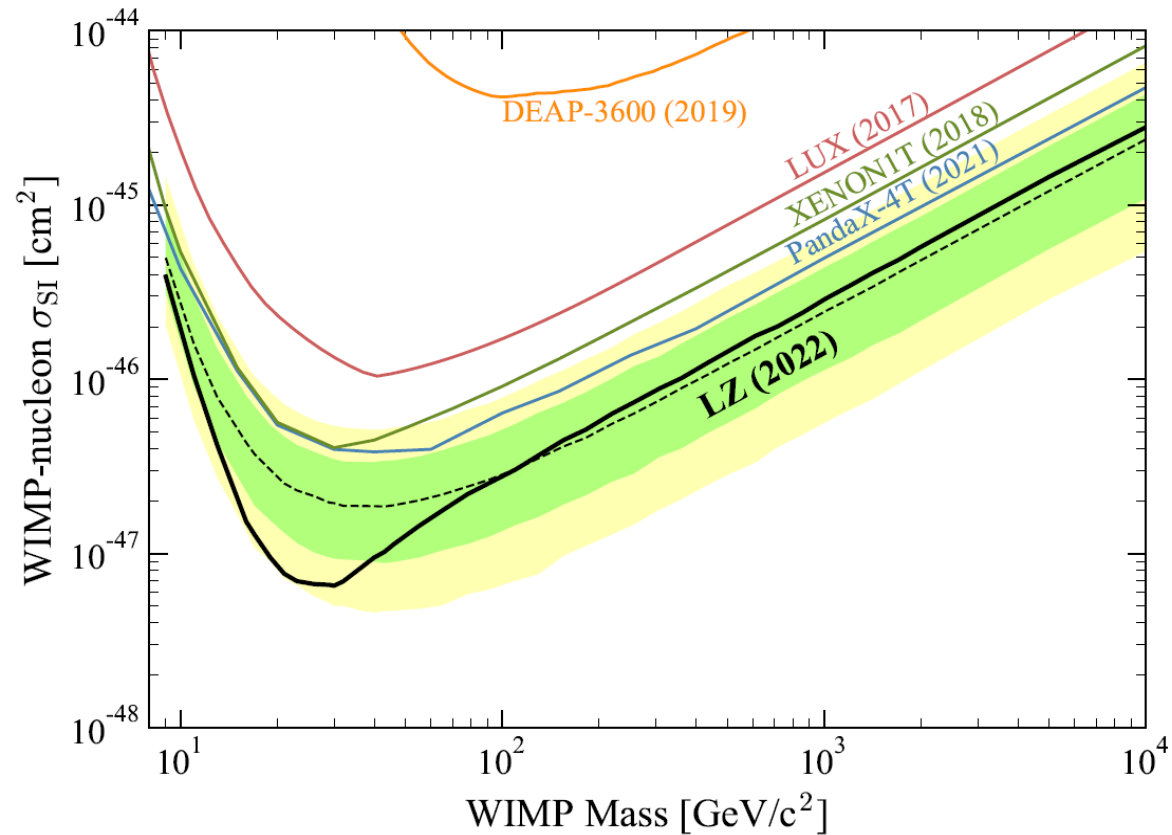


Differential recoil rate

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The equation includes a red question mark above the $\frac{d\sigma}{dE_R}$ term and a blue question mark below the $m_A m_{\text{DM}}$ denominator.

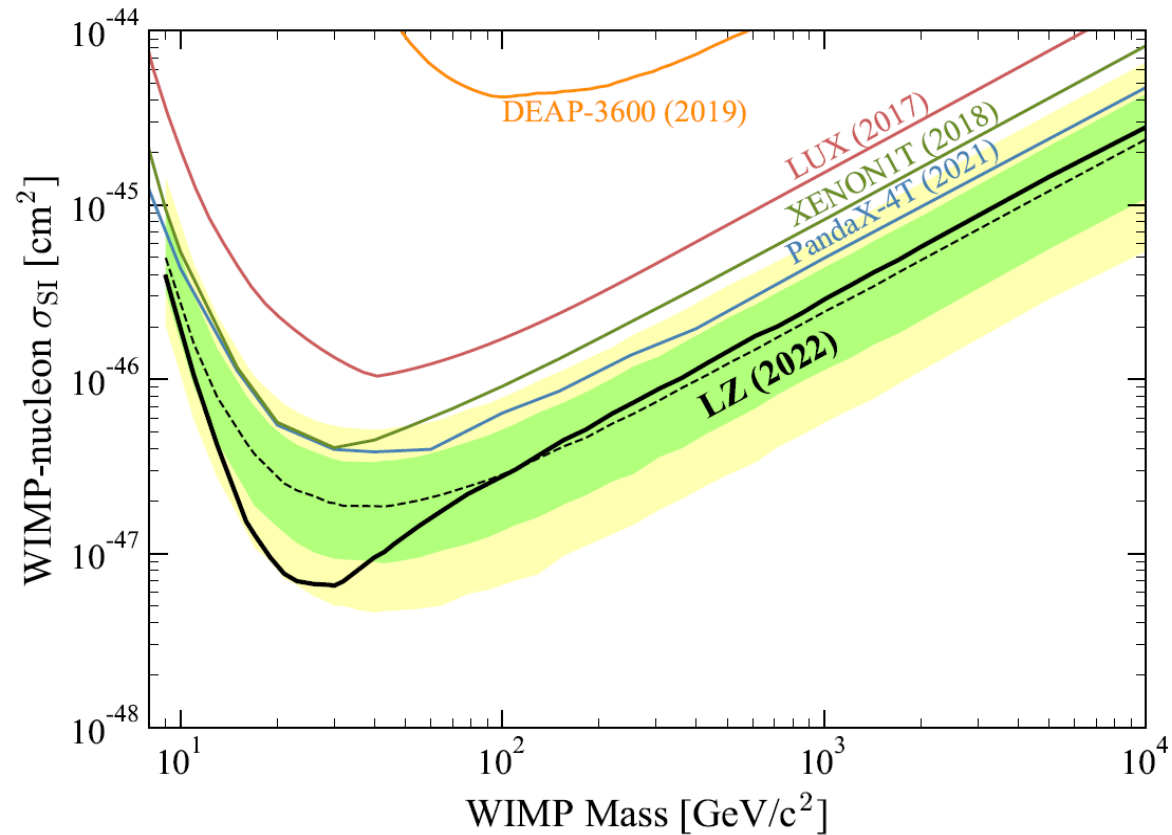
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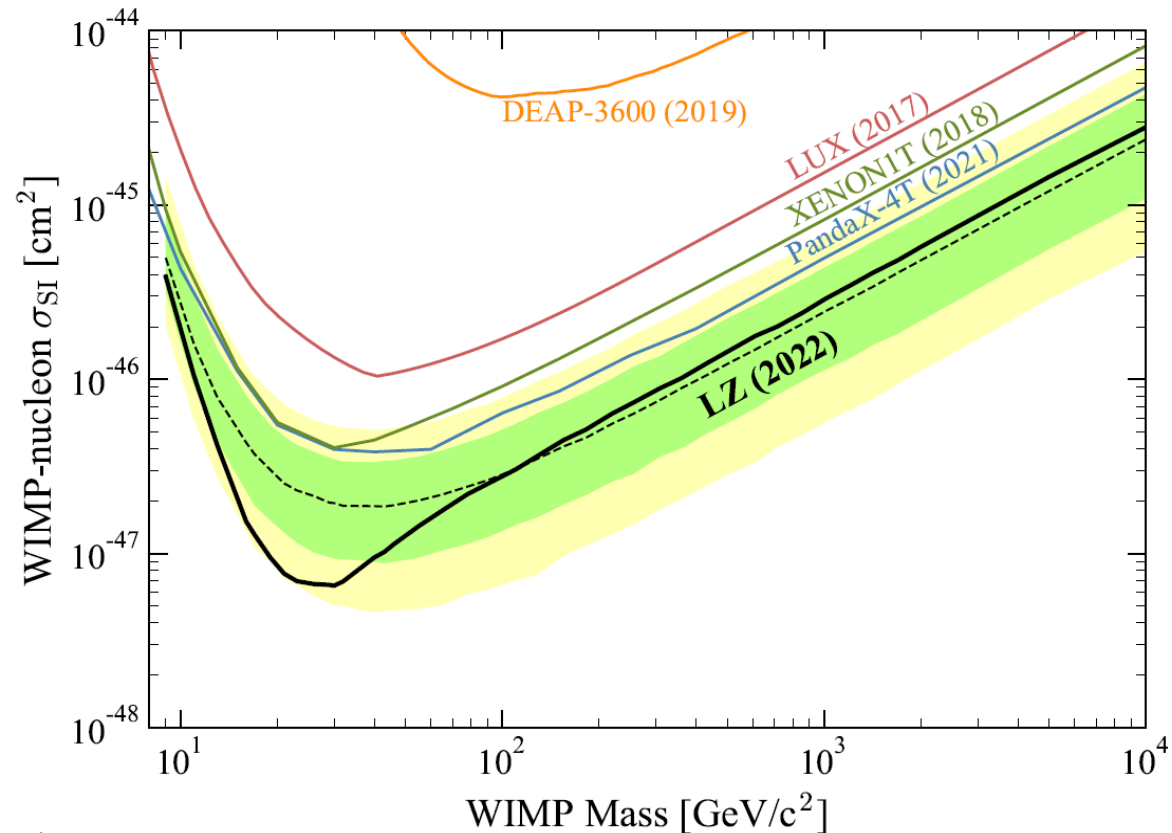
Probing the scattering with nucleons



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Probing the scattering with nucleons



Smallprint:

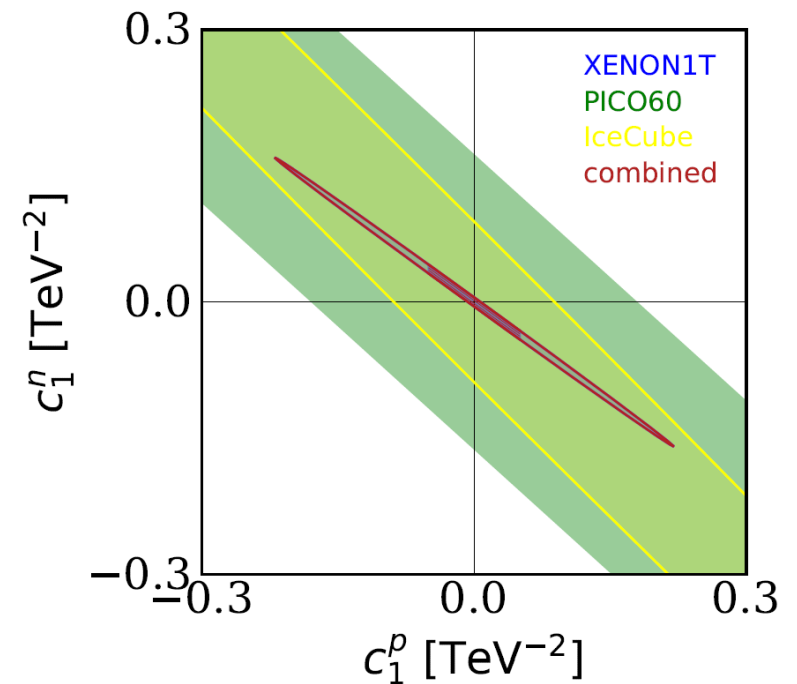
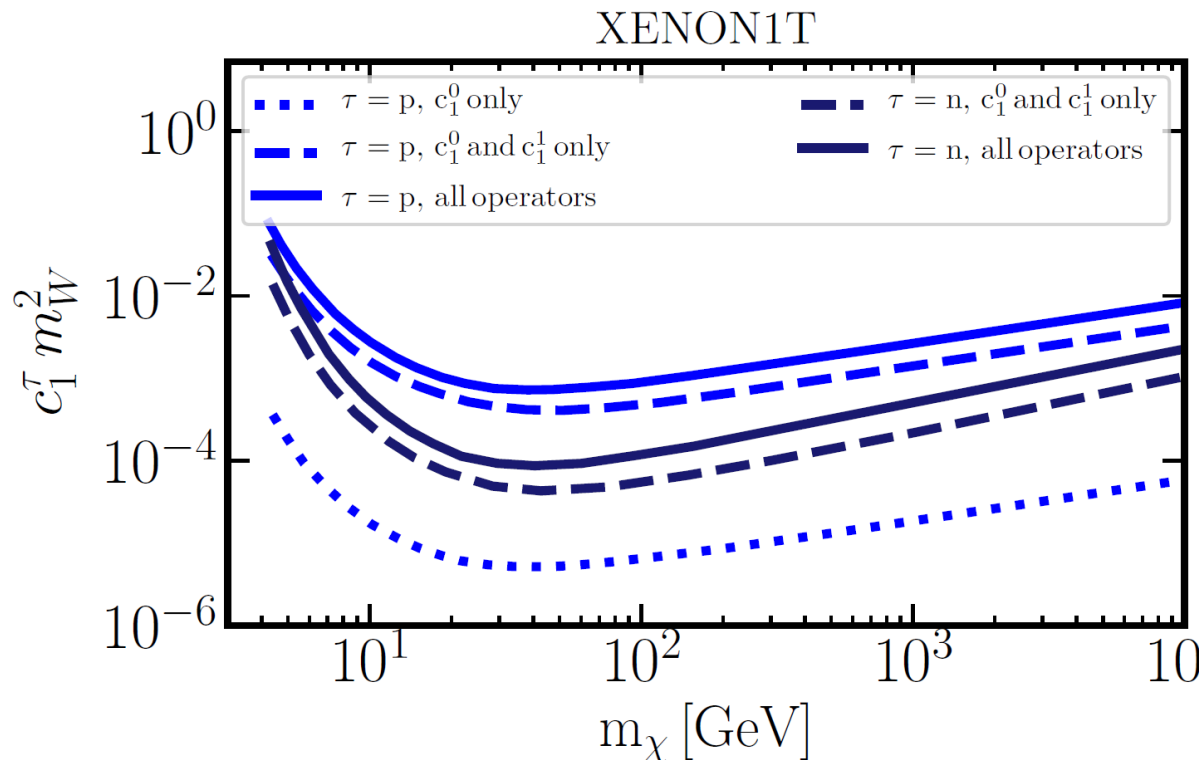
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- Local DM density $\rho=0.3 \text{ GeV}/\text{cm}^3$.
- DM velocity distribution given by a Maxwell-Boltzmann, truncated at the escape velocity.

Probing the scattering with nucleons

Consider the Hamiltonian of the SI interaction:

$$\mathcal{H} = c_p(\bar{\chi}p)(\bar{p}\chi) + c_n(\bar{\chi}n)(\bar{n}\chi)$$

The two interactions can interfere.



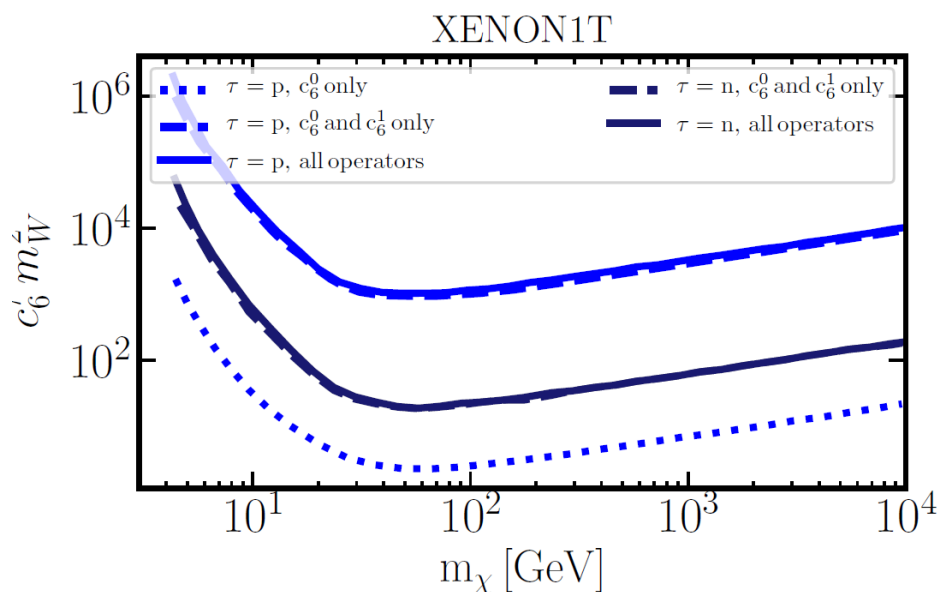
Brenner et al '22

Probing the scattering with nucleons

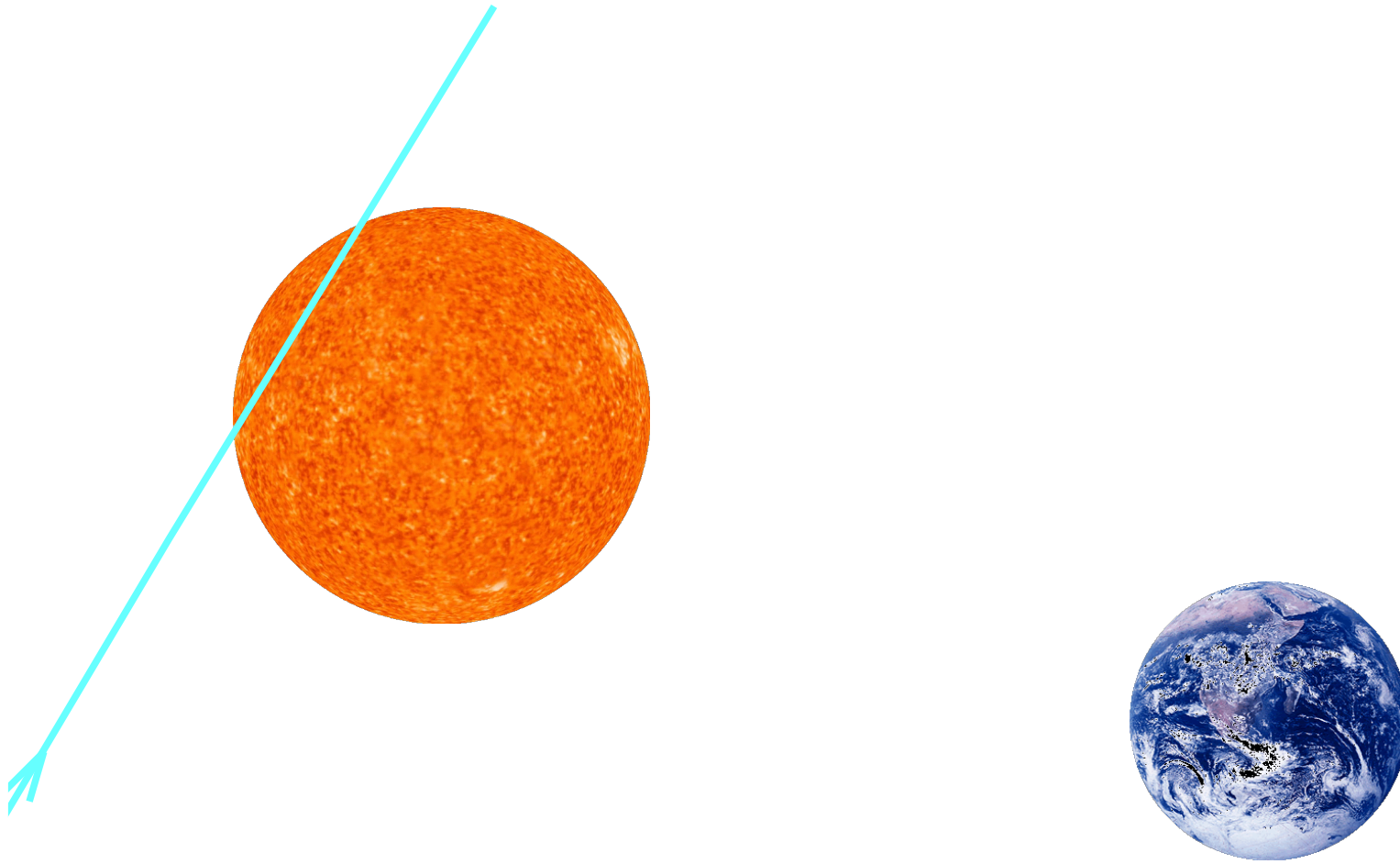
$\mathcal{O}_1 = 1_\chi 1_N$	$\mathcal{O}_9 = i\vec{S}_\chi \cdot (\vec{S}_N \times \frac{\vec{q}}{m_N})$
$\mathcal{O}_3 = i\vec{S}_N \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$	$\mathcal{O}_{10} = i\vec{S}_N \cdot \frac{\vec{q}}{m_N}$
$\mathcal{O}_4 = \vec{S}_\chi \cdot \vec{S}_N$	$\mathcal{O}_{11} = i\vec{S}_\chi \cdot \frac{\vec{q}}{m_N}$
$\mathcal{O}_5 = i\vec{S}_\chi \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$	$\mathcal{O}_{12} = \vec{S}_\chi \cdot (\vec{S}_N \times \vec{v}^\perp)$
$\mathcal{O}_6 = (\vec{S}_\chi \cdot \frac{\vec{q}}{m_N})(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$	$\mathcal{O}_{13} = i(\vec{S}_\chi \cdot \vec{v}^\perp)(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$
$\mathcal{O}_7 = \vec{S}_N \cdot \vec{v}^\perp$	$\mathcal{O}_{14} = i(\vec{S}_\chi \cdot \frac{\vec{q}}{m_N})(\vec{S}_N \cdot \vec{v}^\perp)$
$\mathcal{O}_8 = \vec{S}_\chi \cdot \vec{v}^\perp$	$\mathcal{O}_{15} = -(\vec{S}_\chi \cdot \frac{\vec{q}}{m_N})((\vec{S}_N \times \vec{v}^\perp) \cdot \frac{\vec{q}}{m_N})$

Probing the scattering with nucleons

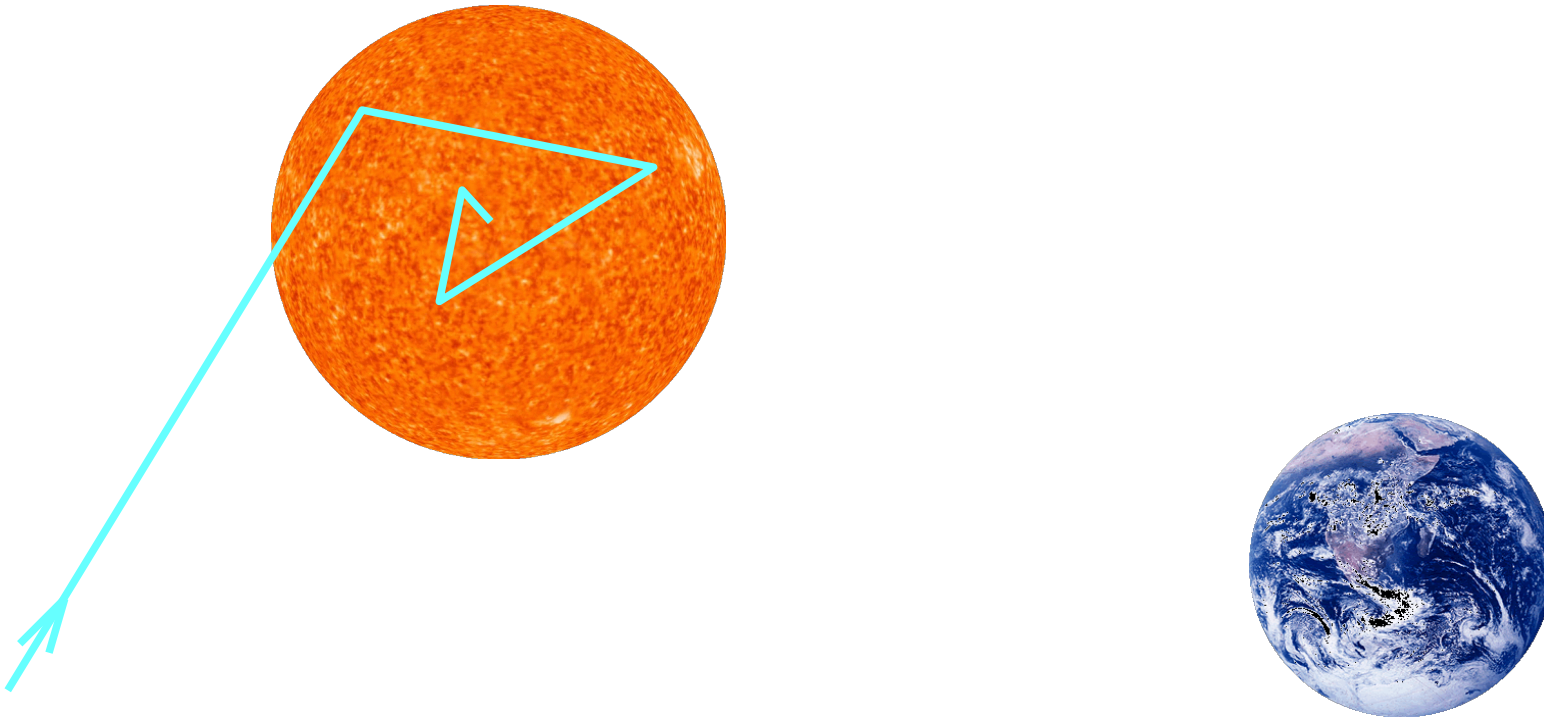
$\mathcal{O}_1 = 1_X 1_N$	$\mathcal{O}_9 = i\vec{S}_X \cdot (\vec{S}_N \times \frac{\vec{q}}{m_N})$
$\mathcal{O}_3 = i\vec{S}_N \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$	$\mathcal{O}_{10} = i\vec{S}_N \cdot \frac{\vec{q}}{m_N}$
$\mathcal{O}_4 = \vec{S}_X \cdot \vec{S}_N$	$\mathcal{O}_{11} = i\vec{S}_X \cdot \frac{\vec{q}}{m_N}$
$\mathcal{O}_5 = i\vec{S}_X \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$	$\mathcal{O}_{12} = \vec{S}_X \cdot (\vec{S}_N \times \vec{v}^\perp)$
$\mathcal{O}_6 = (\vec{S}_X \cdot \frac{\vec{q}}{m_N})(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$	$\mathcal{O}_{13} = i(\vec{S}_X \cdot \vec{v}^\perp)(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$
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Using the Sun to probe scattering with nucleons

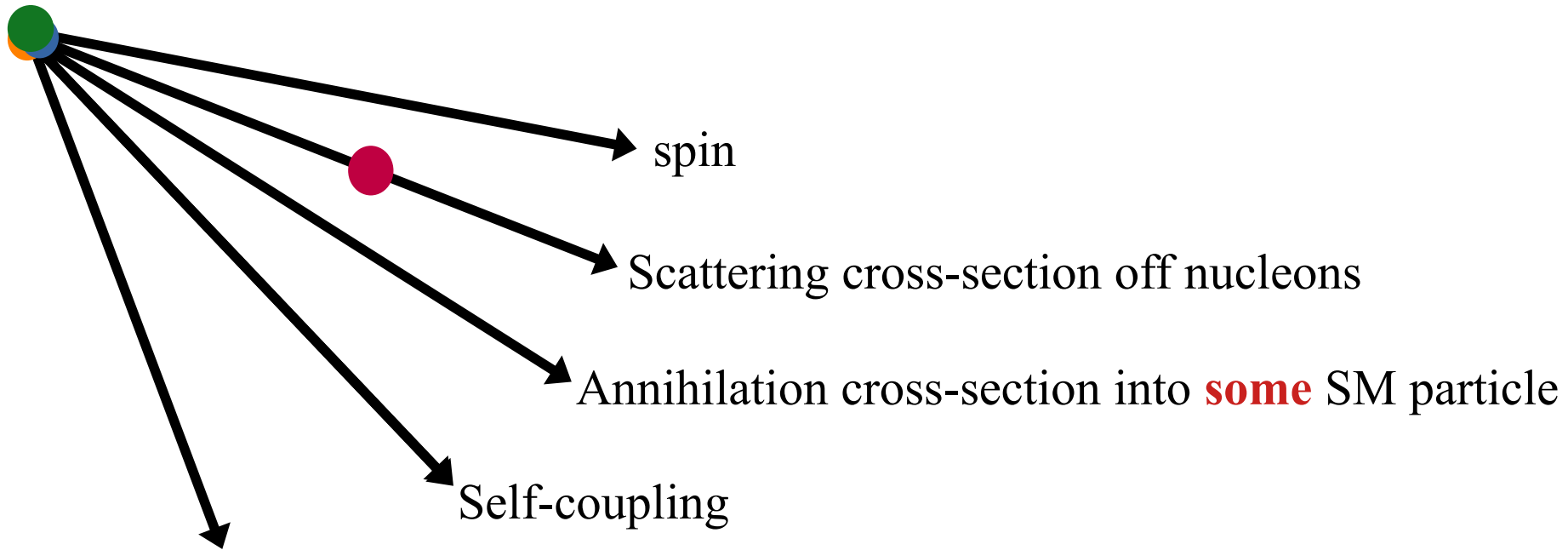


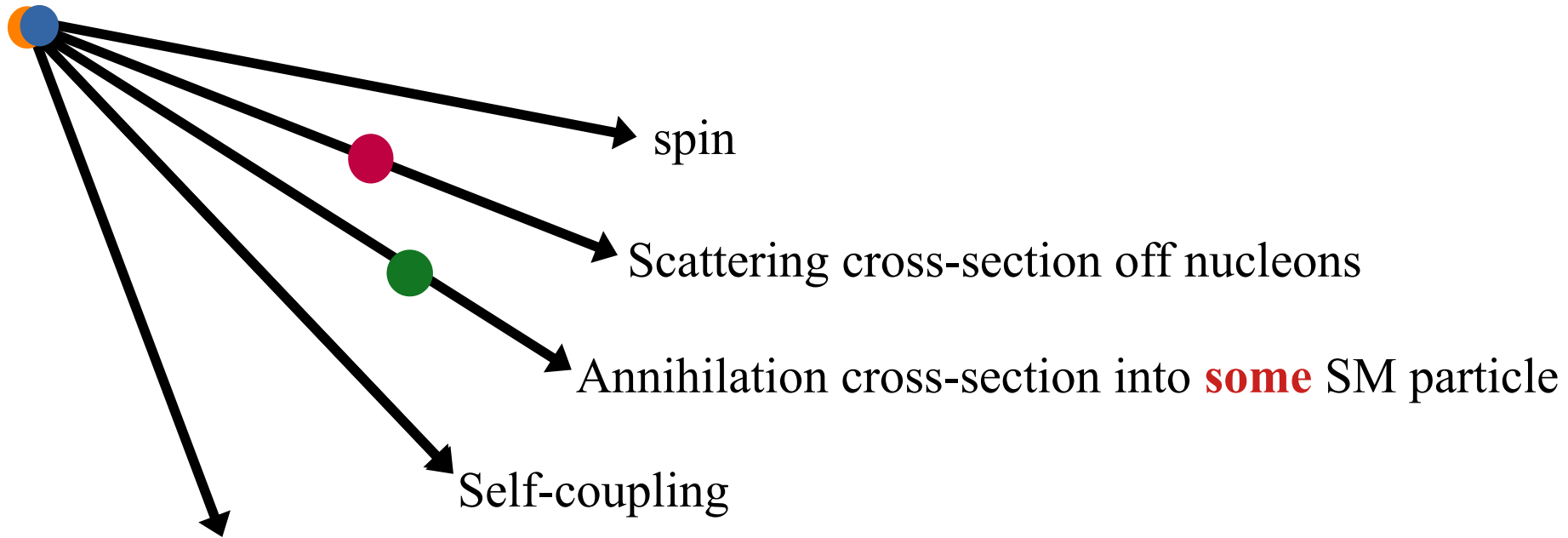
Using the Sun to probe scattering with nucleons



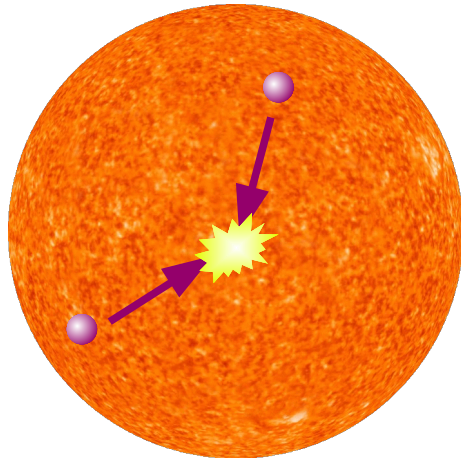
Using the Sun to probe scattering with nucleons



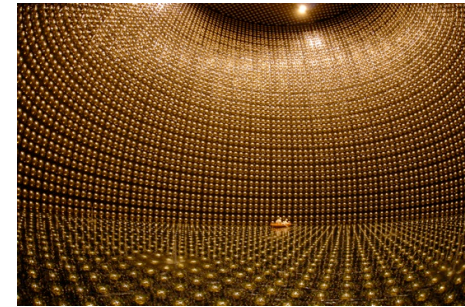
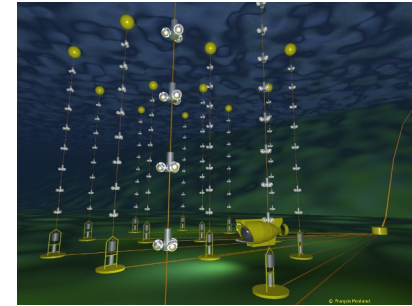
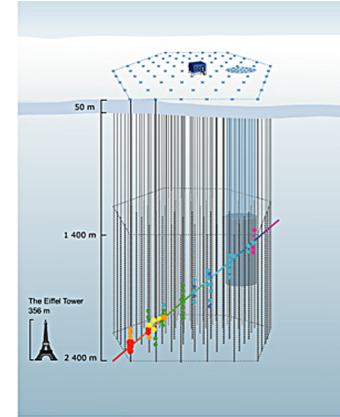
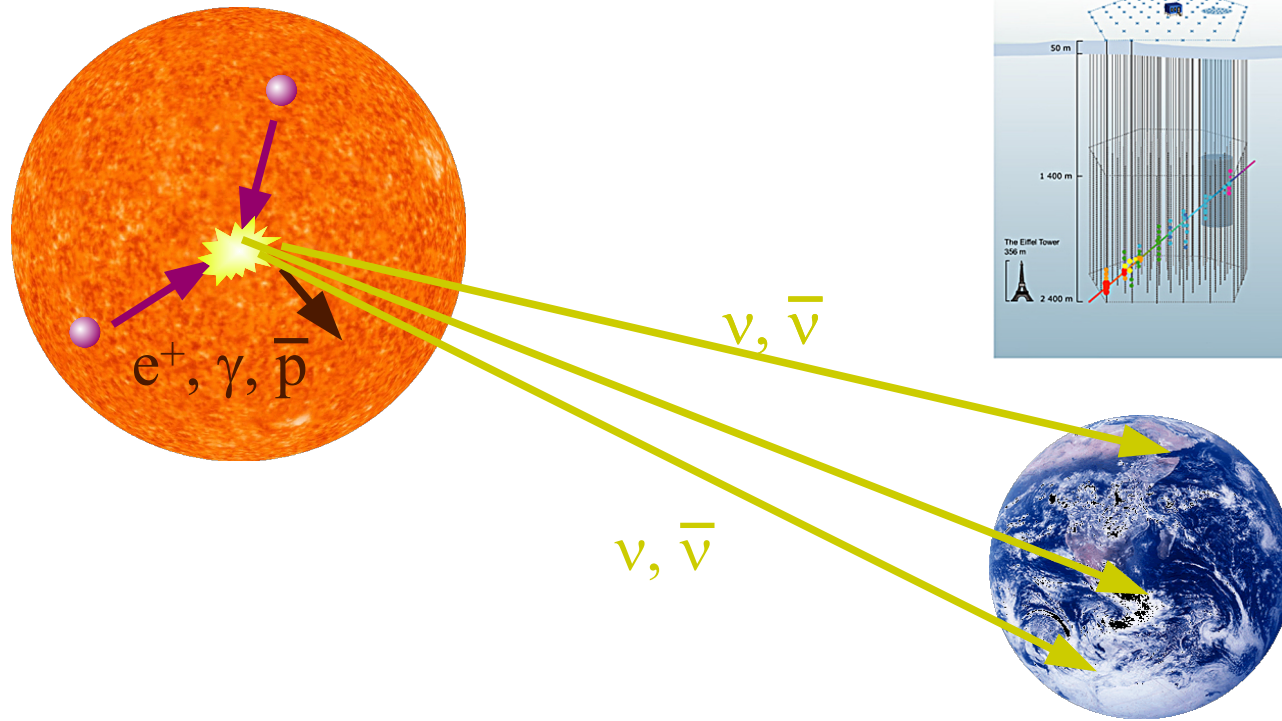




Using the Sun to probe scattering with nucleons

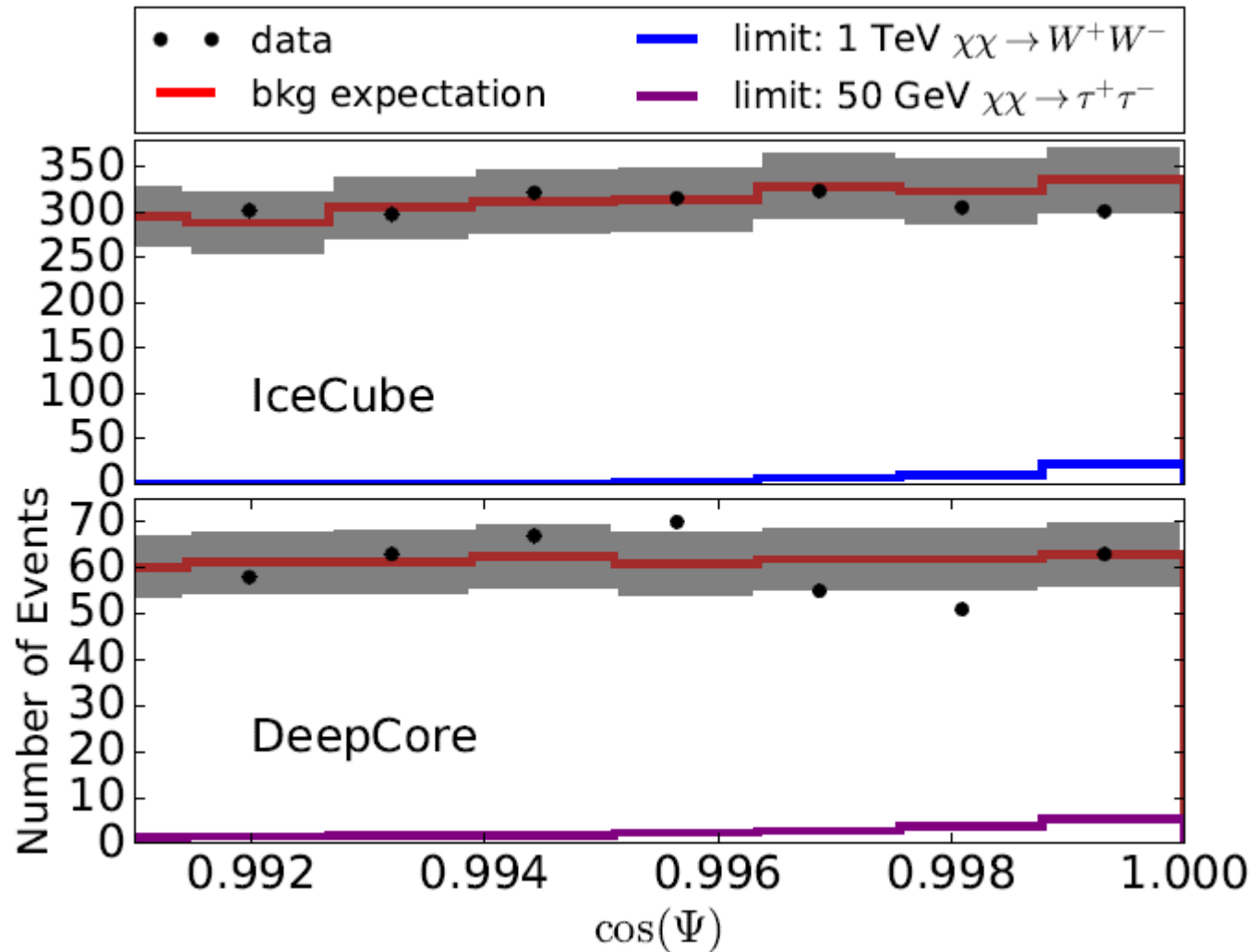


Using the Sun to probe scattering with nucleons

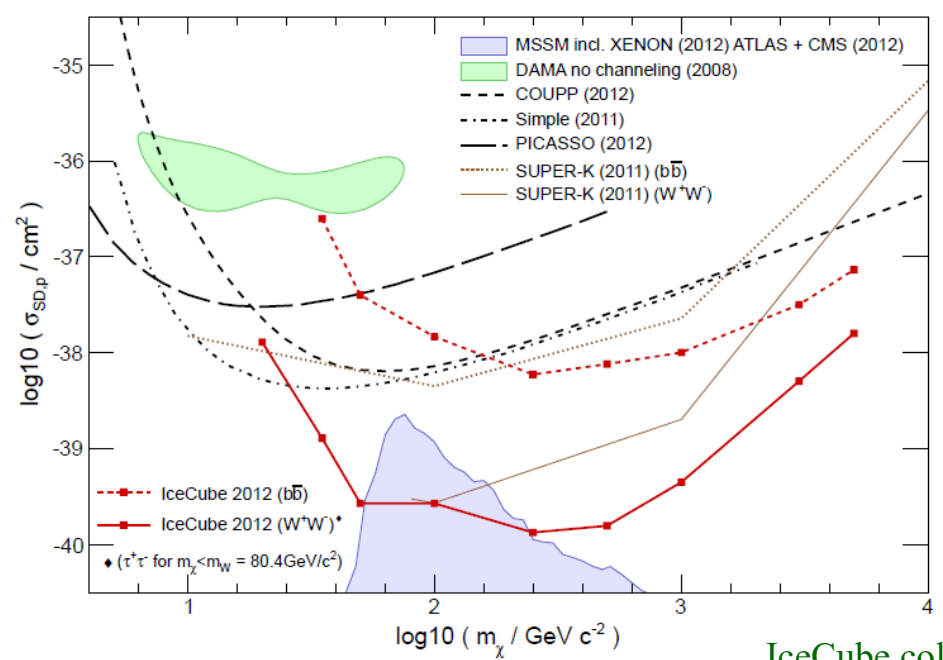
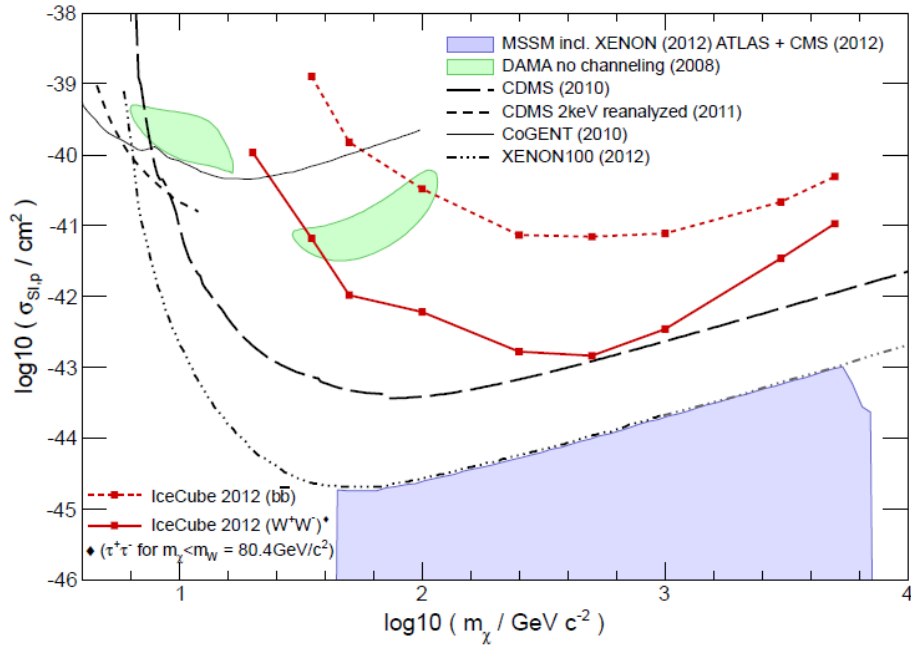


Using the Sun to probe scattering with nucleons

Observations consistent with the background-only hypothesis

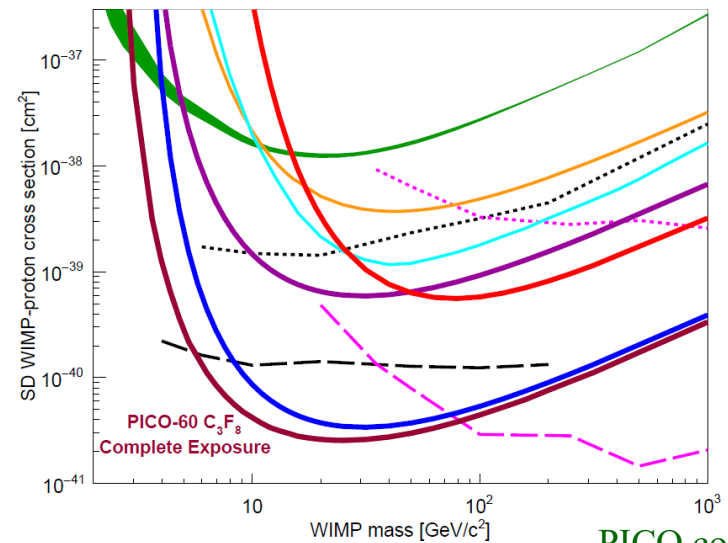
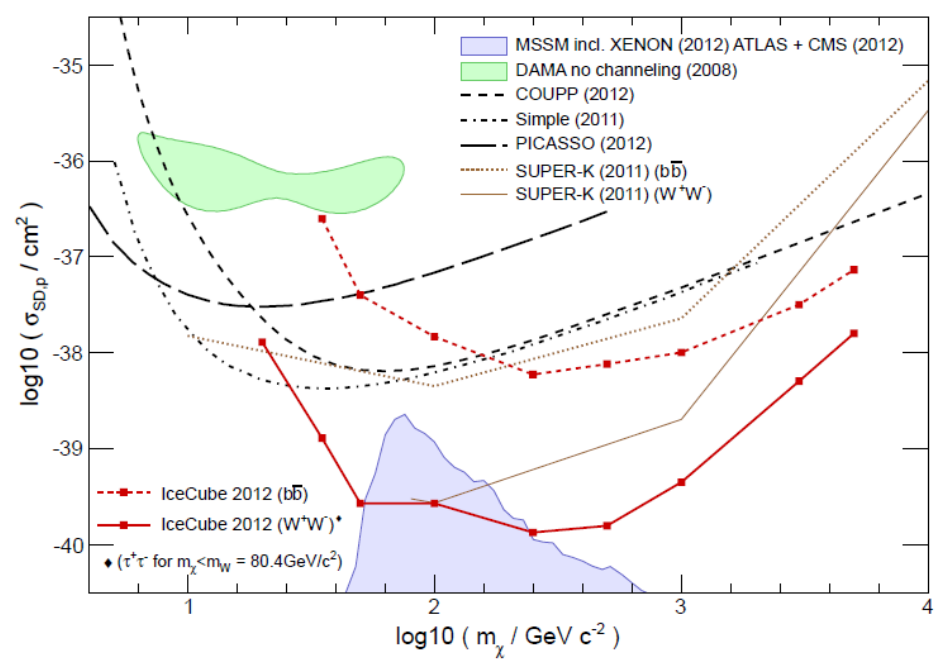
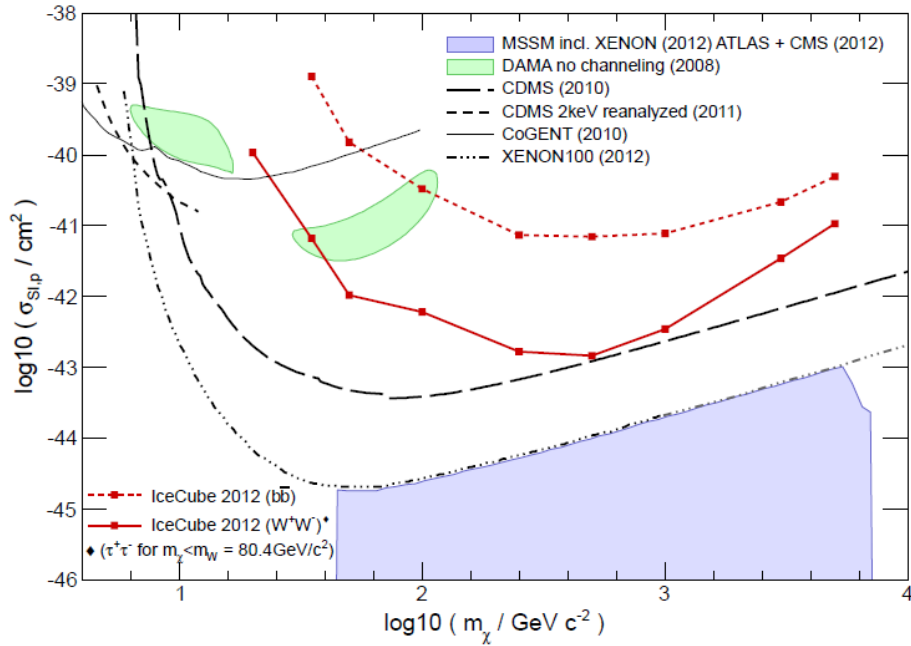


Using the Sun to probe scattering with nucleons



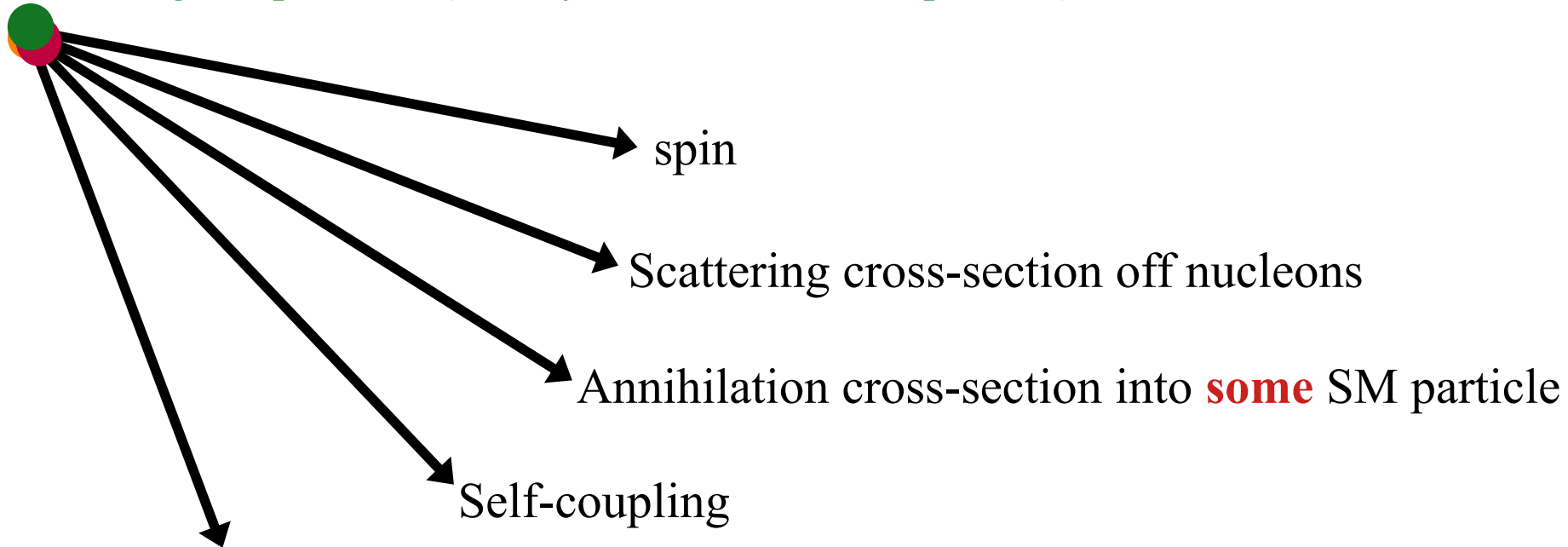
IceCube coll'12

Using the Sun to probe scattering with nucleons



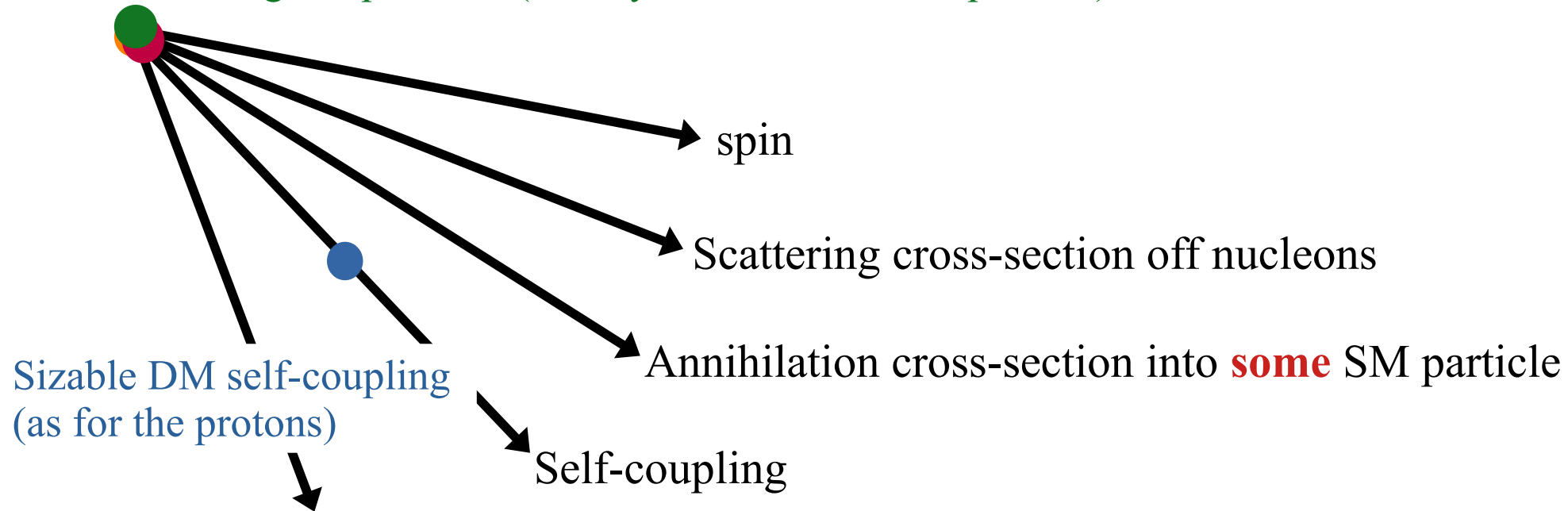
Probing alternative scenarios: ADM

“Asymmetric dark matter”: a symmetry protects DM against annihilations into lighter particles (as baryon number for the protons)



Probing alternative scenarios: ADM

“Asymmetric dark matter”: a symmetry protects DM against annihilations into lighter particles (as baryon number for the protons)



Probing alternative scenarios: ADM

Protons do not annihilate.

Protons have strong self-interactions

Protons form stars

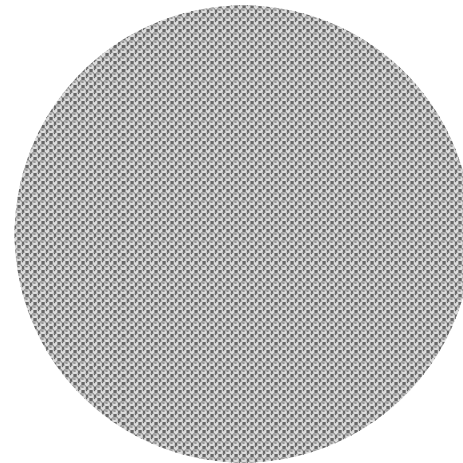


Probing alternative scenarios: ADM

Protons do not annihilate.
Protons have strong self-interactions
Protons form stars



DM does not annihilate.
DM has strong self-interactions
DM form dark stars



Probing alternative scenarios: ADM

Density profile of dark stars calculable from the Klein-Gordon equation in curved spacetime (for bosonic DM) and the Einstein equations:

Colpi et al'86

$$g^{\mu\nu} \nabla_{\mu} \nabla_{\nu} \phi - m^2 \phi - \lambda |\phi|^2 \phi = 0$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G T_{\mu\nu}$$

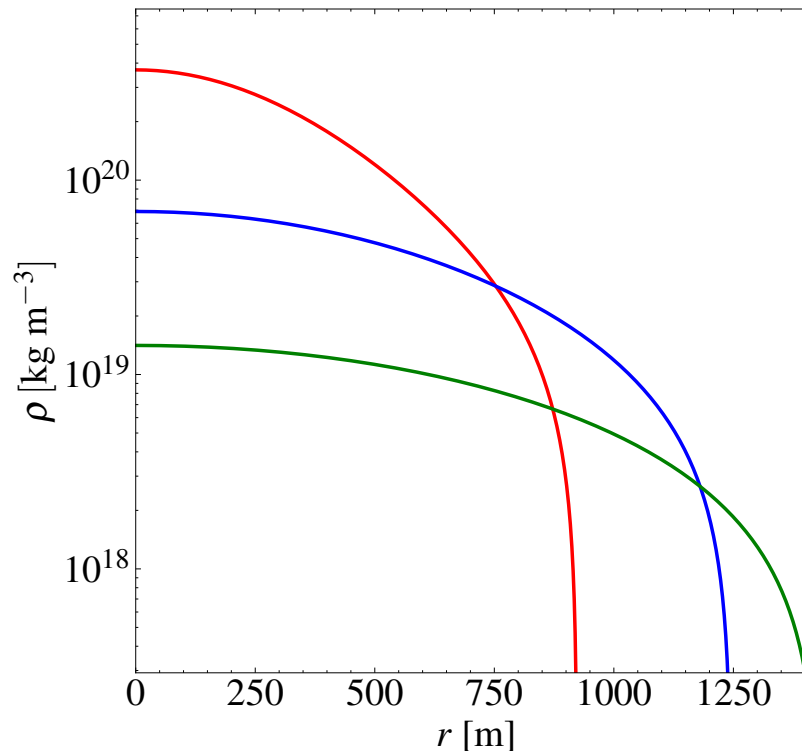
Probing alternative scenarios: ADM

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(For $m=1$ GeV, $\lambda=1$)

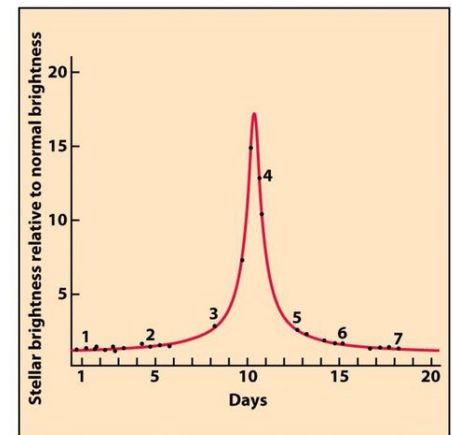
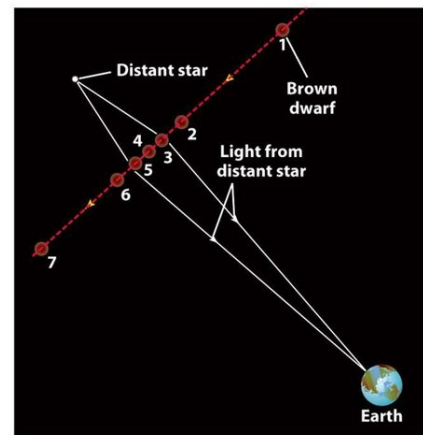
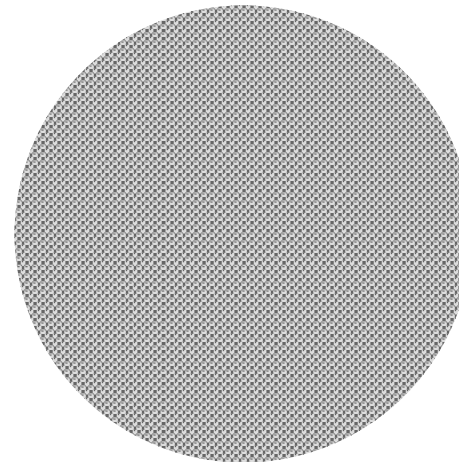
Dark stars are very compact objects

Probing alternative scenarios: ADM

Protons do not annihilate.
Protons have strong self-interactions
Protons form stars

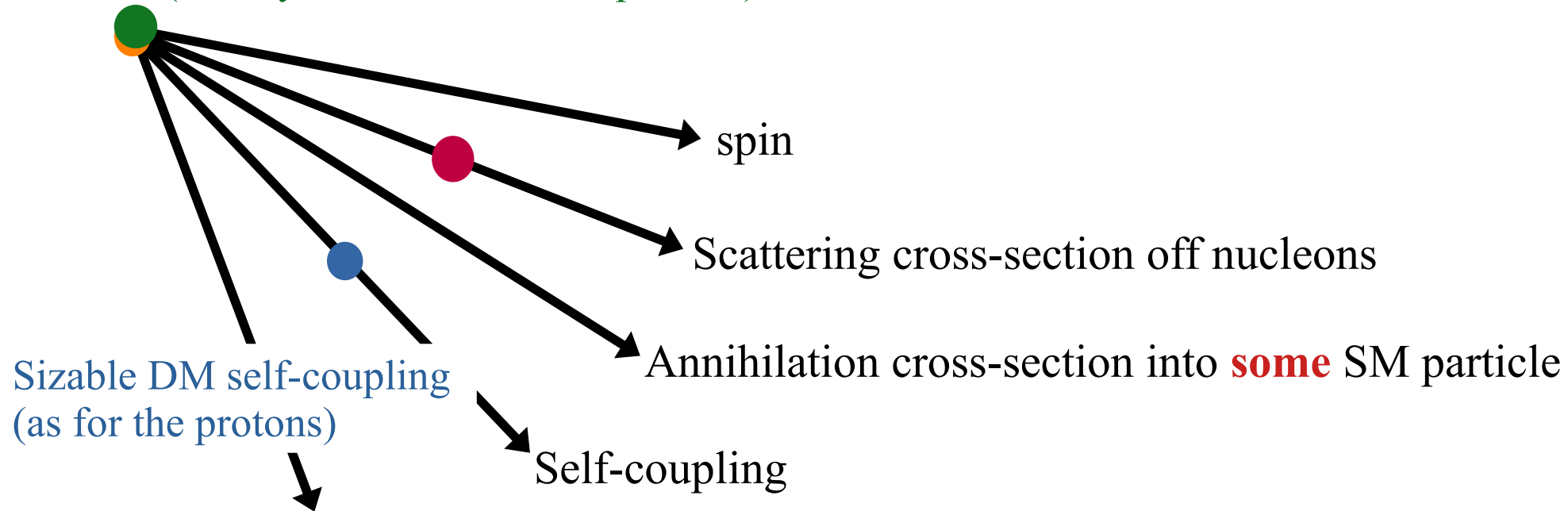


DM does not annihilate.
DM has strong self-interactions
DM form dark stars

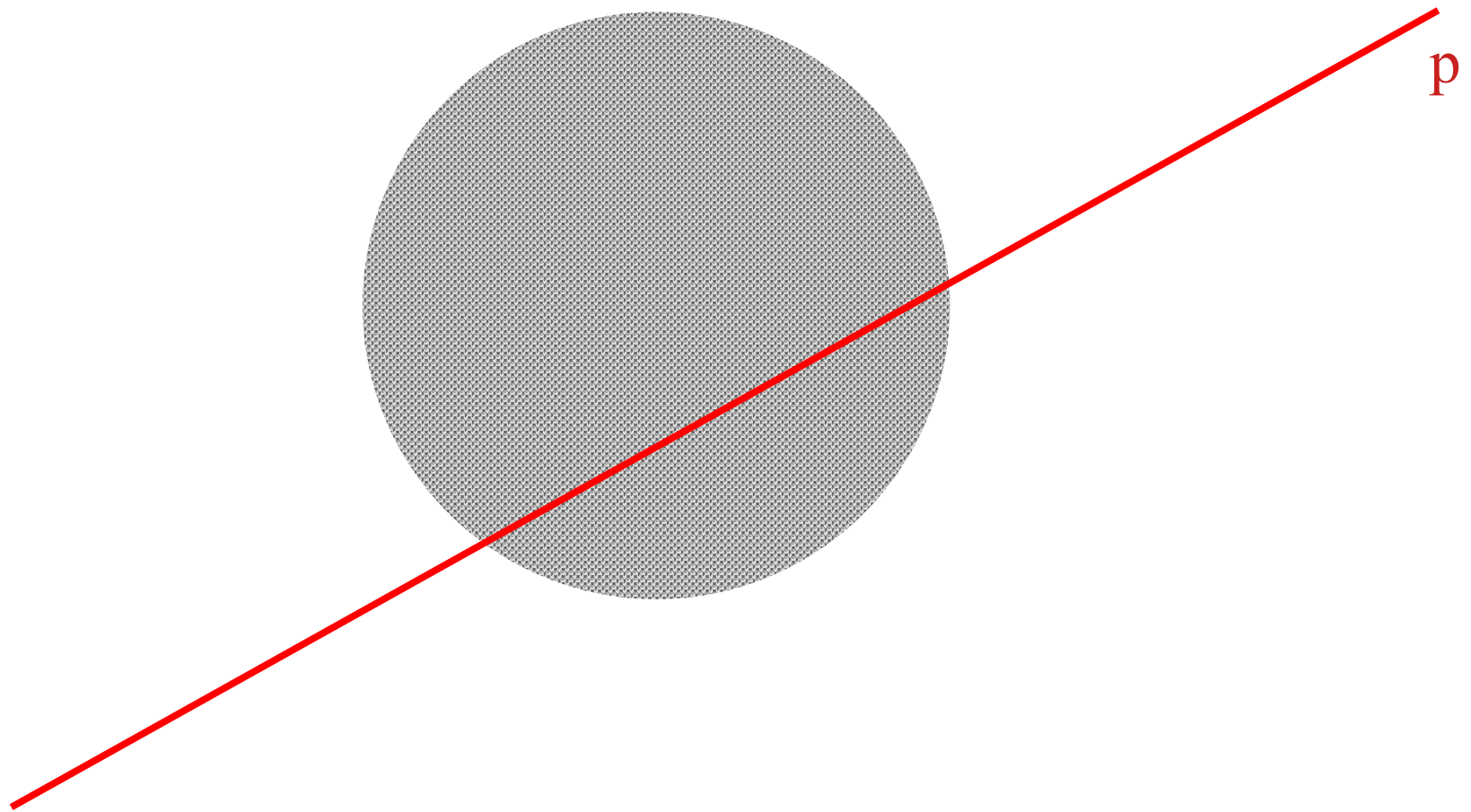


Using DSs to probe scattering with nucleons

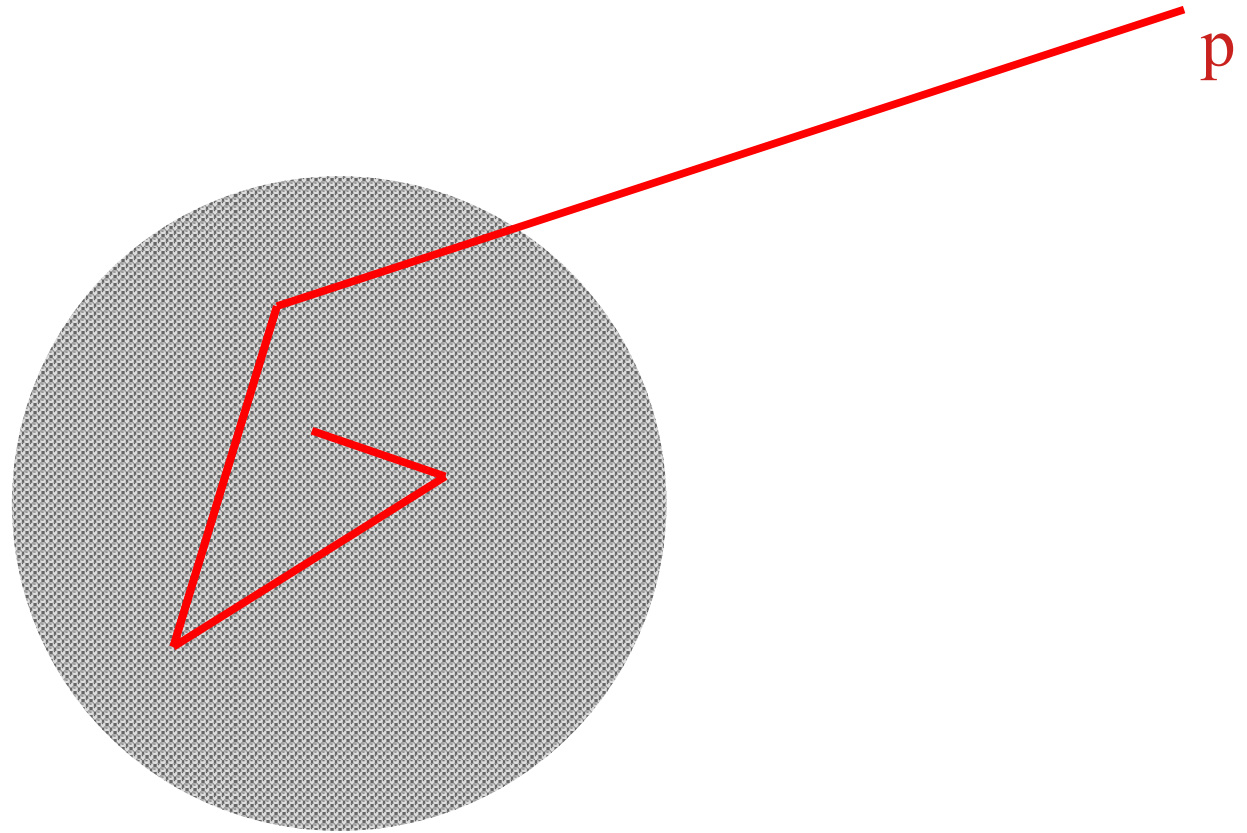
“Asymmetric dark matter”: a symmetry protects DM against annihilations
(as baryon number for the protons)



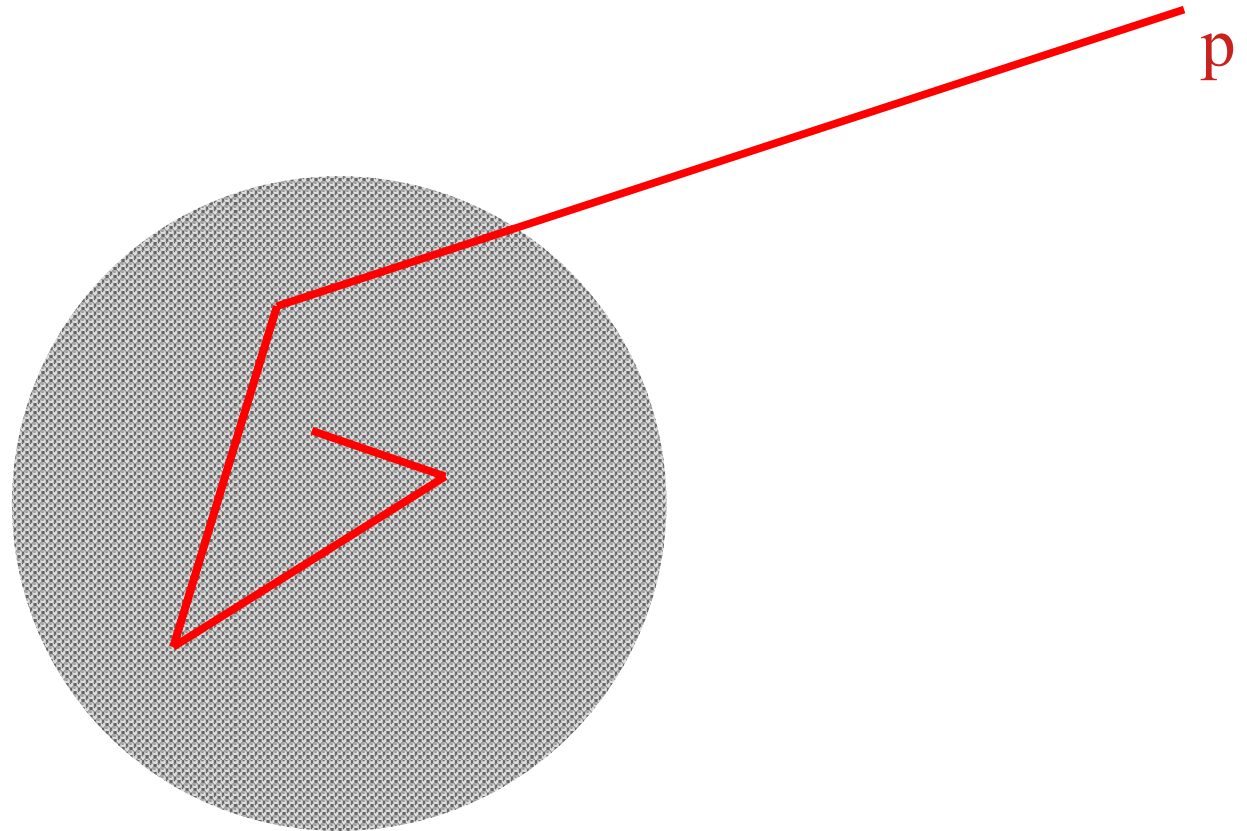
Using DSs to probe scattering with nucleons



Using DSs to probe scattering with nucleons

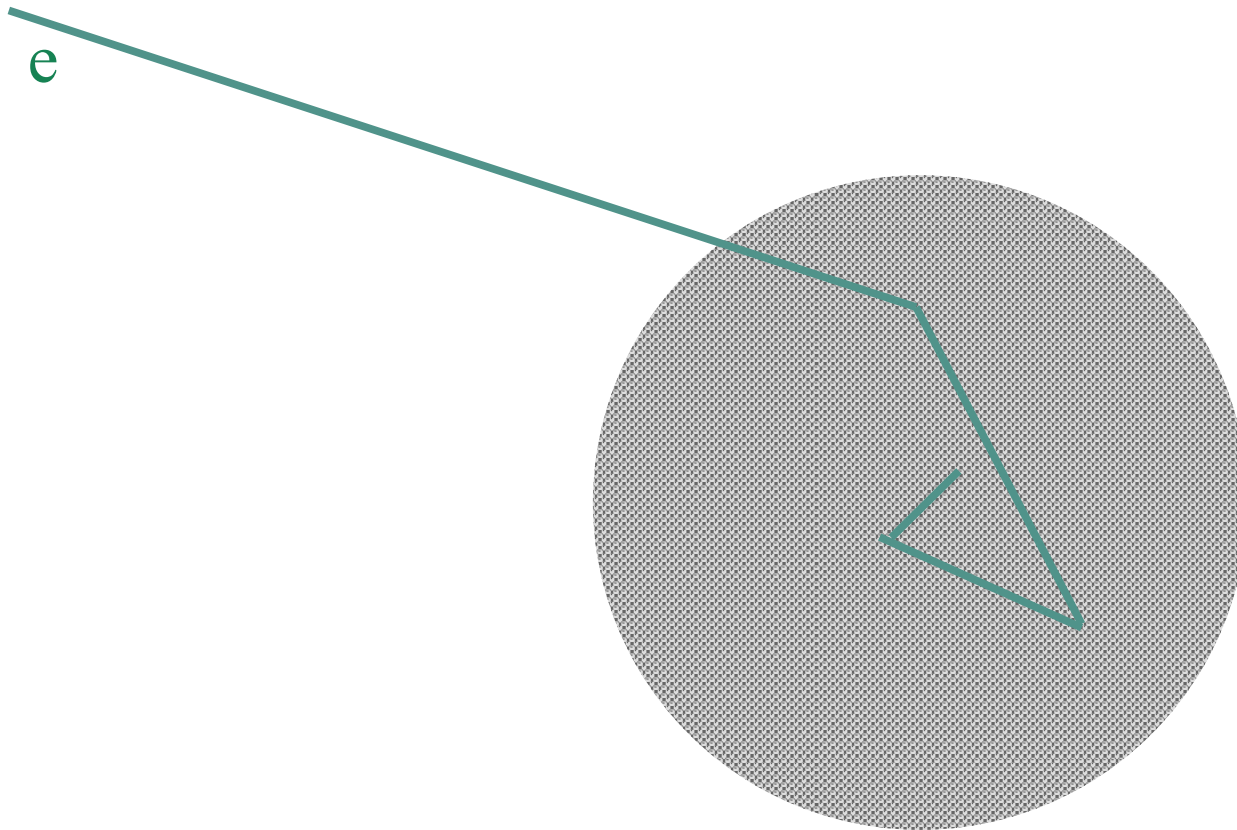


Using DSs to probe scattering with nucleons

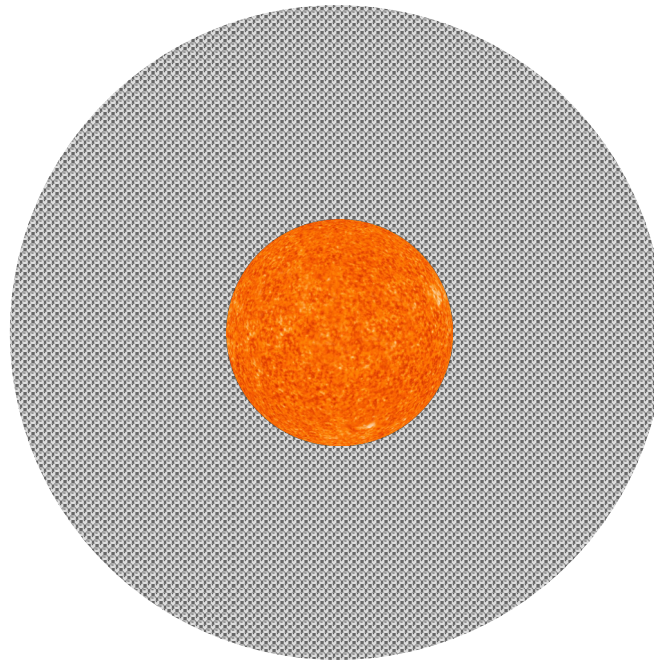


$$C \sim (10^{18} \text{ s}^{-1}) \left[\left(\frac{\sigma}{10^{-45} \text{ cm}^2} \right) \left(\frac{n_p}{10^{-5} \text{ cm}^{-3}} \right) \right] \left(\frac{R_{\text{DS}}}{1 \text{ km}} \right)^3 \left(\frac{M}{M_{\odot}} \right)^{-\frac{3}{2}}$$

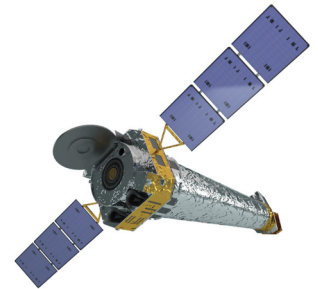
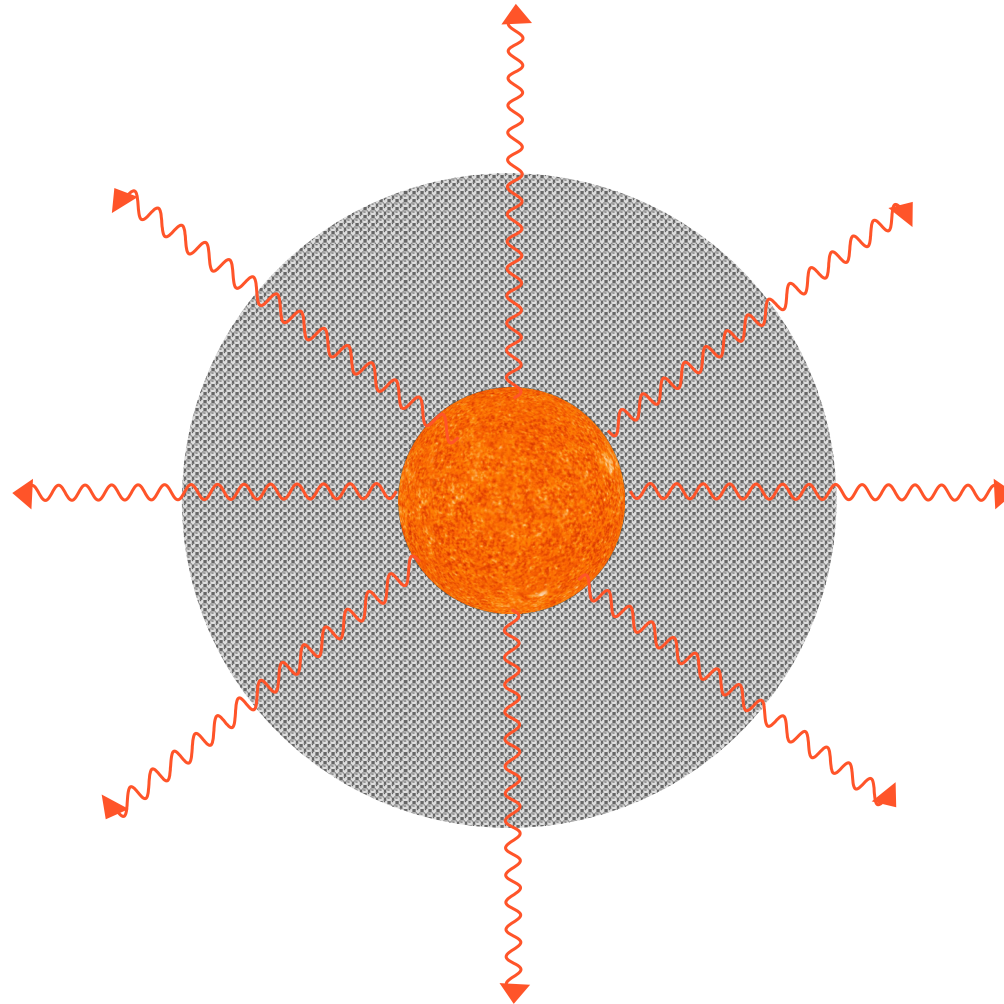
Using DSs to probe scattering with nucleons



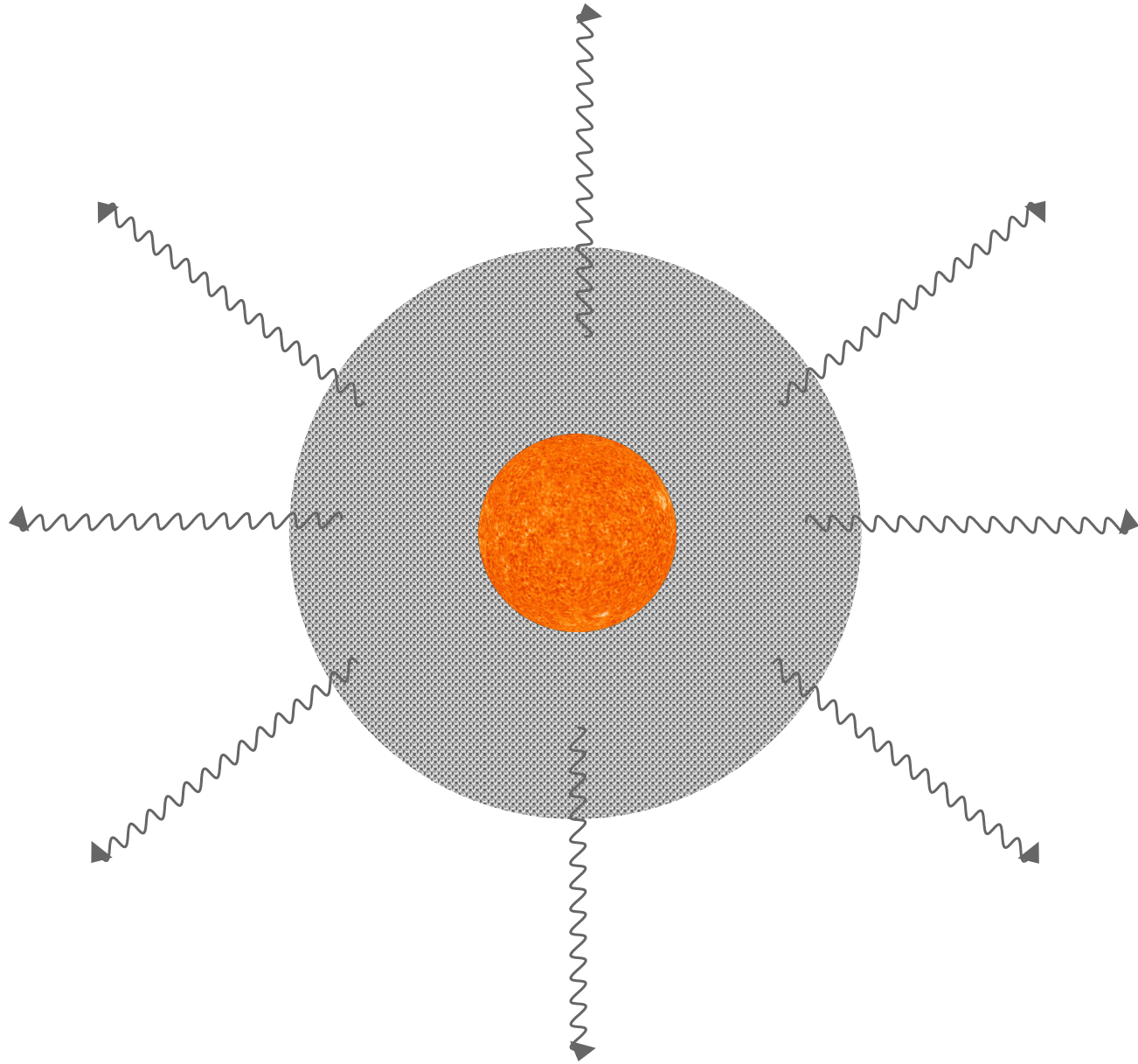
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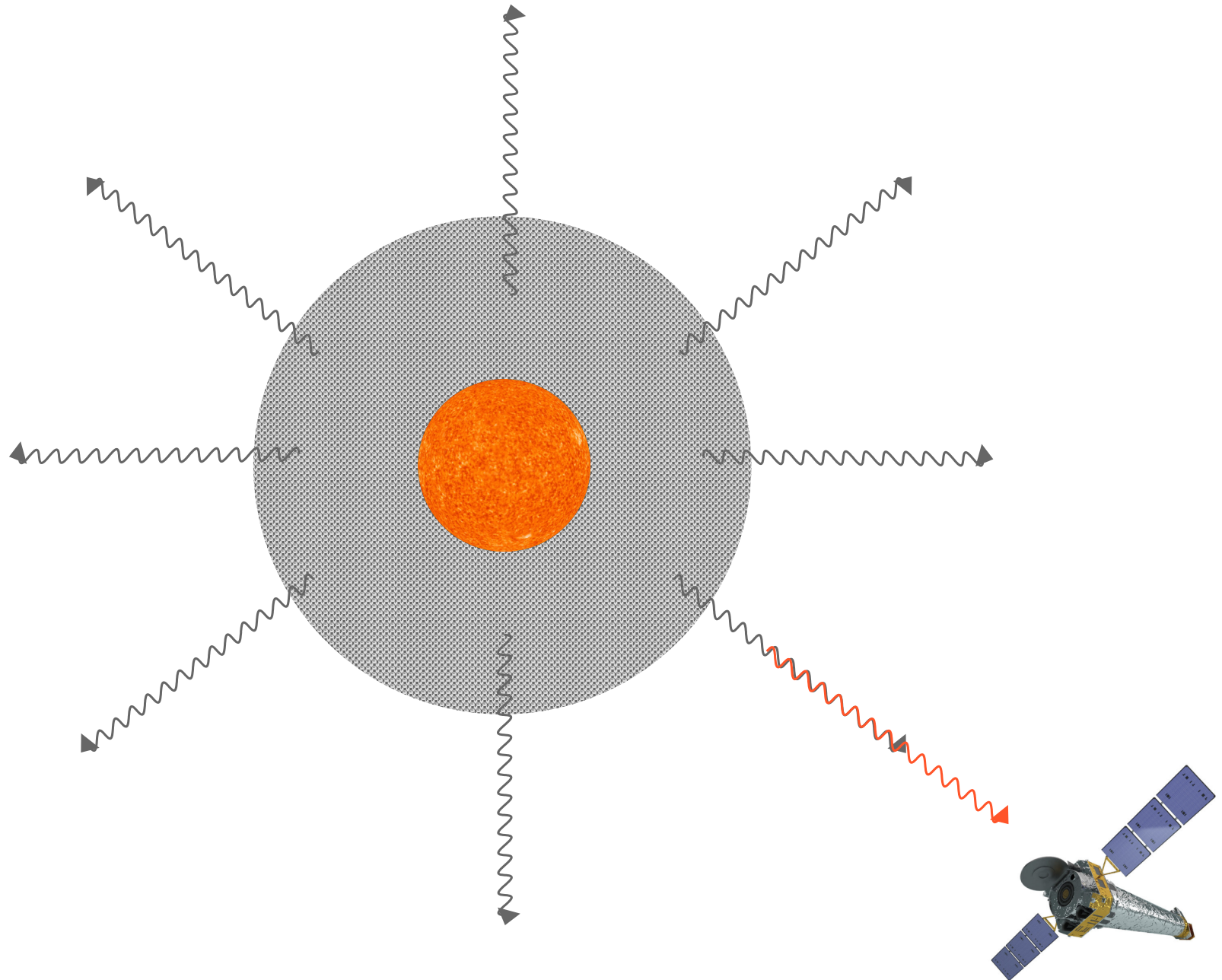
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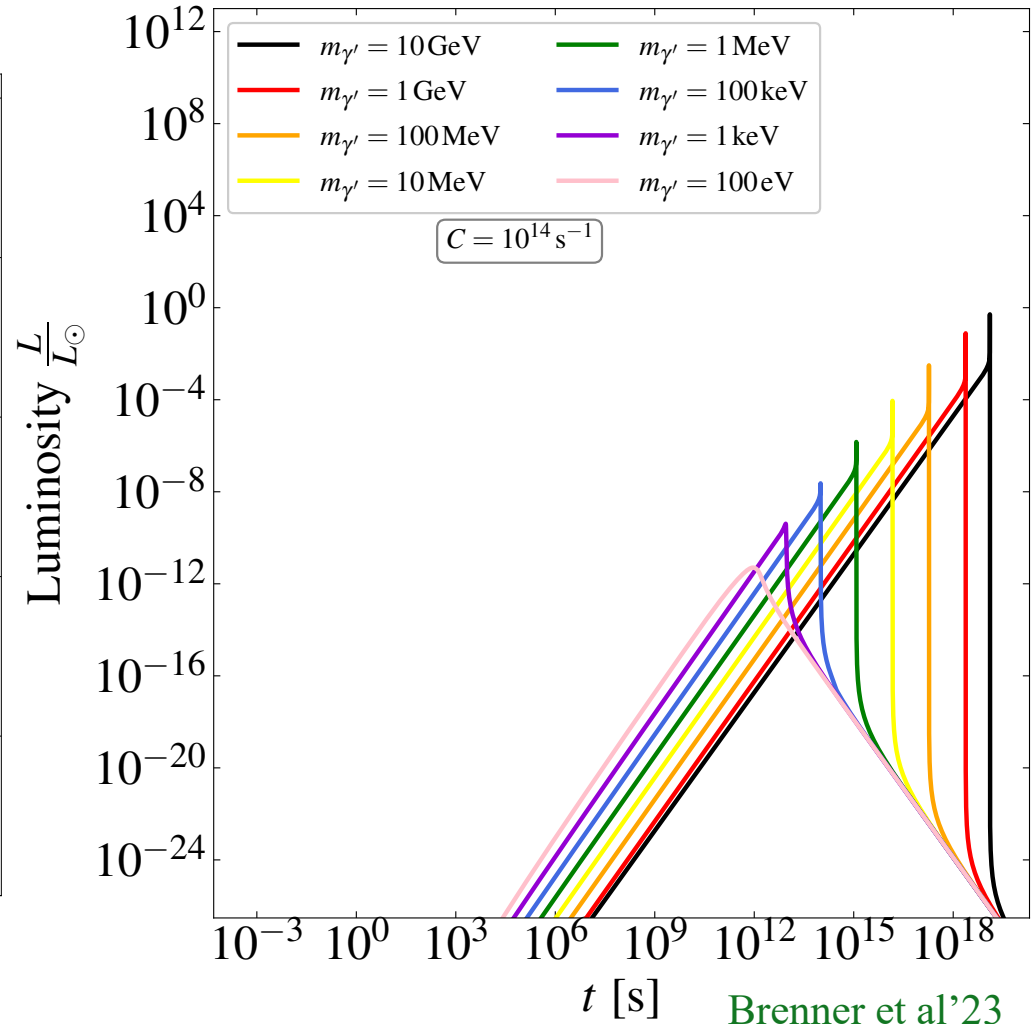
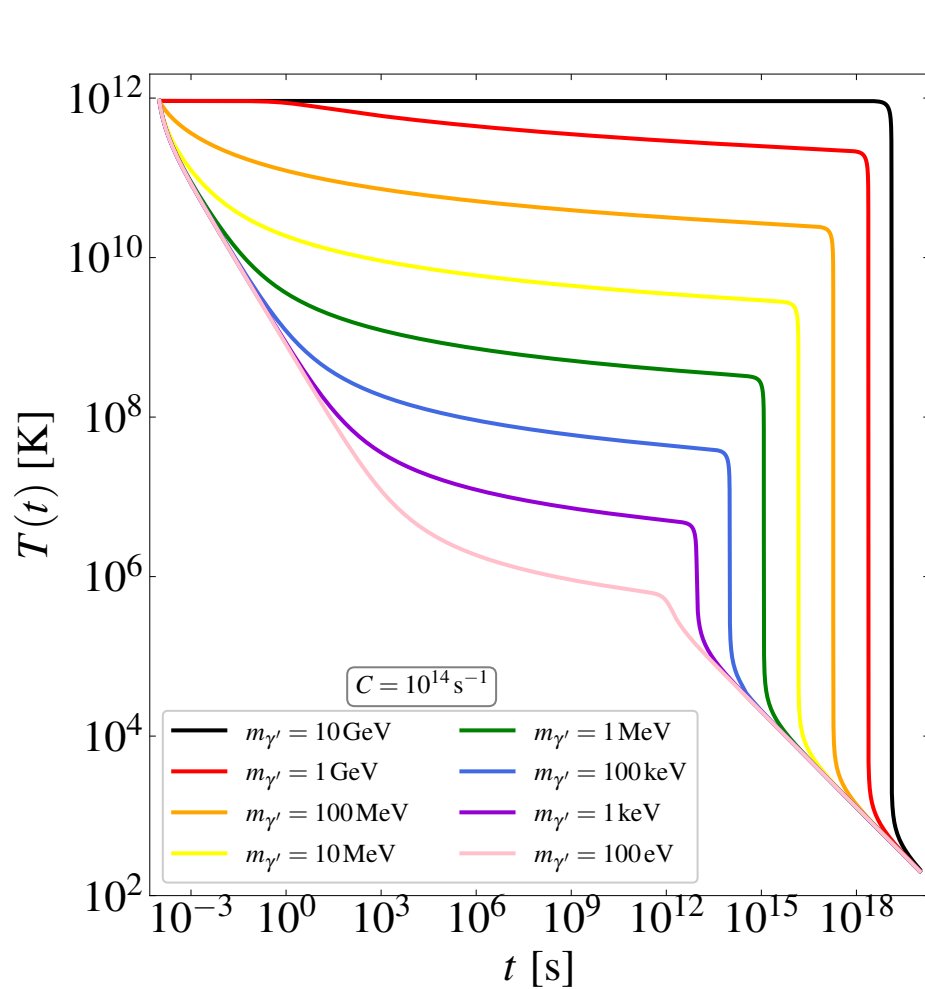
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Using DSs to probe scattering with nucleons



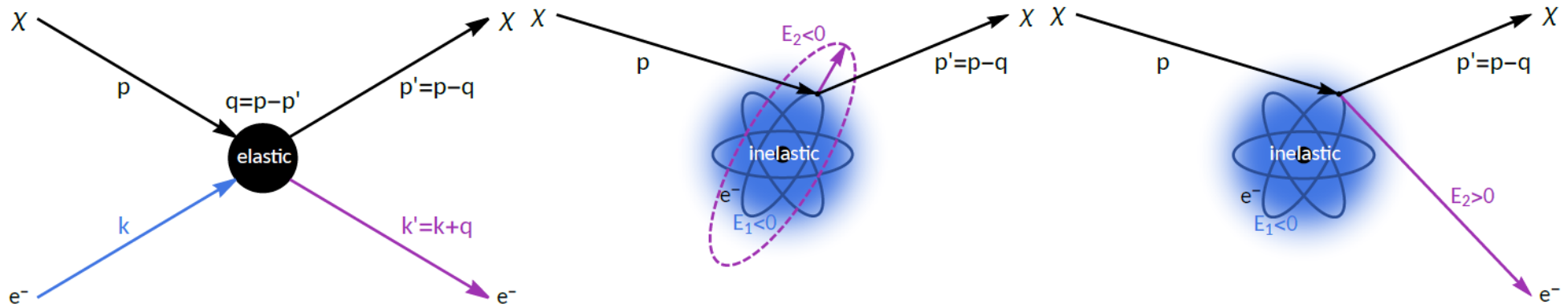
Using DSs to probe scattering with nucleons



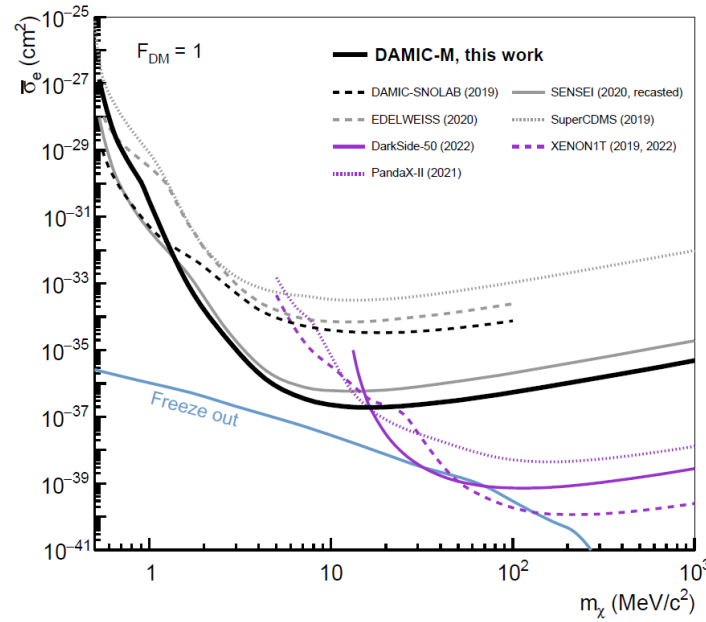
The “dark star” could be very luminous in gamma-rays

If $L=10^{-4} L_{\text{sol}}$, current instruments could detect up to 100 sources, if 1% of the dark matter is in the form of dark stars.

Probing the scattering with electrons



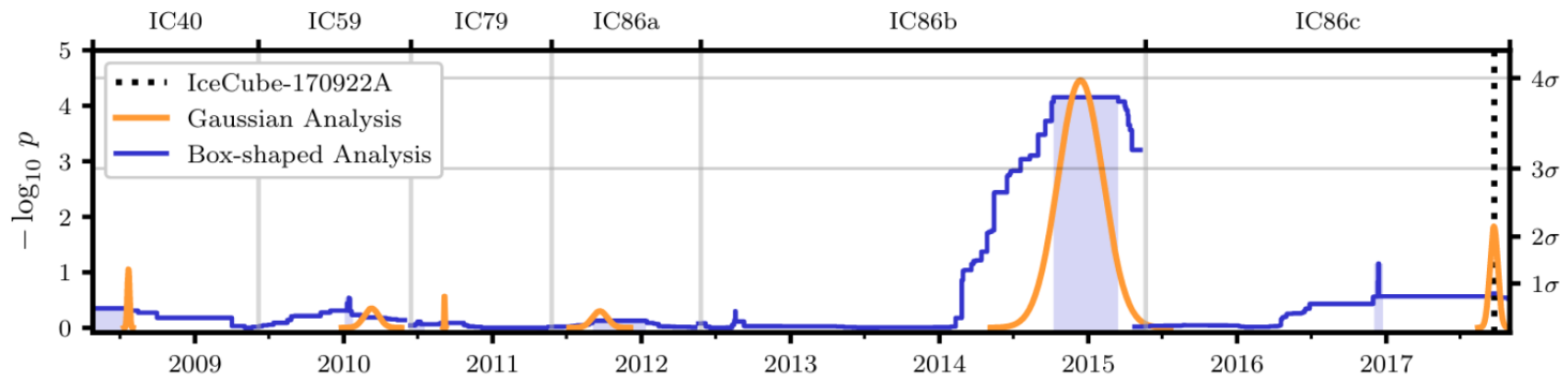
From Catena et al'20



Probing the scattering with neutrinos

The neutrino event IceCube-170922A was coincident in direction and time with a gamma-ray flare from the blazar TXS 0506+056, located 1.75 Gpc away from the Earth..

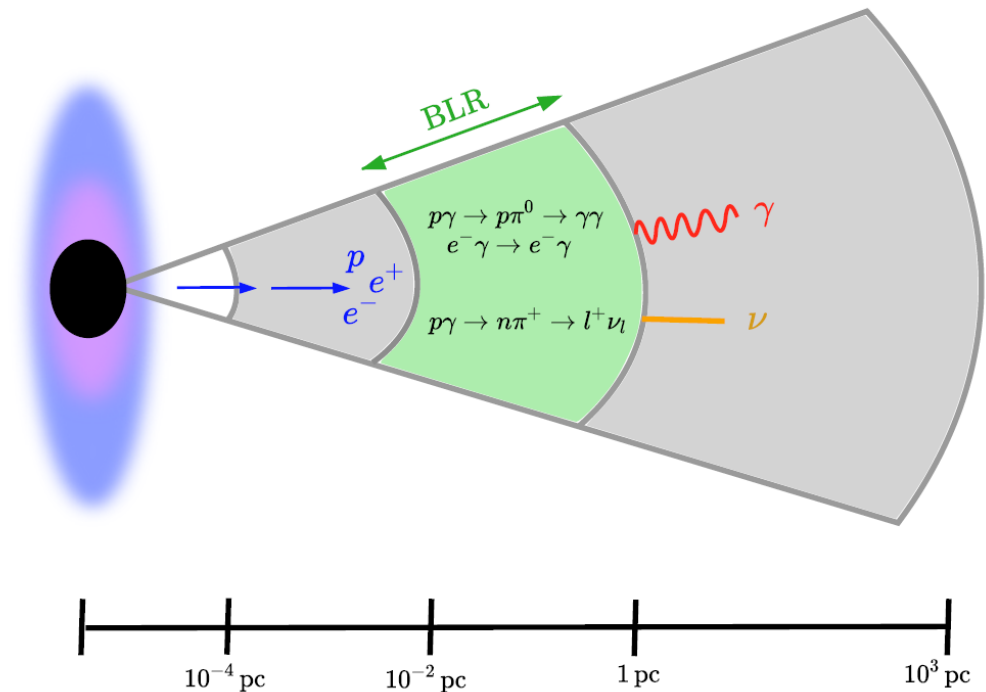
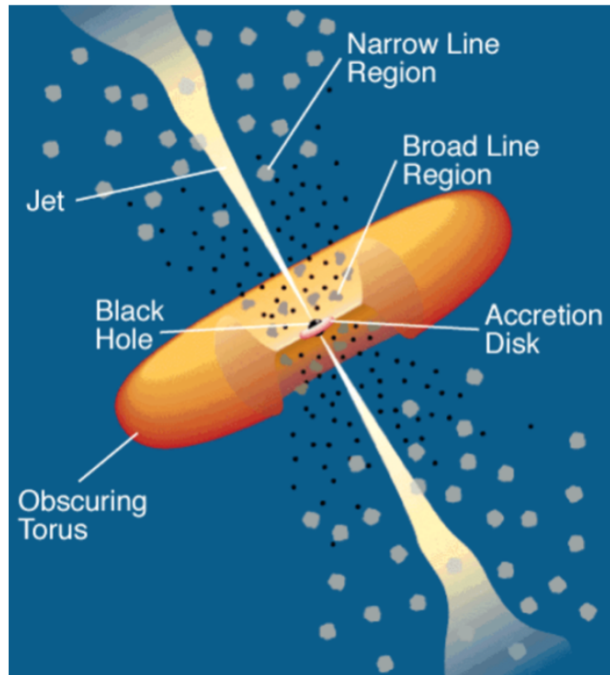
Archival data found 13 ± 5 events coincident with TXS 0506+056.



First known source of high energy astrophysical neutrinos

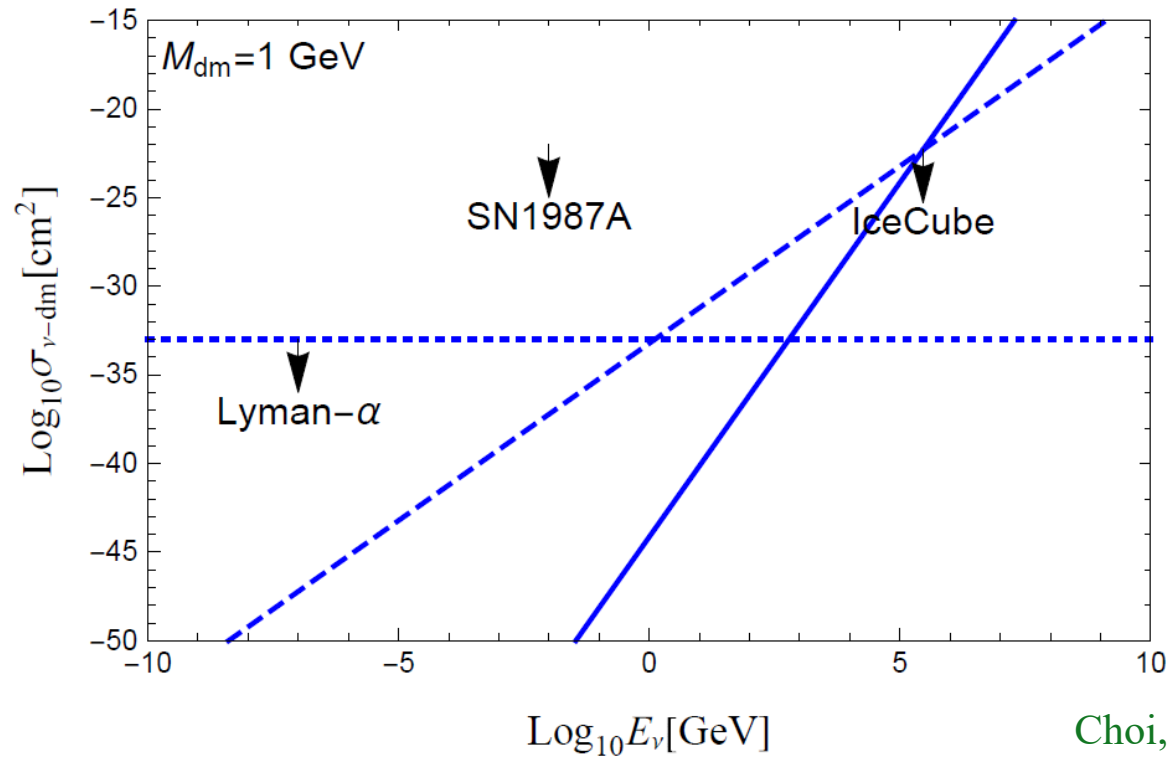
Probing the scattering with neutrinos

The neutrino and photon fluxes can be qualitatively well reproduced in leptohadronic models.



Neutrinos propagate through the intergalactic medium and through the Milky Way before reaching us. If the dark matter neutrino cross-section is large, the neutrino flux would be attenuated.

Probing the scattering with neutrinos



Choi, Kim, Rott '19
Kelly, Machado '19

Probing the scattering with neutrinos

In the center of the blazar it is located a supermassive black hole, with mass $\sim 3 \times 10^8 M_{\text{sun}}$.

The gravity of the black hole produces a “spike” in the dark matter distribution Gondolo, Silk'99, Peebles '72, Quinlan, Hernquist, Sigurdsson '95

$$\rho(r) = \rho_0 \left(\frac{r_0}{r} \right)^\gamma \longrightarrow \rho_{\text{sp}} \sim \rho_R \left(\frac{R_{\text{sp}}}{r} \right)^{\gamma_{\text{sp}}}$$

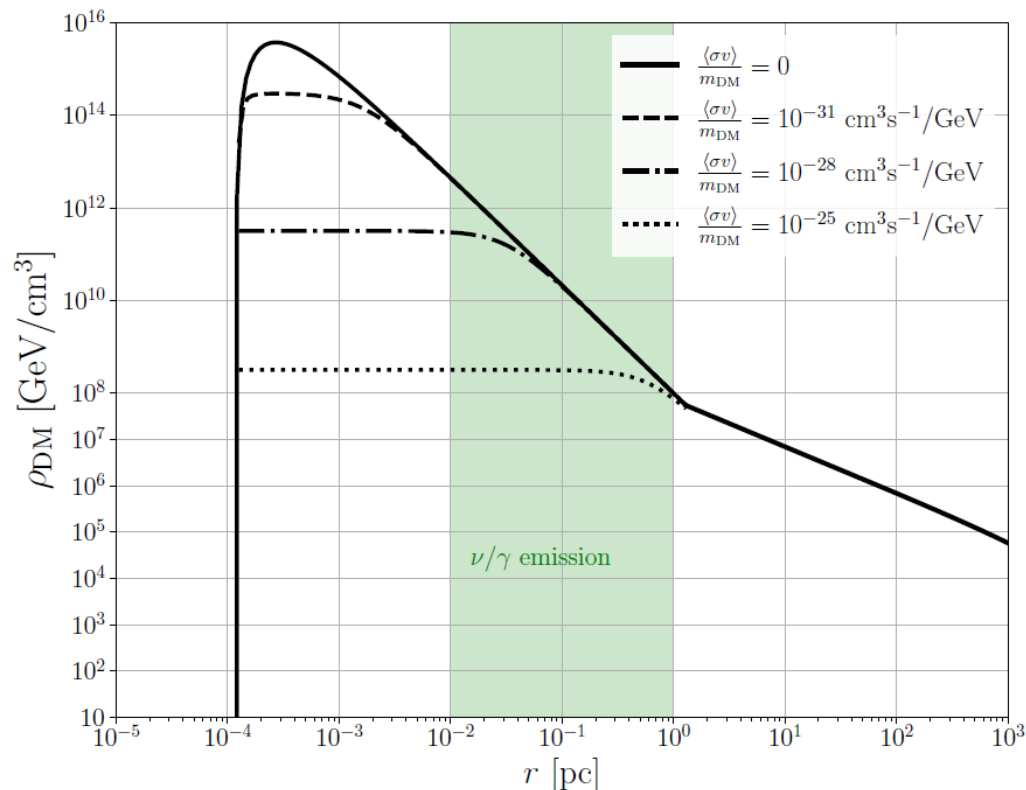
$$\gamma_{\text{sp}} = \frac{9 - 2\gamma}{4 - \gamma}$$

Probing the scattering with neutrinos

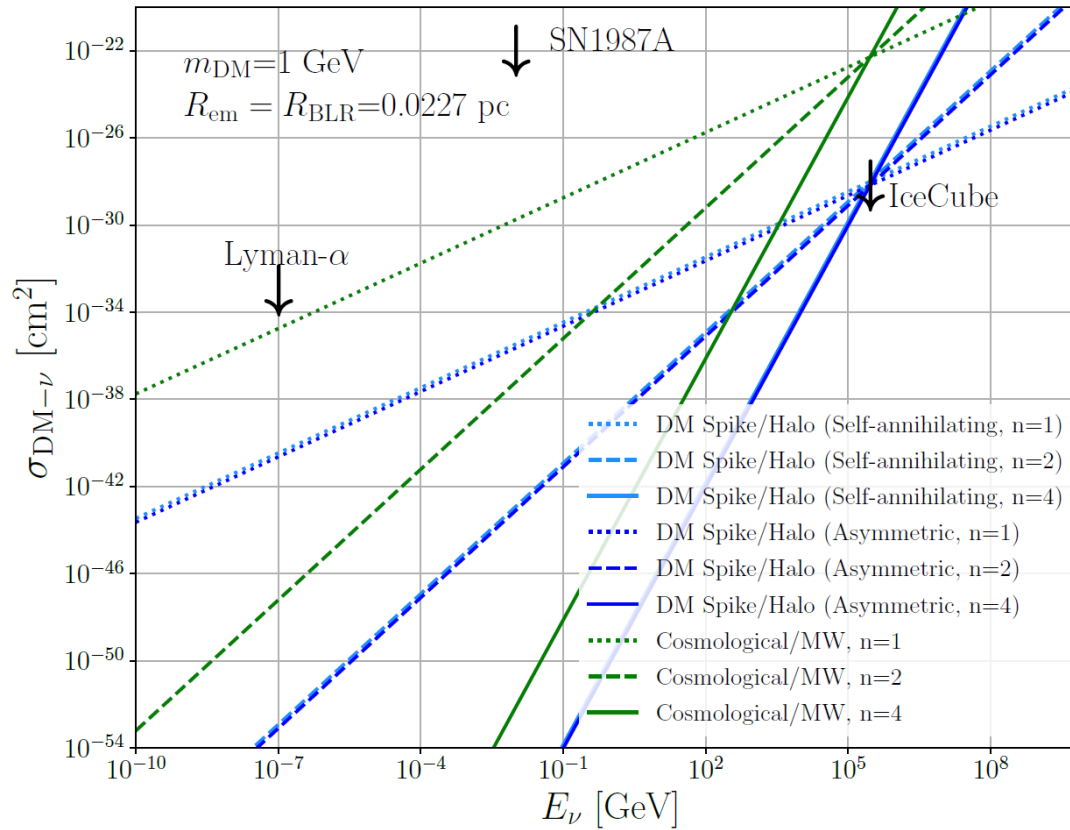
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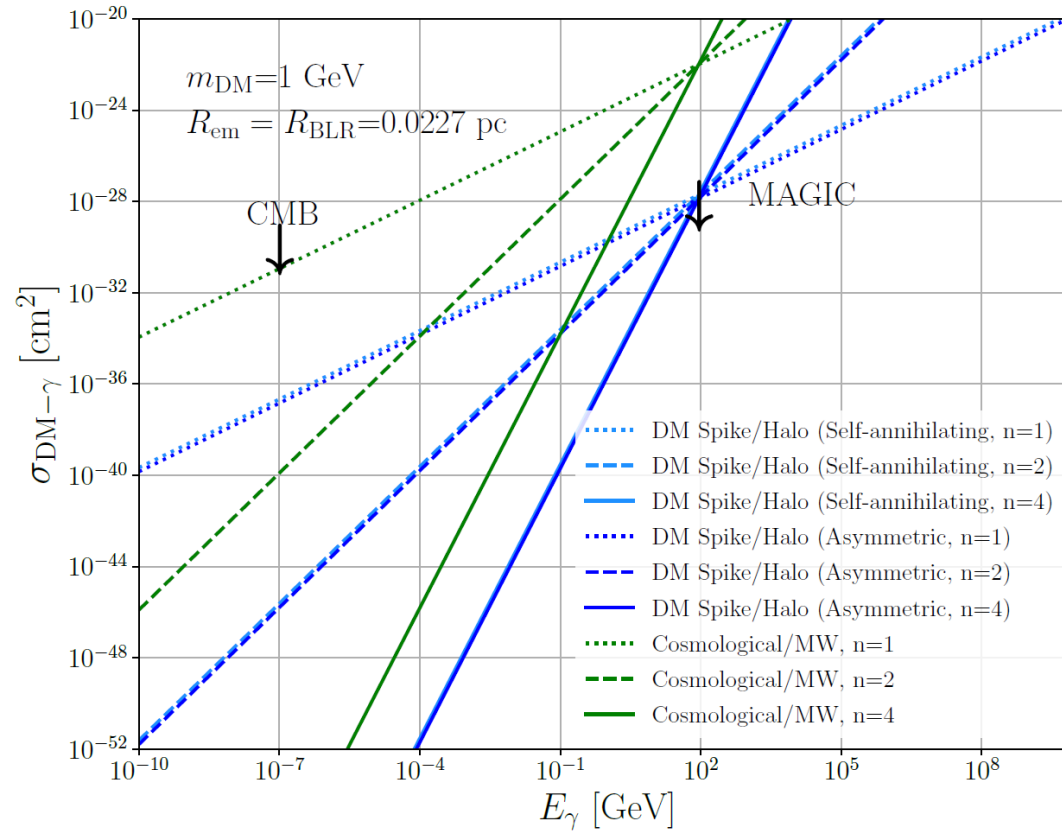
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Probing the scattering with neutrinos



Probing the scattering with photons



Conclusions

- After 40+ years of search, there is still no concluding evidence that dark matter is made of elementary particles.
- “Traditional” searches put some tension on some WIMP scenarios. Many other scenarios still poorly constrained by data.
- Better experiments are needed, but also new ideas for dark matter detection.
- Astronomical objects (compact dark stars, active galactic nuclei, etc.) open new opportunities to detect non-gravitational signals of dark matter.