Dark matter searches ECMS, LHC CERN

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(ON BEHALF OF THE CMS COLLABORATION)



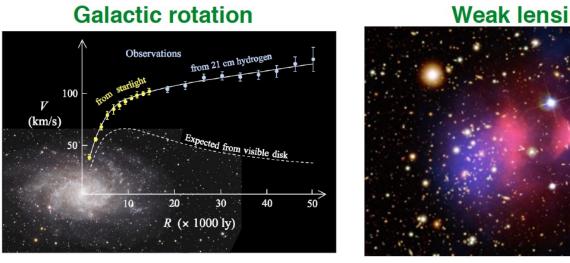
ICHEPAP 2023 11th -15th Dec 2023

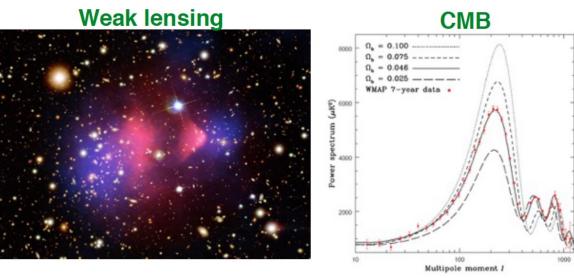


Introduction



Existence of dark matter known through its gravitational interactions





Underlying nature of dark matter (DM) remains unknown

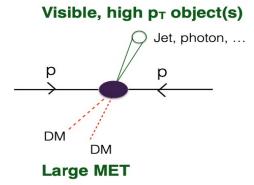
There is a well established case for weakly interacting DM particles (WIMPs)

Such particles may be produced in high energy p-p collisions at LHC!!

Dark Matter

MET

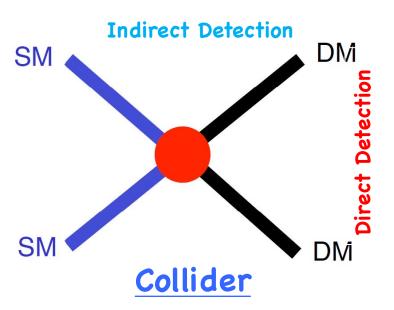
- Favorite collider candidate: WIMP
- Weakly interacting, massive, stable





General collider strategy

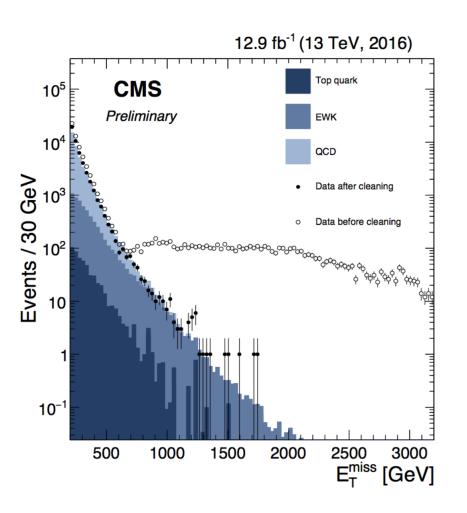
X = 7, W, Z, H, 9, E



Non-interacting particles cause momentum imbalance in the transverse plane of the beam

Challenges with Missing Transverse Momentum





Spurious detector signals can cause fake MET signatures that must be identified and suppressed.

Anomalous high MET can be due to:

- Particles striking sensors in the ECAL photodetectors
- Beam halo particles
- ECAL dead cells (real energy to have been missed)
- Noise in photodiode & readout box electronics in HCAL

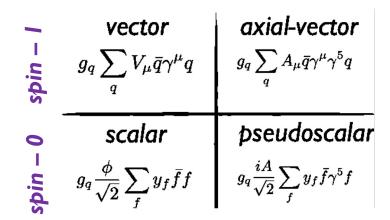


Disclaimer: Presenting sub-set of DM results

Simplified Dark Matter Models

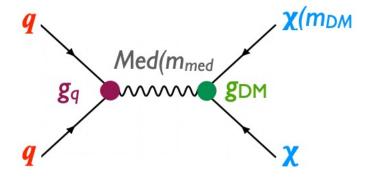


- Dark matter production in pp collisions described using "simplified models"
- Capture the essential features of a variety of DM signals through a minimal set of parameters
- √ s channel mediators



√ t – channel mediators

LHC DM Forum arXiv:1507.00966



Parameters:

- Spin/Parity of mediator
- Mediator mass (M_{med})
- DM Mass (M_{DM})
- Mediator coupling to DM (g_{DM})
- Mediator coupling to quarks (g_q)

CMS Result

MonoJet

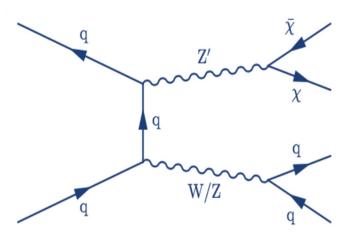


- ✓ Search for physics with particles that decay invisibly in association with a jet
- ✓ Search is performed in MonoJet and MonoV categories and combined

MonoJet: Jets come from the fragmentation of a single quark/gluon

q g_q χ χ q χ χ

MonoV: Jets come from a hadronically decaying V (W/Z) boson



Several new physics models predict such an experimental signature at the detector

MonoJet

CMS Result



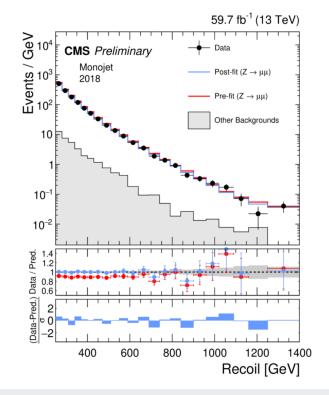
Signal Region: Jets + MET

Background: Z/W+jet, top, dibosons, multijet

• At least one high pT central jet

- Veto events with leptons (e, μ , τ) and photons γ
- MET (Hadronic Recoil) > 250 GeV
- Events are broadly categorized in mono-jet and mono-V based on leading jet pT
- \circ Mono-V: Jet pT (AK8) > 250 GeV
- o **Mono-jet:** Jet pT (AK4) > 100 (150) GeV

We employ semi-data driven technique, supported by statistically independent <u>control</u> regions ($1e/\mu$, $2e/\mu$, t, γ), to constrain the normalization of SM backgrounds



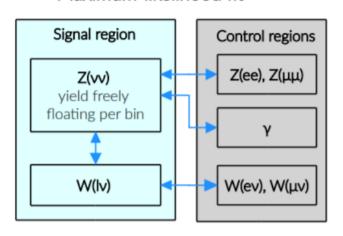
Data/background predictions in Control Regions

MonoJet Signal Extraction



- Signal extraction is based on MET distribution, an unconstrained parameter per recoil bin per category and year.
 - Monojet: 22 bins Mono-V: 7 bins
- SR and CRs are linked together with binned transfer factors (TF).
- The TFs are constrained by simulations, within theory and experimental uncertainties.
 - experimental uncertainties on TFs.
 - mixed QCD-EWK corrections and NNLO QCD on TFs

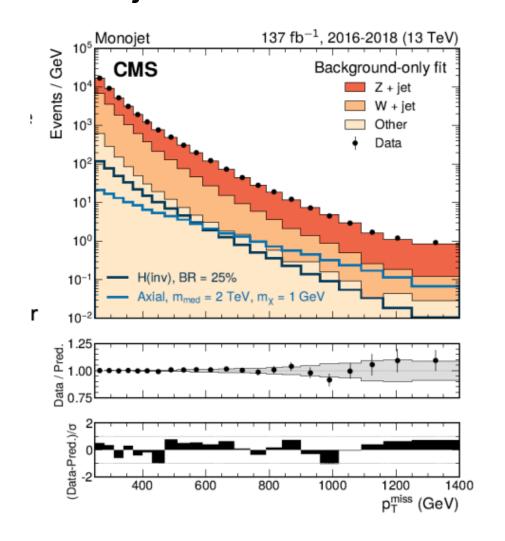
Maximum-likelihood fit.



- Theory uncertainties for the V ratios and V+jets corrections are from the <u>Lindert et al paper</u>.
- This fit model is replicated in each category/year, then likelihoods across the categories/years are combined.

Monojet in SR





Selection

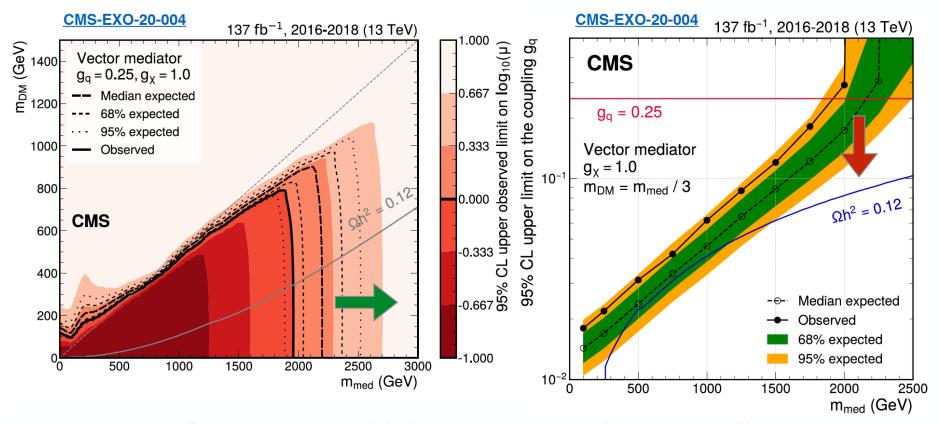
p_T^{miss} > 250 GeV, monojet+ monoV Madgraph generator for major bkgs 22 bins 5 CR (2 W, 2 Z, I photon)

regions split by year

3x22 free floating parameters for monoJet + 3x7 for monoV

MonoJet Results

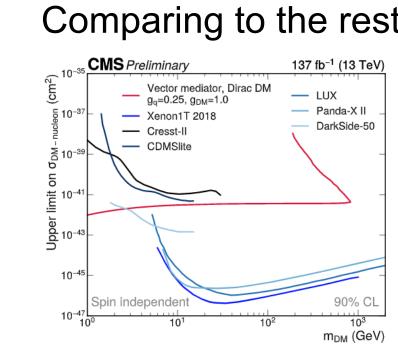


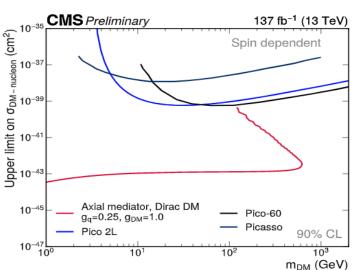


Goal is to probe higher masses and lower couplings

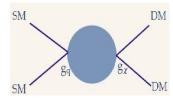
Comparing to the rest of the world



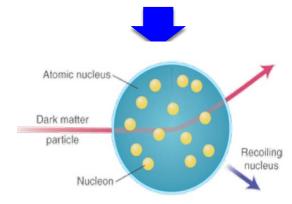




Comparison to direct detection experiments



Axial-vector mediator results are translated into 90% CL exclusion limits on the spin-dependent WIMP-nucleon scattering cross section σ SD as a function of the WIMP mass



ATLAS provides WIMP annihilation rate as a function of WIMP mass [backup]

BHAWNA GOMBER 14th Dec 2023

Mono-Z Search

TOFFICE TO SERVICE TO

Signature: Z(II)+MET

Model for interpretations:

- Simplified model
- ► 2HDM+a

 p_T^{miss} for simplified, m_T for 2HDM+a model

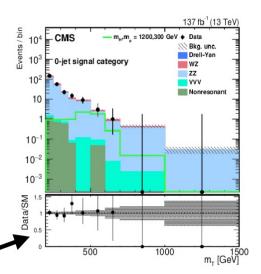
$$m_{\rm T} = \sqrt{2p_{\rm T}^Z p_{\rm T}^{\rm miss} (1 - cos(\Delta \phi_{\ell\ell - \vec{p}_{\rm T}^{\rm miss}}))}$$

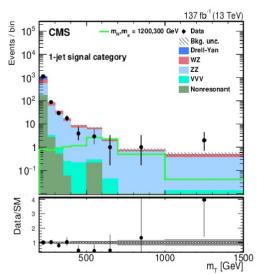
Basic selection:

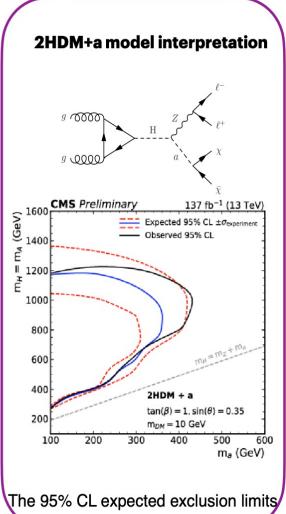
- p_{T(I1)} > 25 GeV, p_{T(I2)} > 20 GeV
- ► |m_{||} m_z| < 15 GeV
- ▶ p_Tmiss > 100 GeV

Backgrounds:

- Drell-Yan, WZ, ZZ, VVV
- Dedicated Control regions to model the background







Mono-Z Comparison

Eur.Phy.J.C.81(2021) 13

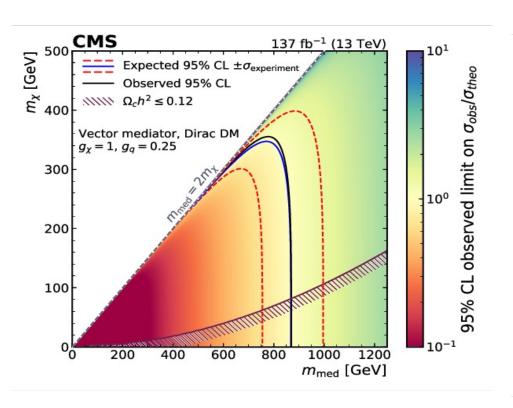


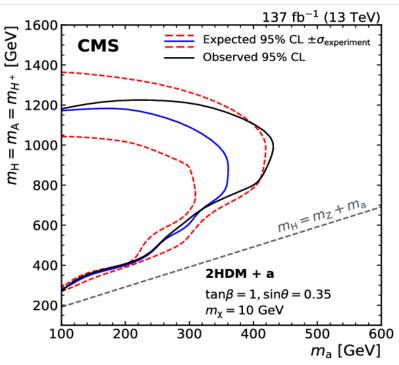
❖ For Simplified DM model (Vector)

Mediator mass excluded upto 800 GeV

For 2HDM+a model

Maximal exclusion $m_a = 350 \text{ GeV}$ and $M_H = 1.2 \text{ TeV}$

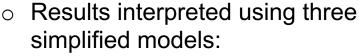




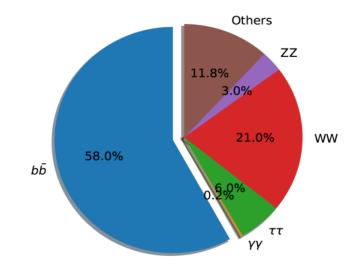
Mono - Higgs

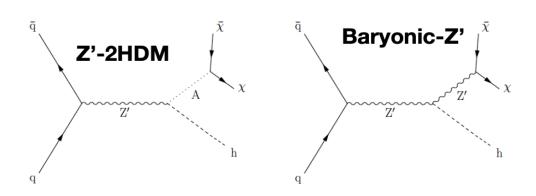


- After the discovery of the Higgs boson (125 GeV) it is possible to probe the dark matter using this new handle
- New massive particle mediates the Higgs-DM interaction
- Search performed in five decay channels and statistically combined
 - \circ $b\overline{b}$, $\gamma\gamma$, WW, ZZ, and $\tau\tau$



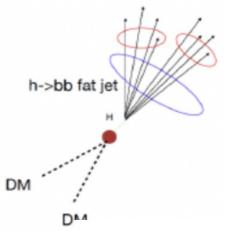
- Z' 2HDM
- Baryonic Z'
- o 2HDM + a





Mono-Higgs

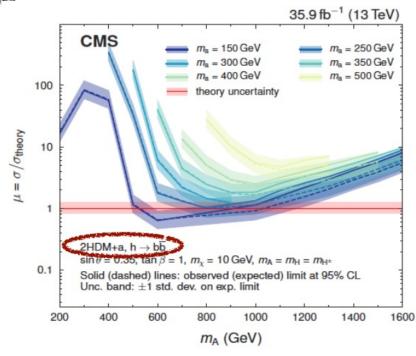
Final states orthogonal to each other

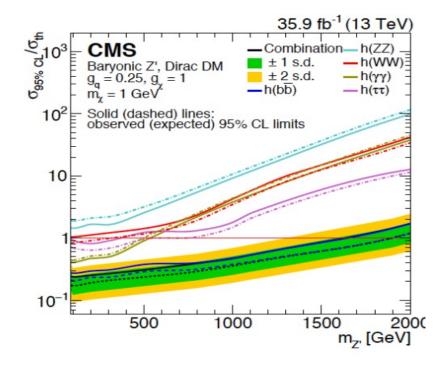


$H \rightarrow bb$	Most sensitive
2HDM+a	CA15 jets
Baryonic – Z'	CA15 jets
Z' – 2HDM	AK8 jets

Decay channel	Final state or category
$h \rightarrow bb$	AK8 jet (Z'-2HDM) CA15 jet (Baryonic Z')
$h\to\gamma\gamma$	$p_{\mathrm{T}}^{\mathrm{miss}} \in 50130\mathrm{GeV}$ $p_{\mathrm{T}}^{\mathrm{miss}} > 130\mathrm{GeV}$
$h \to \tau \tau$	$egin{array}{l} au_{ m h} au_{ m h} \ ext{e} au_{ m h} \end{array}$
$h \to WW \\$	eνμν
$h \to ZZ$	4e 4μ 2e2μ

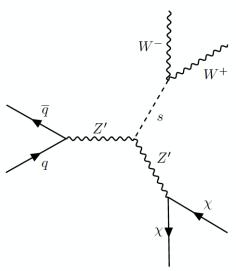
JHEP 03 (2020) 025





Dark Higgs (WW)





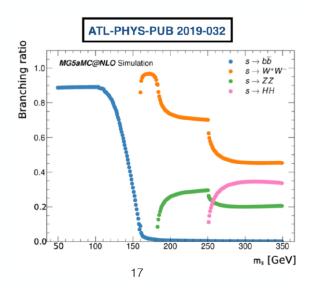
Model parameters are:

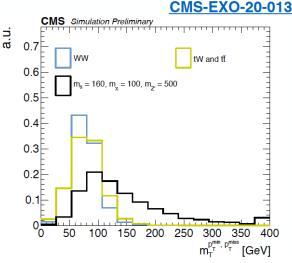
- DM mass: mχ,
- Z' mass: mZ',
- dark Higgs mass: ms,
- Z' couplings to quarks (gq)
- Z' couplings to DM (gx),
- the mixing angle between SM and the dark Higgs bosons (sin θ).

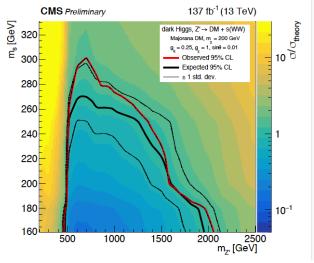
Dark Higgs (WW)

DM particles acquire their mass through their interactions with a Dark Higgs boson.

Signal extraction: 3D ML fit to $\Delta R(II)$, mII, mT(Imin + pT)







Higgs Invisible



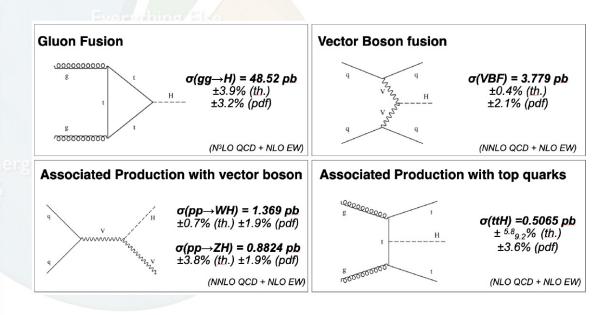
Strong experimental evidence for DM from astrophysical observations. Most studied class of theories predict DM to be a weakly interacting massive particle.

Higgs Portal

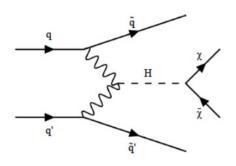
Known Higgs decay branching fractions allow decays to invisible particles ~20%

Produce the Higgs

All Higgs production modes can be studied to tag the event

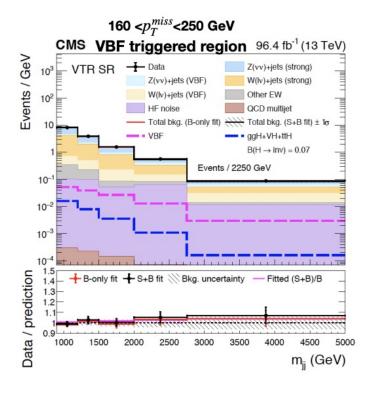


VBF Higgs Invisible



 p_T^{miss} >250 GeV CMS MET triggered region 101 fb⁻¹ (13 TeV) Events / GeV MTR SR Z(vv)+jets (VBF) W(lv)+jets (strong) W(lv)+jets (VBF) QCD multijet ggH+VH+ttH $B(H \rightarrow inv) = 0.07$ Events / 1500 GeV 10^{-2} 10^{-3} Data / prediction B-only fit + S+B fit Bkg. uncertainty 1500 2500 3000 3500 mii (GeV) New trigger strategy: using jet properties from VBF production in addition to p_T^{miss} trigger.

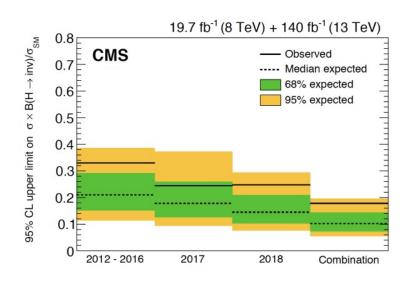
Using V+jets and $\gamma + jets$ CRs to constrain major backgrounds (Z(vv)+jet and W(lv)+jets).

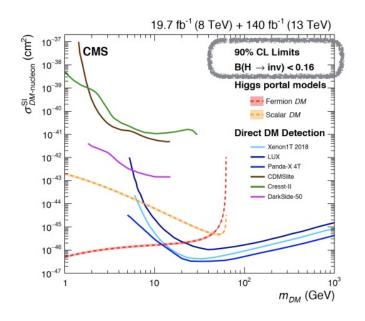


VBF Higgs Invisible



- Combination of Run 1 and Run2
 - 95% CL upper limit on the in BR (H→invisible) < 0.18 (0.10)





Constraints are compatible with SM H → invisible branching ratio.

Constraints on spin independent DM-neucleon cross-section

Dark photon in VBF Higgs

CMS-EXO-20-005



Search for dark photon in VBF Higgs boson events.

F Higgs boson events.

Year

Triggering object

2016

VBF photon

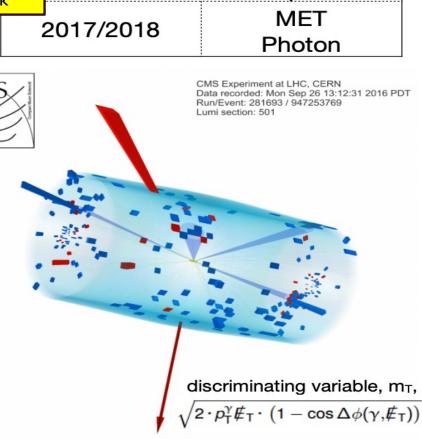
MET

Photon

h(125 GeV) is of particular interest but higher masses are also considered.

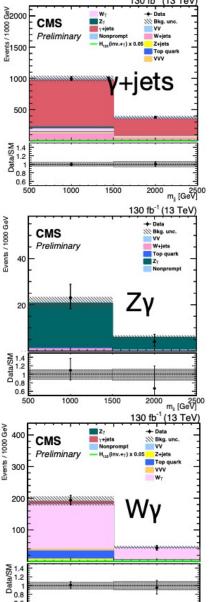
Dedicated CRs for major background: **W+jets, Wγ, Zγ, γ+jets**

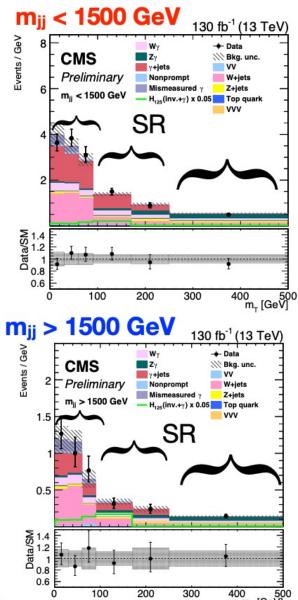
Simultaneous fit of SR and CRs



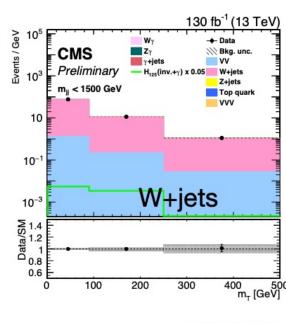
mjj

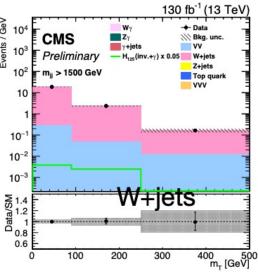






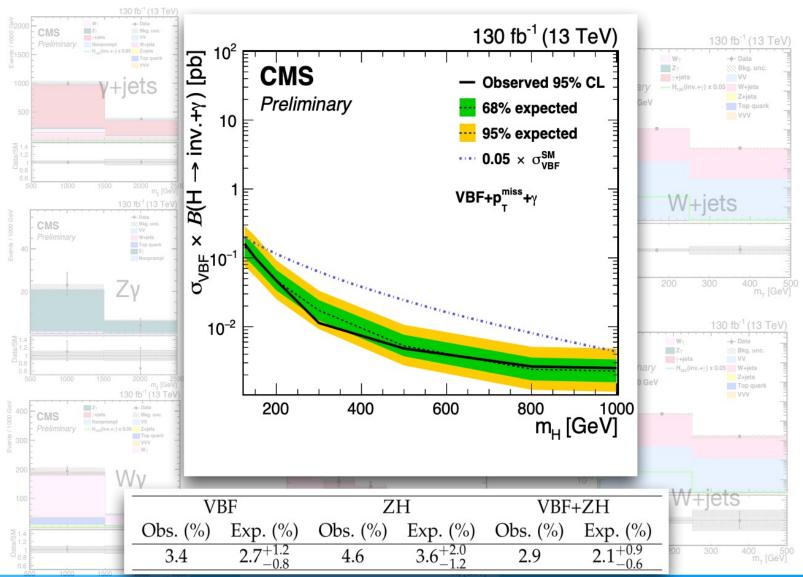






Dark photon in VBF Higgs





Summary



- **❖** Performing a variety of searches for new phenomena at the LHC, including searches for dark matter, which provide access to the phase space
- ❖ Presented a few new results for CMS, all of which use the full Run2 results
- ❖ No signal observed yet, but more to do! Need to look everywhere
- More exciting Run2 results will be coming out in the comign months
 Stay tuned for this and the upcoming Run3!

Thank You

Backup



MonoZ: Fitting Strategy



Fitting Strategy

- ATLAS:
 - nonresonant and WZ production normalized from data
 - ► ZZ production **not** normalized from data, replying on simulation post-facto
- ► CMS:
 - nonresonant production normalized from data
 - ► Z + jets events in 0 jet and 1 jet categories normalized from data
 - ► WZ and ZZ production estimated from data using a single normalization factor
 - ► large EWK correction uncertainties considered
 - ▶ VV shape: 3 additional nuisances ($\pm 10/20/30\%$) at low ($80 < p_{\rm T}^{\rm miss} < 200 {\rm GeV}$), medium ($200 < p_{\rm T}^{\rm miss} < 400 {\rm GeV}$), and high $p_{\rm T}^{\rm miss}$ ($p_{\rm T}^{\rm miss} > 400 {\rm GeV}$) to allow for independent leverage in the fit
- ► Arguable, the last two set of uncertainties are a matter of choice, having both of them is a rather conservative approach

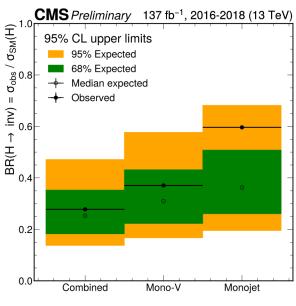
OTHER INTREPRETATIONS : CMS MONOJET



CMS PAS EXO-20-004

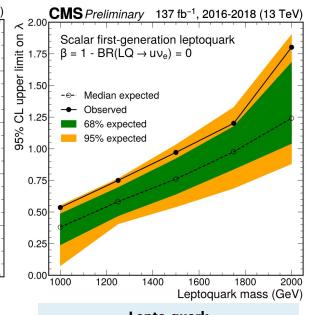
ATLAS: arXiv:2102.10874

NEW



Higgs Portal

 $B(H\rightarrow inv) < 27.8 (25.3)$

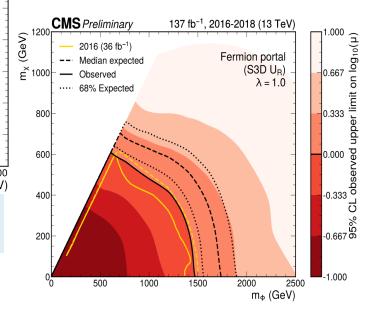


Lepto-quark

Pair production: dominates below λ < 1 TeV while single above 1 TeV

Fermion Portal: DM t-channel

Mediater mass below ~1.5 TeV excluded for $m_{DM} = 1 \text{ TeV}$

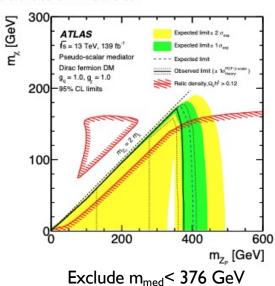


Monojet comparison with ATLAS

ATLAS-EXOT-2018-06 CMS-EXO-20-004







137 fb⁻¹, 2016-2018 (13 TeV)

---- Median expected

---- Gbserved

---- 68% expected

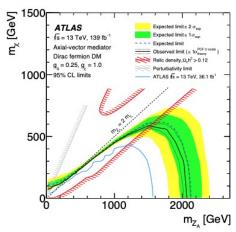
95% expected

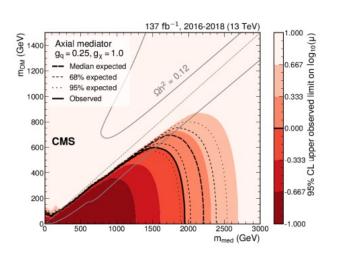
Pseudoscalar mediator $pp \rightarrow \chi \chi + j / Z(qq)$ $m_{DM} = 1 \text{ GeV}, g_q = 1.0, g_\chi = 1.0$ m_{med} (GeV)

CMS has significantly better limits in pseudo-scalar mass exclusion

Exclude m_{med} < 470 GeV (490 GeV expected).







CMS and ATLAS pretty much similar limits for spin-1, exclude mediator mass upto 1.95 (2.1) TeV, for CMS(ATLAS), respectively

- CMS produces exclusion in coupling which ATLAS doesn't

Monophoton ATLAS

TOF HOUSE

Table 5: Observed and expected yields from SM backgrounds obtained from the 'simplified shape fit' described in Section 5. The normalisation factors obtained from the fit are also shown. The uncertainty includes both the statistical and systematic uncertainties. The individual uncertainties can be correlated and do not necessarily add in quadrature to equal the total background uncertainty.

	SRE1	SRE2	SRE3	SRI4
Observed events	3023	1164	679	427
Expected SM events	3070 ± 130	1182 ± 75	680 ± 53	448 ± 42
$Z(\rightarrow \nu\nu)\gamma$	1910 ± 110	758 ± 65	468 ± 49	306 ± 40
$W(o\ell u)\gamma$	394 ± 22	159 ± 15	71.0 ± 8.2	56.7 ± 7.1
$Z(o\ell\ell)\gamma$	33.2 ± 2.4	9.32 ± 0.89	4.26 ± 0.48	2.69 ± 0.37
γ + jets	87 ± 35	11.9 ± 4.8	2.7 ± 1.1	3.0 ± 1.2
Fake photons from e	511 ± 48	188 ± 18	100.9 ± 9.5	59.7 ± 5.6
Fake photons from jets	136 ± 28	56 ± 29	33 ± 16	20 ± 11
$k_{Z\gamma}$	0.99 ± 0.08	0.89 ± 0.09	0.90 ± 0.11	0.86 ± 0.12
$k_{W\gamma}$	0.81 ± 0.09	0.84 ± 0.11	0.74 ± 0.11	0.85 ± 0.13
$k_{\gamma+ m jets}$	0.82 ± 0.21			

Table 6: Summary of the uncertainties (%) in the background estimate for inclusive SRs after the background-only fit and for exclusive SRs after the 'simplified shape fit'. The individual uncertainties can be correlated and do not necessarily add in quadrature to equal the total background uncertainty.

SRI1	SRI2	SRI3	SRI4	SRE1	SRE2	SRE3
[%]	[%]	[%]	[%]	[%]	[%]	[%]
3.5	4.8	6.2	9.5	4.3	6.3	7.8
2.4	3.6	5.3	8.5	3.3	5.0	6.7
1.4	2.5	2.8	4.1	1.4	3.6	3.7
1.6	2.2	2.5	2.7	2.2	2.2	2.3
2.1	2.0	2.0	2.0	2.3	2.3	2.1
1.0	1.2	1.3	1.4	1.0	1.0	1.2
0.8	0.6	0.7	0.9	0.9	0.9	0.6
0.7	0.8	0.9	1.0	0.6	0.7	0.9
0.1	0.4	0.7	0.9	0.5	0.2	0.5
0.5	0.5	0.5	0.5	0.5	0.5	0.5
0.3	0.2	0.2	0.2	0.3	0.2	0.2
	[%] 3.5 2.4 1.4 1.6 2.1 1.0 0.8 0.7 0.1 0.5	[%] [%] 3.5 4.8 2.4 3.6 1.4 2.5 1.6 2.2 2.1 2.0 1.0 1.2 0.8 0.6 0.7 0.8 0.1 0.4 0.5 0.5	[%] [%] [%] 3.5 4.8 6.2 2.4 3.6 5.3 1.4 2.5 2.8 1.6 2.2 2.5 2.1 2.0 2.0 1.0 1.2 1.3 0.8 0.6 0.7 0.7 0.8 0.9 0.1 0.4 0.7 0.5 0.5 0.5	[%] [%] [%] [%] 3.5 4.8 6.2 9.5 2.4 3.6 5.3 8.5 1.4 2.5 2.8 4.1 1.6 2.2 2.5 2.7 2.1 2.0 2.0 2.0 1.0 1.2 1.3 1.4 0.8 0.6 0.7 0.9 0.7 0.8 0.9 1.0 0.1 0.4 0.7 0.9 0.5 0.5 0.5 0.5	[%] [%] [%] [%] 3.5 4.8 6.2 9.5 4.3 2.4 3.6 5.3 8.5 3.3 1.4 2.5 2.8 4.1 1.4 1.6 2.2 2.5 2.7 2.2 2.1 2.0 2.0 2.0 2.3 1.0 1.2 1.3 1.4 1.0 0.8 0.6 0.7 0.9 0.9 0.7 0.8 0.9 1.0 0.6 0.1 0.4 0.7 0.9 0.5 0.5 0.5 0.5 0.5 0.5	[%] [%] [%] [%] [%] [%] 3.5 4.8 6.2 9.5 4.3 6.3 2.4 3.6 5.3 8.5 3.3 5.0 1.4 2.5 2.8 4.1 1.4 3.6 1.6 2.2 2.5 2.7 2.2 2.2 2.1 2.0 2.0 2.0 2.3 2.3 1.0 1.2 1.3 1.4 1.0 1.0 0.8 0.6 0.7 0.9 0.9 0.9 0.7 0.8 0.9 1.0 0.6 0.7 0.1 0.4 0.7 0.9 0.5 0.2 0.5 0.5 0.5 0.5 0.5 0.5

Monophoton ATLAS



Table 5: Observed and expected yields from SM backgrounds obtained from the 'simplified shape fit' described in Section 5. The normalisation factors obtained from the fit are also shown. The uncertainty includes both the statistical and systematic uncertainties. The individual uncertainties can be correlated and do not necessarily add in quadrature to equal the total background uncertainty.

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Observed events	3023	1164	679	427
Expected SM events	3070 ± 130	1182 ± 75	680 ± 53	448 ± 42
$Z(\rightarrow \nu \nu)\gamma$	1910 ± 110	758 ± 65	468 ± 49	306 ± 40
$W(o\ell u)\gamma$	394 ± 22	159 ± 15	71.0 ± 8.2	56.7 ± 7.1
$Z(o\ell\ell)\gamma$	33.2 ± 2.4	9.32 ± 0.89	4.26 ± 0.48	2.69 ± 0.37
γ + jets	87 ± 35	11.9 ± 4.8	2.7 ± 1.1	3.0 ± 1.2
Fake photons from e	511 ± 48	188 ± 18	100.9 ± 9.5	59.7 ± 5.6
Fake photons from jets	136 ± 28	56 ± 29	33 ± 16	20 ± 11
$k_{Z\gamma}$	0.99 ± 0.08	0.89 ± 0.09	0.90 ± 0.11	0.86 ± 0.12
$k_{W\gamma}$	0.81 ± 0.09	0.84 ± 0.11	0.74 ± 0.11	0.85 ± 0.13
$k_{\gamma+\mathrm{jets}}$	0.82 ± 0.21			

Towards Run3



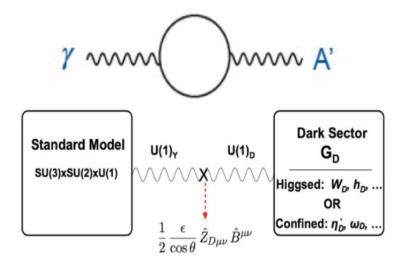
- ❖ MonoZ
 - ❖ No major hurdles, uses standard objects/triggers
 - ❖ Improve ZZ control region (had low statistics in Run2)
 - ❖ Use aggressive approach for theory uncertainty for ZZ
 - ❖ Expand 2HDM+a results, include sin(theta) scan, as done by ATLAS
- **Monojet**
 - ❖ In Run2, electron/photon uncertainties are the leading limitations
 - ❖ Need to do in-house developments for the scale factors
 - ❖ Use of most recent MadGraph which includes NLO QCD & EW corrections, should help with simulation of V+jets background and reduce uncertainties
 - ❖ Better MET trigger turn-on could improve low-mass constraints
 - ❖ For MonoV category, develop trigger which includes boosted tagging
 - * Reinterpretation: Publish full likelihood instead of simplified as ATLAS did for Run2 paper
- Monophoton
 - ❖ Include more intrepretation, ALP could be one
 - ❖ Use new madgraph samples which can include NLO QCD & EW corrections
- * Hinv
 - ❖ VBF
 - ❖ Recovery of the HF jets -- currently detector issues cost ~30% of acceptance
 - ❖ Better trigger based on VBF + MET @ L1 / HLT, including angular cuts to reject QCD online
 - Theoretical constraints for V+jets background
 - ❖ Not available currently, work with theorists needed
 - **ttH**
 - * Re-optimized AK8 taggers for Run-3
 - ❖ Use ML to separate ttH (H->inv) from ttbar

Dark Photon



Add U(1)_d from hidden sector

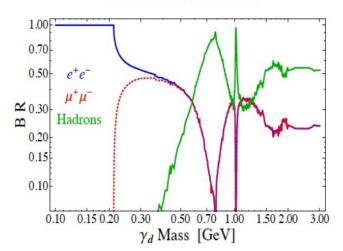
- Connection between dark sector and SM
- Couple with SM via kinetic mixing, ε is kinetic mixing coefficient
- Massive gauge boson $(A/Z_d/\gamma_d)$
- ϵ and mass of $(A/Z_d/\gamma_d)$ are key parameters



Search Strategy based on life:

- Small: Prompt, resolved/collimated decay e.g. LJ
- Medium: Resolved/collimated decay e.g. delay LJ, displaced muons
- Long: stable particles, MET signature at colliders

From arXiv:1002.2952



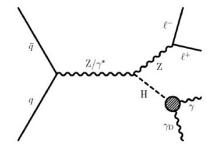
Dark Photon: ZH Channel

- Probing a Higgs portal model with dark sector
 - $H \rightarrow \gamma \gamma_d$ where γ_d is massless dark photon
 - γ_d couples to Higgs through hidden charge sector
 - M_T of photon-MET system is used as discriminating variable
 - Dominate background normalized in control region

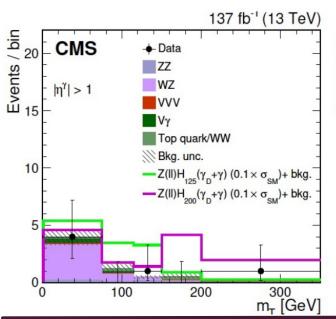
JHEP 10 (2019) 139

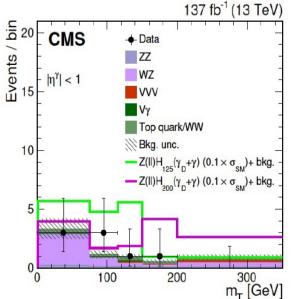
EXO-19-007

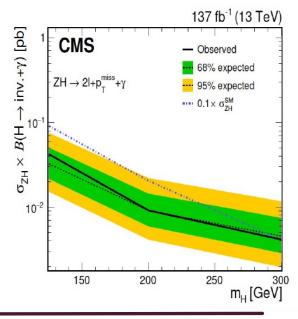
$$m_{\rm T} = \sqrt{2p_{\rm T}^{\rm miss}E_{\rm T}^{\gamma}[1-\cos(\Delta\phi_{\vec{p}_{\rm T}^{\rm miss},\vec{E}_{\rm T}^{\gamma}})]}$$



Relatively clean final state Limit on BR < 4.6% at 95% CL for SM $H(\gamma + Inv.)$







MonoHiggs bb



	0 lepton	1 muon	2 leptons	
Aim	Signal regions	$t\bar{t}$ and W+HF control region	Z+HF control region	
Fitted observable	m_h distribution	Muon charge (2 b -tag) Yields ($\geq 3 b$ -tag)	Yields	
b-tag multiplicities	resolved (small- R jets): 2, ≥ 3 merged (variable- R track jets):			
	2 (inside Higgs candidate), ≥3 (2 inside Higgs candidate)			
$E_{ m T}^{ m miss}$ proxy	$E_{ m T}^{ m miss}$	$E_{ m T, lep. invis.}^{ m miss}$	$E_{ m T, lep. invis.}^{ m miss}$	
	resolved: [150, 200], [200, 350] and [350, 500] GeV			
Bins in $E_{\rm T}^{\rm miss}$ proxy	Other merged regions: $[500, 750]$ and $[750, \infty)$ GeV			

VBF+photon + MET



ւառուբ	raking years.						
	Data-taking year	2016	2017/2018				
	Variable	VBF+ γ	Single photon	$p_{ m T}^{ m miss}$			
	Number of photons		≥1 photon				
	$p_{ m T}^{\gamma}$	>80 GeV	>230 GeV	> 80 GeV			
	Number of leptons		0				
	$p_{ m T}^{ m miss}$	>100 GeV	>140 GeV	>140 GeV			
	Jet counting		2-5				
	$m_{ m ii}$		>500 GeV				
	$ \mathring{\Delta\eta}_{ } $		>3.0				
	$\eta_{j_1} \H imes \eta_{j_2}$		<0				
	.'^. ' ~		4.0				

Table 2: Summary of the binning choice in the SRs and CRs.

Region	Bins	Range (GeV)
SR , $m_{ij} < 1500 GeV$	6	[0,30,60,90,170,250,inf]
$SR, m_{jj}^{"} \ge 1500 GeV$	6	[0,30,60,90,170,250,inf]
W + jets CR, $m_{\rm ij} < 1500{\rm GeV}$	3	[0,90,250,inf]
$W + \text{jets CR}$, $m_{\text{ij}} \ge 1500 \text{GeV}$	3	[0,90,250,inf]
$Z(\ellar\ell) + \gamma$ CR, $m_{ m ij} < 1500{ m GeV}$	1	[0,inf]
$Z(\ell \bar{\ell}) + \gamma$ CR, $m_{ii}^{"} \geq 1500\mathrm{GeV}$	1	[0,inf]
$W(\rightarrow \ell \nu) + \gamma \tilde{CR}$, $m_{ii} < 1500 \text{GeV}$	1	[0,inf]
$W(\rightarrow \ell \nu) + \gamma CR$, $m_{ii} \ge 1500 \text{GeV}$	1	[0,inf]
$\gamma + ext{jets CR}$, $m_{ ext{ii}} < 15 ilde{0}0 ext{GeV}$	1	[0,inf]
$\gamma + \mathrm{jets} \ \mathrm{CR}$, $m_{\mathrm{ij}}^{''} \geq 1500 \mathrm{GeV}$	1	[0,inf]

2HDM+Amodel parameters



The phenomenology of the model is fully determined by 14 independent parameters: the masses of the Higgs bosons h, H, A, and H^{\pm} ; the mass of the mediator a; the mass of the DM particle χ ; the Yukawa coupling strength between the mediator and the DM particle, g_{χ} ; the electroweak VEV, v; the ratio of the VEVs of the two Higgs doublets, $\tan \beta$; the mixing angles of the CP-even and CP-odd weak eigenstates, α and θ , respectively; and the three quartic couplings between the scalar doublets and the mediator $(\lambda_{P1}, \lambda_{P2}, \lambda_3)$.

The values of some of these parameters are heavily constrained by both electroweak and flavour measurements as well as phenomenological considerations, such as the requirement that the Higgs potential is stable. Further parameter choices are driven by the desire to simplify the phenomenology of the model and reduce the space of independent parameters to be scanned by experimental searches. A summary of the parameter choices and the benchmark scans shown in this note is given in the following. A detailed description of the 2HDM+*a* benchmark scans recommended by the LHC Dark Matter Working Group is given in Ref. [22].

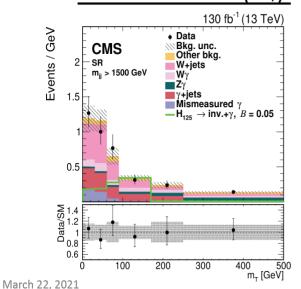
The following parameter settings are common to all benchmark scans shown in Section 6. The coupling g_{χ} is set to unity with a negligible effect on the shapes of the kinematic distributions of interest. As mentioned above, the alignment limit $(\cos(\beta - \alpha) = 0)$ is assumed, and hence $m_h = 125$ GeV and v = 246 GeV. The quartic coupling $\lambda_3 = 3$ is chosen to ensure the stability of the Higgs potential for our choice of the masses of the heavy Higgs bosons which are themselves fixed to the same value $(m_A = m_H = m_{H^{\pm}})$ to simplify the phenomenology and evade the constraints from electroweak precision measurements [21]. The other quartic couplings are also set to 3 in order to maximise the trilinear couplings between the CP-odd and the CP-even neutral states

Dark photon: VBF H + photon + p_T^{miss} _{CMS-EXO-20-005}

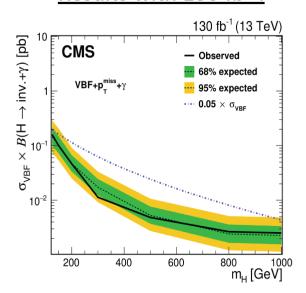


- Another way to search for DM: look for decays of a Higgs boson as the mediator to massless dark photon + SM particles
- Signature: H produced via VBF whose decay produces an isolated photon, p_Tmiss, and two forward jets
- First search for decays to undetected particle and isolated photon in the VBF channel

Transverse mass (m_T) :



Results with 130 fb⁻¹:



H-> $\gamma \gamma_d$ benchmark

mediator γ γ_D

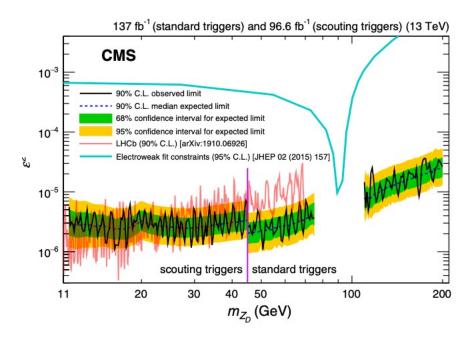
<u>Combination with analysis where H</u> <u>produced in association with a Z boson</u>

For SM-like 125 GeV H boson:

	VBF	ZH	VBF+ZH
Observed 95% CL limits	0.034	0.046	0.029
on B(H $\rightarrow \gamma \gamma_d$)			
		arXiv:1	908.02699

Dark photon model





PhysRevLett.124.131802